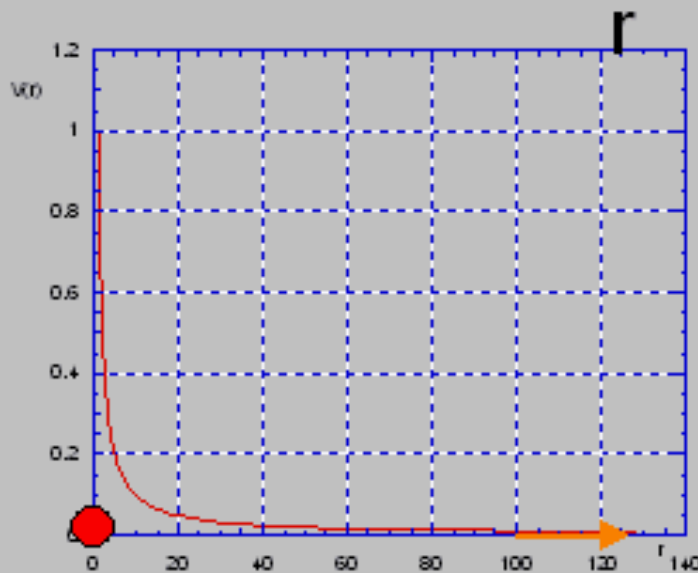


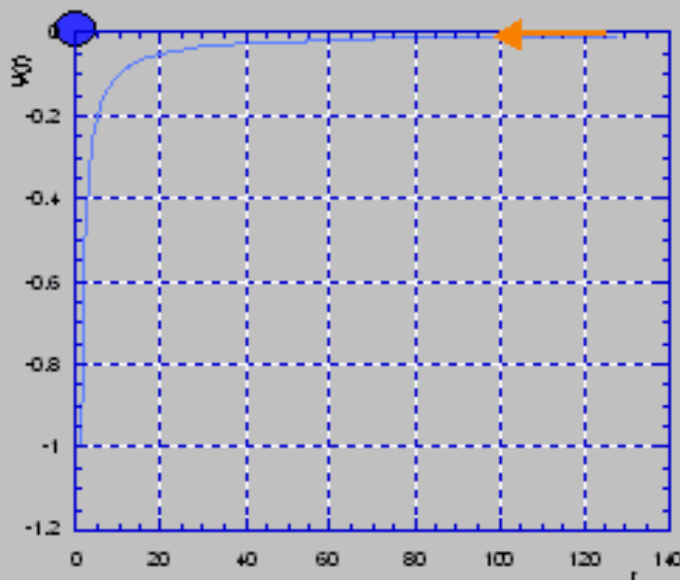
Campo elettrico e potenziale

$$\mathbf{E} = -\nabla V$$

Carica puntiforme



Carica +va:
 V decrescente
Derivata -va
con segno - OK!



Carica -va:
 V crescente
Derivata +va
con segno - OK!

Dal potenziale al campo

Es.: carica puntiforme

$$\mathbf{E} = -\nabla V$$

$$E_x = -\frac{\partial V}{\partial x}, \text{ etc}$$

$$V = \frac{q}{4\pi\epsilon_0} \frac{1}{r}, r = (x^2 + y^2 + z^2)^{1/2}$$

$$\frac{\partial V}{\partial x} = \frac{dV}{dr} \frac{\partial r}{\partial x} = -\frac{q}{4\pi\epsilon_0} \frac{1}{r^2} \frac{2x}{2(x^2 + y^2 + z^2)^{1/2}}$$

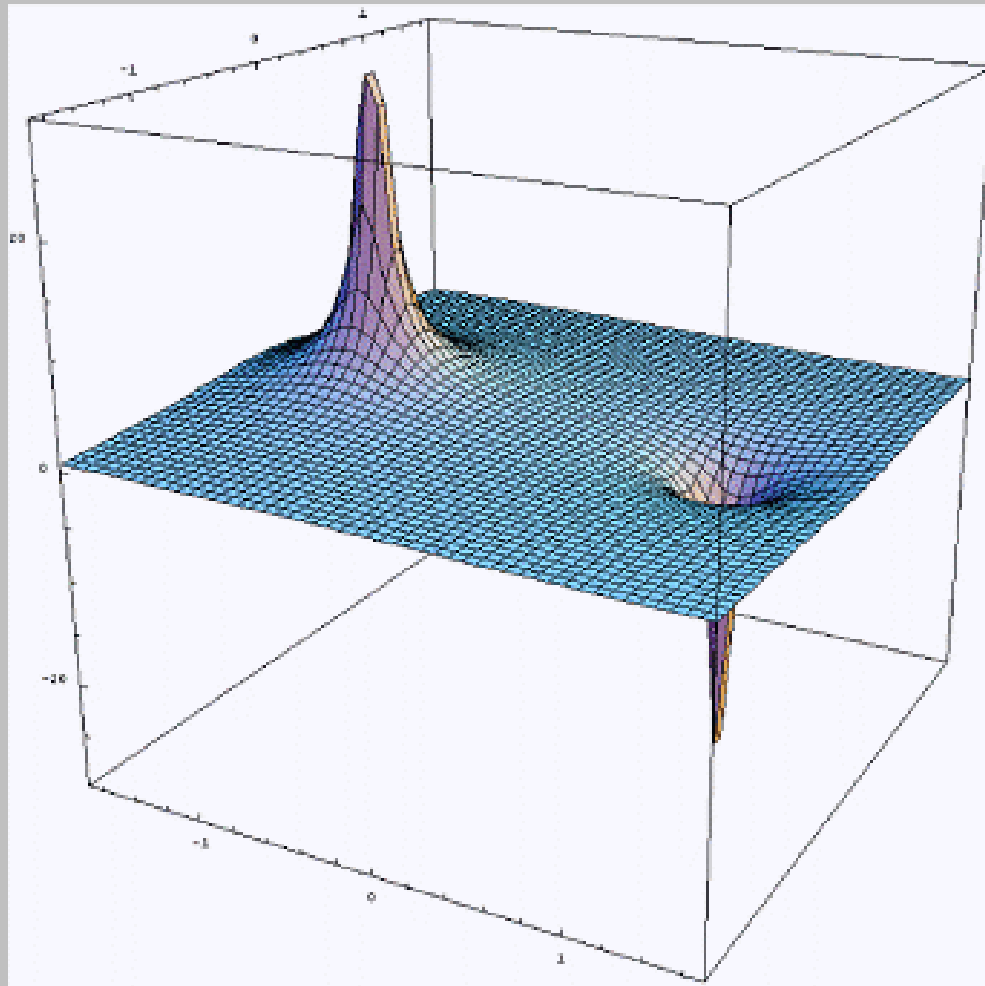
$$E_x = -\left(-\frac{q}{4\pi\epsilon_0} \frac{x}{r^3}\right), E_y, E_z \text{ analoghi}$$

$$\mathbf{E} = E_x \hat{\mathbf{i}} + E_y \hat{\mathbf{j}} + E_z \hat{\mathbf{k}} = \frac{q}{4\pi\epsilon_0} \frac{x}{r^3} \hat{\mathbf{i}} + \frac{q}{4\pi\epsilon_0} \frac{y}{r^3} \hat{\mathbf{j}} + \frac{q}{4\pi\epsilon_0} \frac{z}{r^3} \hat{\mathbf{k}}$$

$$\mathbf{E} = \frac{q}{4\pi\epsilon_0} \frac{\mathbf{r}}{r^3} = \frac{1}{4\pi\epsilon_0} \frac{q}{r^2} \hat{\mathbf{r}}$$

Campo coulombiano

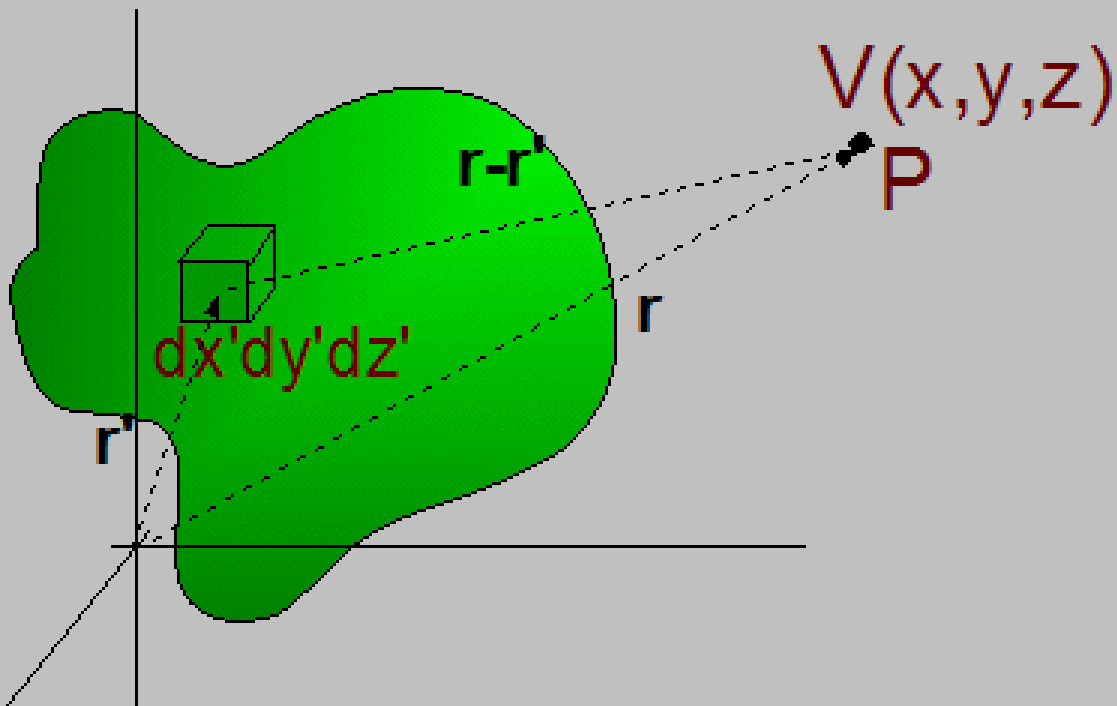
Potenziale: 2 cariche opposte



Potenziale in funzione di (x,y)

(da V.Gracco, Fis. generale II)

Potenziale: distribuzione continua



$$dq = \rho(x', y', z') dx' dy' dz'$$

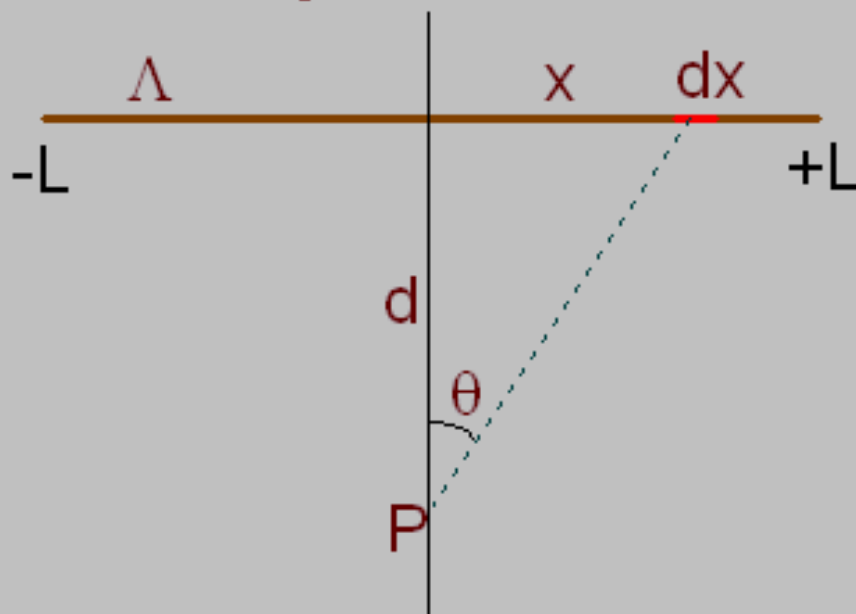
$$dV = \frac{1}{4\pi\epsilon_0} \frac{dq}{|\mathbf{r} - \mathbf{r}'|} = \frac{1}{4\pi\epsilon_0} \frac{\rho dx' dy' dz'}{|\mathbf{r} - \mathbf{r}'|}$$

$$|\mathbf{r} - \mathbf{r}'|^2 = (x - x')^2 + (y - y')^2 + (z - z')^2$$

$$V = \int dV = \int_{\text{volumen della carica}} \frac{1}{4\pi\epsilon_0} \frac{\rho dx' dy' dz'}{\left[(x - x')^2 + (y - y')^2 + (z - z')^2 \right]^{3/2}}$$

**Simile a c.elettrico, ma piu' semplice:
una sola funzione (V= scalare)**

Esempio: filo finito



$$dV = \frac{1}{4\pi\epsilon_0} \frac{dq}{r} = \frac{1}{4\pi\epsilon_0} \frac{\Lambda dx}{r}$$

$$dV = \frac{1}{4\pi\epsilon_0} \frac{\Lambda dx}{(d^2 + x^2)^{1/2}}$$

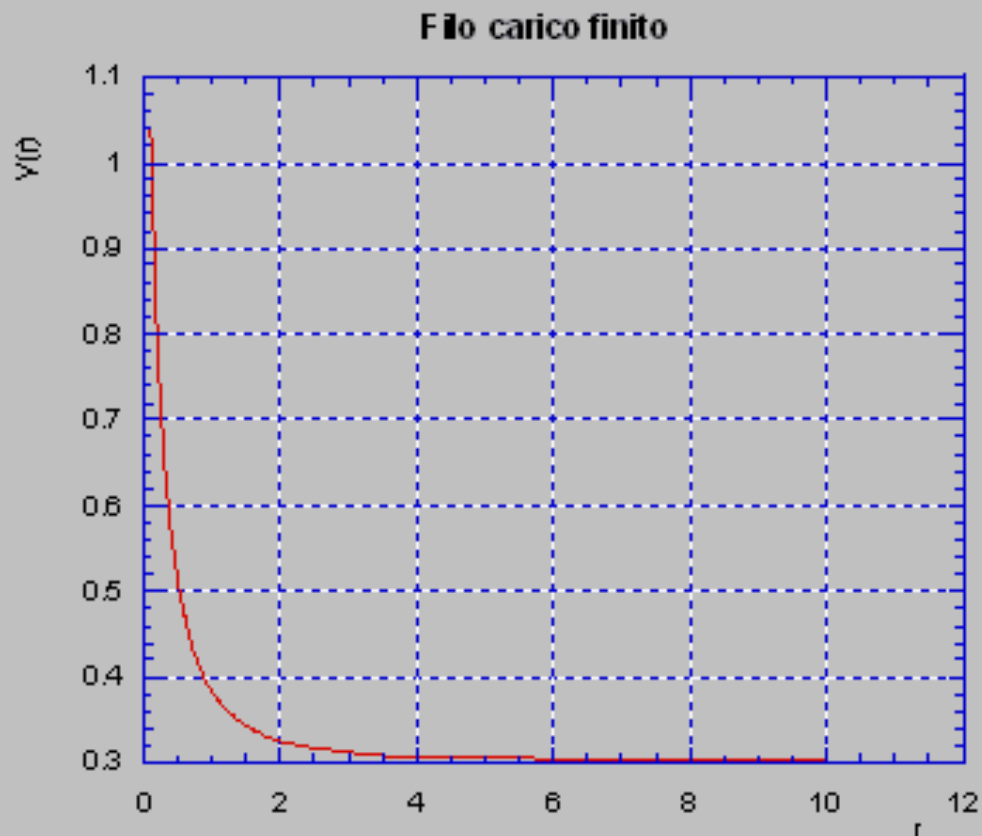
$$\rightarrow V = \frac{2\Lambda}{4\pi\epsilon_0} \int_0^L \frac{dx}{(d^2 + x^2)^{1/2}} = \frac{2\Lambda}{4\pi\epsilon_0} \ln \left(x + \sqrt{x^2 + d^2} \right) \Big|_0^L$$

$$\rightarrow V = \frac{\Lambda}{2\pi\epsilon_0} \left[\ln \left(L + \sqrt{L^2 + d^2} \right) - \ln d \right]$$

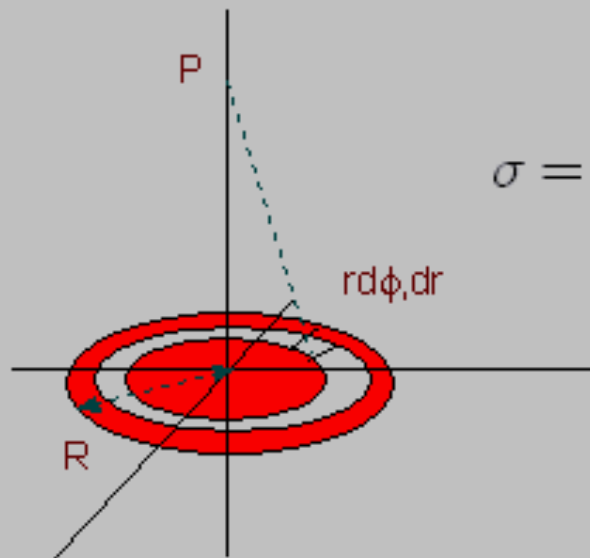
$$\rightarrow V = \frac{\Lambda}{2\pi\epsilon_0} \left[\ln \frac{L + \sqrt{L^2 + d^2}}{d} \right]$$

Filo carico finito

Grafico del potenziale: $V(r)$ vs r



Esempio: disco carico



$$\sigma = \frac{Q}{\pi R^2}$$

$$dq = \sigma dA = \sigma dr r d\varphi$$

$$u^2 = r^2 + z^2$$

$$dV = \frac{1}{4\pi\epsilon_0} \frac{dq}{u} = \frac{1}{4\pi\epsilon_0} \frac{\sigma dr r d\varphi}{(r^2 + z^2)^{1/2}}$$

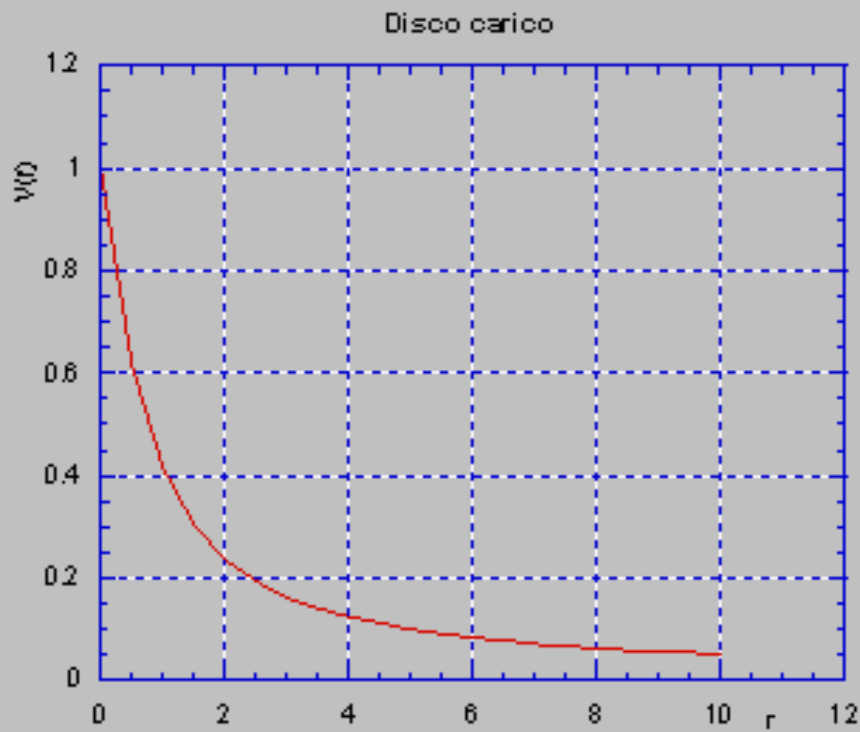
$$\rightarrow V = \frac{1}{4\pi\epsilon_0} \iint_{\text{disco}} \frac{\sigma dr r d\varphi}{(r^2 + z^2)^{1/2}}$$

$$\rightarrow V = \frac{2\pi\sigma}{4\pi\epsilon_0} \int_0^R \frac{r dr}{(r^2 + z^2)^{1/2}} =$$

$$V = \frac{\sigma}{2\epsilon_0} (r^2 + z^2)^{1/2} \Big|_0^R = \frac{\sigma}{2\epsilon_0} \left[(R^2 + z^2)^{1/2} - z \right]$$

Disco carico

Grafico del potenziale: $V(r)$ vs r



Campo di un dipolo elettrico:

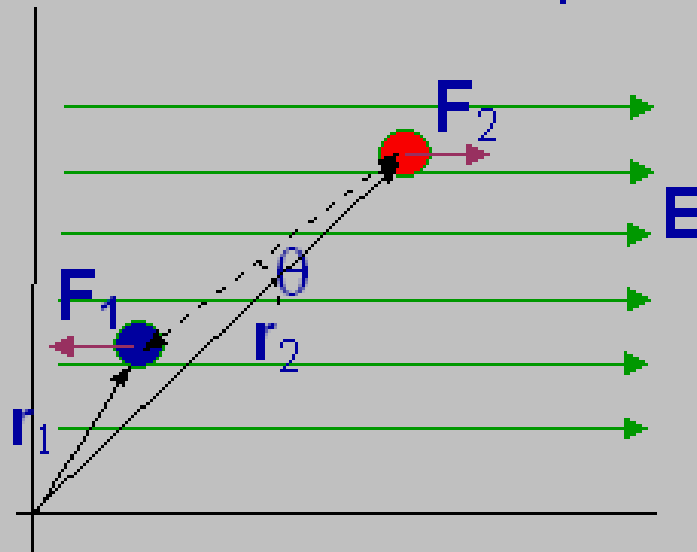
già visto a inizio lezioni per il caso particolare di punti sull'asse di simmetria

Tuttavia, interessante calcolarlo in modo generale

Per questo: v. nota Potenziali e Campi di Dipolo

Forza su un dipolo elettrico

Dipolo immerso in campo uniforme:



Forza totale = 0

Coppia: momento meccanico

$$\mathbf{M} = \mathbf{r}_1 \times \mathbf{F}_1 + \mathbf{r}_2 \times \mathbf{F}_2 = \mathbf{r}_1 \times (-\mathbf{F}) + \mathbf{r}_2 \times \mathbf{F}$$

$$\rightarrow \mathbf{M} = (\mathbf{r}_2 - \mathbf{r}_1) \times \mathbf{F} = \mathbf{a} \times (q\mathbf{E}) = (q\mathbf{a}) \times \mathbf{E} = \mathbf{p} \times \mathbf{E}$$

Posizione di equilibrio: $\mathbf{M}=0$

$$\rightarrow \mathbf{p} \parallel \mathbf{E}$$

Energia potenziale di un dipolo in un campo esterno

Spostamento angolare infinitesimo
lavoro del momento meccanico:

$$\begin{aligned}dL &= \mathbf{F}_1 \cdot d\mathbf{s}_1 + \mathbf{F}_2 \cdot d\mathbf{s}_2 \\ &= \mathbf{F}_1 \cdot \mathbf{r}_1 d\theta + \mathbf{F}_2 \cdot \mathbf{r}_2 d\theta \\ &= \mathbf{F} \cdot (\mathbf{r}_1 d\theta - \mathbf{r}_2 d\theta) = Md\theta\end{aligned}$$

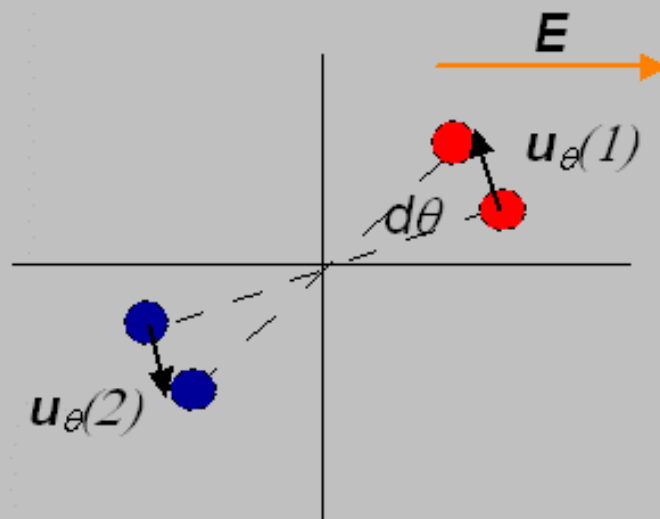
Spostamento finito da θ_0 a θ :

$$\begin{aligned}L &= \int dL = \int_{\theta_0}^{\theta} Md\theta = \int_{\theta_0}^{\theta} -pE \sin \theta d\theta \\ &\rightarrow L = pE (\cos \theta - \cos \theta_0)\end{aligned}$$

Variazione en. potenziale:

$$\begin{aligned}L &= pE (\cos \theta - \cos \theta_0) \\ L &= -\Delta U = -(U(\theta) - U(\theta_0)) \\ &\rightarrow U(\theta) = -pE \cos \theta = -\mathbf{p} \cdot \mathbf{E}\end{aligned}$$

Calcolo lavoro elementare



$$dL = dL_1 + dL_2 = \mathbf{F}_1 \cdot d\mathbf{s}_1 + \mathbf{F}_2 \cdot d\mathbf{s}_2$$

$$\mathbf{F}_1 = -\mathbf{F}_2 = \mathbf{F} = q\mathbf{E}$$

$$d\mathbf{s}_1 = |\mathbf{r}_1| d\theta \hat{\mathbf{u}}_\theta = \frac{d}{2} d\theta \hat{\mathbf{u}}_\theta(1)$$

$$d\mathbf{s}_2 = |\mathbf{r}_2| d\theta \hat{\mathbf{u}}_\theta = \frac{d}{2} d\theta \hat{\mathbf{u}}_\theta(2)$$

$$\rightarrow dL = \mathbf{F} \cdot |\mathbf{r}_1| d\theta \hat{\mathbf{u}}_\theta(1) - \mathbf{F} \cdot |\mathbf{r}_2| d\theta \hat{\mathbf{u}}_\theta(2) = \mathbf{F} \cdot [\hat{\mathbf{u}}_\theta(1) - \hat{\mathbf{u}}_\theta(2)] \frac{d}{2} d\theta$$

$$\rightarrow dL = q\mathbf{E} \cdot [\hat{\mathbf{u}}_\theta(1) - \hat{\mathbf{u}}_\theta(2)] \frac{d}{2} d\theta$$

$$\mathbf{E} \cdot \hat{\mathbf{u}}_\theta(1) = |\mathbf{E}| \cos(\theta + \pi/2) = -|\mathbf{E}| \sin \theta$$

$$\mathbf{E} \cdot \hat{\mathbf{u}}_\theta(2) = |\mathbf{E}| \cos(\theta + 3\pi/2) = +|\mathbf{E}| \sin \theta$$

$$\rightarrow dL = q(-2|\mathbf{E}| \sin \theta) \frac{d}{2} d\theta = -|\mathbf{p}||\mathbf{E}| \sin \theta d\theta = -|\mathbf{p} \times \mathbf{E}| d\theta$$

$$\rightarrow dL = -(\mathbf{p} \times \mathbf{E}) d\theta = |\mathbf{M}| d\theta$$