

Elementary Particles I

4 – Quark

Quark Model, Light and Heavy Quarks

The Quark Model

Fundamental hypothesis:

Mesons = Bound states $q\bar{q}$

Baryons, Antibaryons = Bound states $qqq, \bar{q}\bar{q}\bar{q}$

What are states $q, \bar{q}$? They are called *quark, antiquark*

Building blocks of ordinary hadrons:

A new level of structure for the hadronic matter

Quarks fill the fundamental representation of $SU(3)$

Quarks are spin $1/2$, point-like fermions

Guess:

They are never observed as free particles

The only bound states observed are $q\bar{q}, qqq, \bar{q}\bar{q}\bar{q}$ } **Why ?**

Predicting New Particles

Not a new game in town...

In the Thirties:

Pauli: *Neutrino*

Required in order to save energy, angular momentum conservation in nuclear β decay

Observed in 1956 (Reines et al., Nuclear reactor experiment)

Yukawa: *Pion*

Welcome in order to explain the general features of nuclear force

Observed in 1947 (Blackett et al., Cosmic radiation)

Quarks

Fundamental and conjugate irr.rep. of $SU(3)$: **3, 3***

Each made of 3 states

Quantum numbers: From Gell-Mann – Nishijima & $SU(3)$

$$Q = I_3 + Y/2$$

Symbol	Flavor	Spin	Q	B	S	Y	I	I ₃
<i>u</i>	<i>Up</i>	$\frac{1}{2}$	2/3	1/3	0	1/3	1/2	+1/2
<i>d</i>	<i>Down</i>	$\frac{1}{2}$	-1/3	1/3	0	1/3	1/2	-1/2
<i>s</i>	<i>Strange</i>	$\frac{1}{2}$	-1/3	1/3	-1	-2/3	0	0

} isospin doublet
isospin singlet

Quarks are predicted to carry fractional charge, baryon number!

Should they show up as free particles, would be easy to detect :

Expect unusual electromagnetic rates $\propto Q^2$

Expect bound states with fractional mass numbers $\propto B$

Mesons and Baryons

Hadrons: Expected to fill product representations

From our group theory rudiments:

$$\text{Mesons} \quad 3 \otimes 3^* = 1 \oplus 8$$

$$\text{Baryons} \quad 3 \otimes 3 \otimes 3 = 1 \oplus 8 \oplus 8 \oplus 10$$

Expect:

Nonets of mesons with given spin, parity

Singlets, octets, decuplets of baryons, as above

Quarks & Antiquarks: $\mathbf{3}$ & $\mathbf{3}^*$

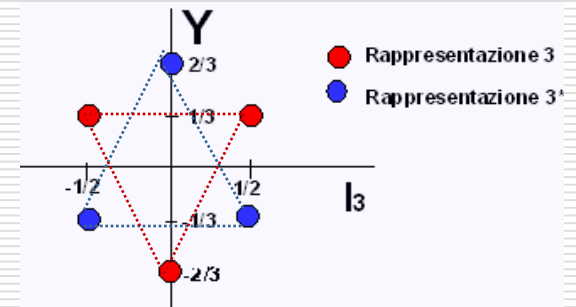
More quantum numbers

Relative space parity = -1 (Fermions)

Absolute space parity = +1 quarks, -1 antiquarks (Conventional)

Flavor	Spin	Q	B	S	Y	I	I_3
<i>Up</i>	$\frac{1}{2}$	$\frac{2}{3}$	$\frac{1}{3}$	0	$\frac{1}{3}$	$\frac{1}{2}$	$+\frac{1}{2}$
<i>Down</i>	$\frac{1}{2}$	$-\frac{1}{3}$	$\frac{1}{3}$	0	$\frac{1}{3}$	$\frac{1}{2}$	$-\frac{1}{2}$
<i>Strange</i>	$\frac{1}{2}$	$-\frac{1}{3}$	$\frac{1}{3}$	-1	$-\frac{2}{3}$	0	0

Flavor	Spin	Q	B	S	Y	I	I_3
<i>Anti-Up</i>	$\frac{1}{2}$	$-\frac{2}{3}$	$-\frac{1}{3}$	0	$-\frac{1}{3}$	$\frac{1}{2}$	$-\frac{1}{2}$
<i>Anti-Down</i>	$\frac{1}{2}$	$+\frac{1}{3}$	$-\frac{1}{3}$	0	$-\frac{1}{3}$	$\frac{1}{2}$	$+\frac{1}{2}$
<i>Anti-Strange</i>	$\frac{1}{2}$	$+\frac{1}{3}$	$-\frac{1}{3}$	+1	$+\frac{2}{3}$	0	0



A Couple of Subtle Points - I

Q: *Why are isospin 3rd components swapped for antiquarks?*

A: *Want to stick to Gell-Mann – Nishijima for them too*

$$u = \begin{pmatrix} 1 \\ 0 \end{pmatrix}, d = \begin{pmatrix} 0 \\ 1 \end{pmatrix}$$

$$Q = I_3 + \frac{Y}{2} = I_3 + \frac{B}{2}$$

Required in order to deal with $qqq, q\bar{q}, \overline{qqq}$
E.g. all present in the same process

$$Q(\bar{u}) = -\frac{2}{3} = I_3(\bar{u}) + \frac{B(\bar{u})}{2} = I_3(\bar{u}) - \frac{1}{6} \rightarrow I_3(\bar{u}) = -\frac{2}{3} + \frac{1}{6} = -\frac{1}{2}$$

$$Q(\bar{d}) = +\frac{1}{3} = I_3(\bar{d}) + \frac{B(\bar{d})}{2} = I_3(\bar{d}) - \frac{1}{6} \rightarrow I_3(\bar{d}) = +\frac{1}{3} + \frac{1}{6} = +\frac{1}{2}$$

A Couple of Subtle Points - II

Q: Why there is a -1 extra phase for u antiquark? $\begin{pmatrix} \bar{d} \\ -\bar{u} \end{pmatrix}$

A: Want to stick to same C-G coefficient for both quarks and antiquarks

Same C-G $\leftrightarrow$ Same I-spin rotation matrices

$$\begin{pmatrix} u' \\ d' \end{pmatrix} = e^{-i\tau_2\theta/2} \begin{pmatrix} u \\ d \end{pmatrix} = \begin{pmatrix} \cos(\theta/2)u - \sin(\theta/2)d \\ \sin(\theta/2)u + \cos(\theta/2)d \end{pmatrix} \quad \text{Rotation of generic state}$$

$$\begin{pmatrix} \bar{u}' \\ \bar{d}' \end{pmatrix} = e^{-i\tau_2\theta/2} \begin{pmatrix} \bar{u} \\ \bar{d} \end{pmatrix} = \begin{pmatrix} \cos(\theta/2)\bar{u} - \sin(\theta/2)\bar{d} \\ \sin(\theta/2)\bar{u} + \cos(\theta/2)\bar{d} \end{pmatrix} \leftarrow \text{Want to have this}$$

$$\rightarrow \begin{pmatrix} \bar{d}' \\ \bar{u}' \end{pmatrix} = \begin{pmatrix} \sin(\theta/2)\bar{u} + \cos(\theta/2)\bar{d} \\ \cos(\theta/2)\bar{u} - \sin(\theta/2)\bar{d} \end{pmatrix} = \begin{pmatrix} \cos(\theta/2)\bar{d} + \sin(\theta/2)\bar{u} \\ -\sin(\theta/2)\bar{d} + \cos(\theta/2)\bar{u} \end{pmatrix}$$

$$\rightarrow \begin{pmatrix} \bar{d}' \\ \bar{u}' \end{pmatrix} = \begin{pmatrix} \cos(\theta/2)\bar{d} - \sin(\theta/2)(-\bar{u}) \\ -\sin(\theta/2)\bar{d} - \cos(\theta/2)(-\bar{u}) \end{pmatrix}$$

$$\rightarrow \begin{pmatrix} \bar{d}' \\ -\bar{u}' \end{pmatrix} = \begin{pmatrix} \cos(\theta/2)\bar{d} - \sin(\theta/2)(-\bar{u}) \\ \sin(\theta/2)\bar{d} + \cos(\theta/2)(-\bar{u}) \end{pmatrix} = e^{-i\tau_2\theta/2} \begin{pmatrix} \bar{d} \\ -\bar{u} \end{pmatrix}$$

Required in order to deal with $qqq, q\bar{q}, \bar{q}q\bar{q}$
E.g. all present in the same process

Indeed, required because mesons *are* made of quark-antiquark pairs

The Light Mesons

Combine 3 quarks with 3 antiquarks: Get 9 combinations

Quark content $u\bar{d}, u\bar{s}, u\bar{u}, d\bar{u}, d\bar{s}, d\bar{d}, s\bar{u}, s\bar{d}, s\bar{s}$

Identified mesons

'State'	Q	I_3	I	S	Remarks	$J^{PC}=0^+$	$J^{PC}=1^-$	$J^{PC}=2^{++}$
$u\bar{d}$	+1	+1	1	0		π^+	ρ^+	a_2^+
$u\bar{s}$	+1	+1/2	$\frac{1}{2}$	+1		K^+	K^{+*}	K^{+**}
$u\bar{u}$	0	0	0,1	0	<i>I-spin undefined</i>	π^0, η, η'	ρ^0, ω, φ	a_2^0, f_2, f_2'
$d\bar{u}$	-1	-1	1	0		π^-	ρ^-	a_2^-
$d\bar{s}$	0	-1/2	1/2	+1		K^0	K^{0*}	K^{0**}
$d\bar{d}$	0	0	0,1	0	<i>I-spin undefined</i>	π^0, η, η'	ρ^0, ω, φ	a_2^0, f_2, f_2'
$s\bar{u}$	-1	-1/2	1/2	-1		K^-	K^{-*}	K^{-**}
$s\bar{d}$	0	+1/2	1/2	-1		$\bar{K}^0$	$\bar{K}^{0*}$	$\bar{K}^{0**}$
$s\bar{s}$	0	0	0	0		π^0, η, η'	ρ^0, ω, φ	a_2^0, f_2, f_2'

L # 0

The Light Mesons: Quantum Numbers

Build isospin eigenstates from $S=0$, $I_3=0$ states:

$$\frac{1}{\sqrt{2}}(u\bar{u} - d\bar{d}), \frac{1}{\sqrt{6}}(u\bar{u} + d\bar{d} - 2s\bar{s}), \frac{1}{\sqrt{3}}(u\bar{u} + d\bar{d} + s\bar{s})$$

Physical particles must have I defined:
 I -spin is a good symmetry

Left with 3 ambiguous states: $I_3=0$

6 unambiguous states are octet members

→ Have 2 octet, 1 singlet ambiguous

$SU(3)$ singlet: Invariant wrt $SU(3)$ rotations

$$\rightarrow \underbrace{\frac{1}{\sqrt{3}}(u\bar{u} + d\bar{d} + s\bar{s})}_{\eta_1}$$

η_1, η_8 cannot be identified with physical particles

$SU(3)$ Octets: 1 $SU(2)$ triplet, 1 $SU(2)$ singlet

$$\rightarrow \underbrace{\frac{1}{\sqrt{2}}(u\bar{u} - d\bar{d})}_{\pi^0}, \underbrace{\frac{1}{\sqrt{6}}(u\bar{u} + d\bar{d} - 2s\bar{s})}_{\eta_8}$$

The Light Mesons: Spin & Parity

$$\mathbf{J} = \mathbf{L} + \mathbf{S}$$

$$P = (-1)^{l+1}$$

$$C = (-1)^{l+s}$$

Ground state $L = 0 \rightarrow J = S$

Singlets $\rightarrow J = 0 \rightarrow P = -1, C = +1 \rightarrow J^{PC} = 0^{-+}$

Triplets $\rightarrow J = 1 \rightarrow P = -1, C = -1 \rightarrow J^{PC} = 1^{--}$

Remark 1:

Very simple and clear, but: Not covariant!

J separation into **L, S** contributions is frame dependent

→ We are assuming small quark speed: Is this correct?

Remark 2:

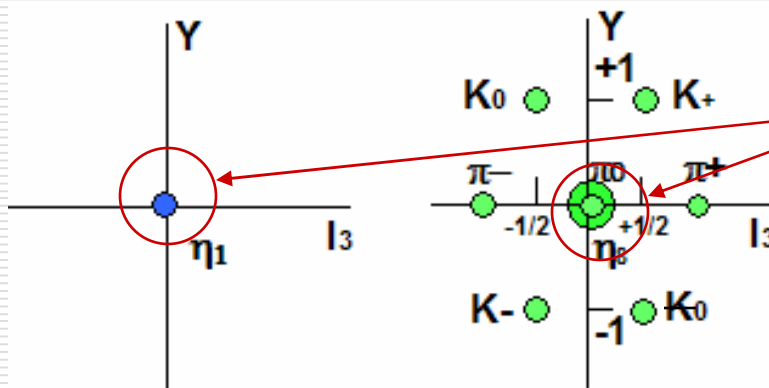
Higher spin multiplets more difficult to explain, due to orbital degrees of freedom

The Light Mesons

Particle identification with SU(3) eigenstates not always straightforward
 Example: Take pseudoscalars

$$|8;1,0\rangle = \frac{1}{\sqrt{2}}(u\bar{u} - d\bar{d}) \rightarrow \pi^0 \quad \text{Must be true because I-spin is a good symmetry}$$

$$\left. \begin{aligned} |8;0,0\rangle &= \frac{1}{\sqrt{6}}(u\bar{u} + d\bar{d} - 2s\bar{s}) \\ |1;0,0\rangle &= \frac{1}{\sqrt{2}}(u\bar{u} + d\bar{d} + s\bar{s}) \end{aligned} \right\} \begin{array}{l} \text{Not identified} \\ \text{Get some insight} \\ \text{from decay modes} \end{array}$$



π^0, η_1, η_8 Central states, $I_3 = Y = Q = 0$

The Good News

Use $SU(2)$ shift operators:

First, π^+ :

$I^- \pi^+ = \sqrt{2} \pi^0$ From definition (and multiplet diagram)

From π^+ wave function:

$$I^- \pi^+ = I^- (u\bar{d}) = d\bar{d} - u\bar{u} \Rightarrow \pi^0 = \frac{d\bar{d} - u\bar{u}}{\sqrt{2}}$$

Repeat for π^0 :

$$I^- \pi^0 = \sqrt{2} \pi^- = I^- \left(\frac{d\bar{d} - u\bar{u}}{\sqrt{2}} \right) = \frac{-d\bar{u} - d\bar{u}}{\sqrt{2}} \Rightarrow \pi^- = -d\bar{u} \quad \text{The - sign!}$$

Isosinglet (with u and d only), is η :

$$I^- \eta = I^- \left(\frac{d\bar{d} + u\bar{u}}{\sqrt{2}} \right) = \frac{-d\bar{u} + d\bar{u}}{\sqrt{2}} = 0$$

Can conclude the π^0 is an octet, don't know about η_1, η_8

The Bad News

Should $SU(3)$ be exact, all particle states would fit to irr.reps

Try to apply mass formula to mesons: Use M^2 instead of M in the mass formula, for reasons not very convincing..

Fermion Lagrangian:	m
Boson Lagrangian:	m^2

Assume the octet member is to be identified with a physical particle
Vectors

Predict

$$m_8^2 = \frac{1}{3}(4m_{K^*}^2 - m_\rho^2) \approx 0.859 \text{ GeV}^2 \quad \Leftrightarrow \quad m_\omega^2 \approx 0.613 \text{ GeV}^2, m_\phi^2 \approx 1.038 \text{ GeV}^2$$

Pseudoscalars

Predict

$$m_8^2 = \frac{1}{3}(4m_K^2 - m_\pi^2) \approx 0.321 \text{ GeV}^2 \quad \Leftrightarrow \quad m_\eta^2 \approx 0.299 \text{ GeV}^2, m_{\eta'}^2 \approx 0.918 \text{ GeV}^2$$

All in all, not very brilliant...

Breaking Everywhere

Since $SU(3)$ is broken, its eigenstates can mix:

Besides *intra*-multiplet (as before), consider *inter*-multiplet mixing

Call H_0 the $SU(3)$ symmetric part of the Hamiltonian:

$$\langle 1|H_0|1\rangle = M_1, \quad \langle 8|H_0|8\rangle = M_8$$

$SU(3)$ breaking can manifest itself in a non-diagonal, singlet-octet mass matrix:

$$M^2 = \begin{pmatrix} M_1^2 & \Delta \\ \Delta & M_8^2 \end{pmatrix}$$

By standard diagonalization find the physical masses:

$$M_{a,b}^2 = \frac{M_1^2 + M_8^2}{2} \pm \sqrt{\frac{(M_1^2 - M_8^2)^2}{4} + \Delta^2}$$

Can infer M_1, Δ :

$$M_1^2 + M_8^2 = \frac{M_a^2 + M_b^2}{2}$$

$$\Delta^2 = \frac{(M_a^2 - M_b^2)^2 - (M_1^2 - M_8^2)^2}{4}$$

How Mixing is Measured

Try to make a sense out of $SU(3)$ breaking

Simple idea: Central states of **1,8** just mix in physical particles

$$\begin{cases} |a\rangle = \sin\theta|1\rangle - \cos\theta|8\rangle \\ |b\rangle = \cos\theta|1\rangle + \sin\theta|8\rangle \end{cases}$$

'Rotation' of states: Must be unitary, phase preserving
→ Just 1 angle

Find the mixing angle:

$$\left. \begin{aligned} H|a\rangle &= M_a|a\rangle \\ H|b\rangle &= M_b|b\rangle \end{aligned} \right\} \rightarrow \left. \begin{aligned} M_a^2 &= M_1^2 - \Delta^2 \cot\theta = M_8^2 - \Delta^2 \tan\theta \\ M_b^2 &= M_1^2 + \Delta^2 \cot\theta = M_8^2 + \Delta^2 \tan\theta \end{aligned} \right\} \rightarrow \tan^2\theta = \frac{M_b^2 - M_1^2}{M_b^2 - M_8^2} = \frac{M_8^2 - M_a^2}{M_1^2 - M_a^2}$$

$\theta_p = -11^\circ$ Pseudoscalars

$\theta_v = +38^\circ$ Vectors

$\theta_T = +32^\circ$ Tensors

Best observed in vector mesons:

Pointing to: $m_u \sim m_d$, $m_s > m_{u,d}$

$m_\omega \approx u\bar{u} + d\bar{d} \rightarrow \omega = 1/\sqrt{2}(u\bar{u} + d\bar{d})$ ρ, ω only u, d quarks: OK mass degenerate

$m_\varphi \approx s\bar{s} \rightarrow \varphi = s\bar{s}$ φ only s quarks: OK decays modes

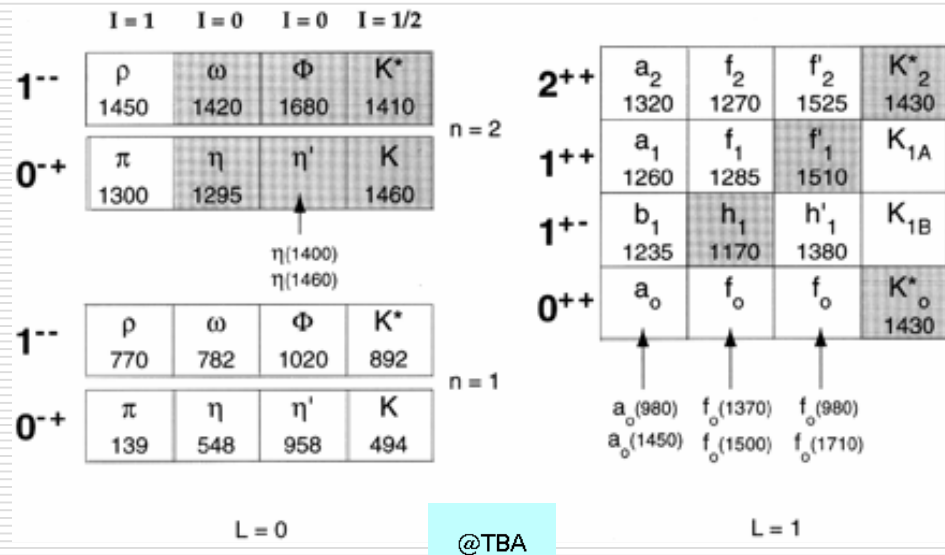
Higher Spin Mesons

Combine non relativistically L, S :

L	S	State	J^{PC}
0	0	1S_0	0^{-+}
0	1	3S_1	1^{--}
1	0	1P_1	1^{-+}
1	1	3P_0	0^{++}
1	1	3P_1	1^{++}
1	1	3P_2	2^{++}
2	0	1D_2	2^{-+}
2	1	3D_1	1^{--}
2	1	3D_2	2^{--}
2	1	3D_3	3^{--}

Remarks:

*States in grey can mix
C is meant for $Q=S=0$*



The Light Baryons - I

Combine 3 quarks: Get $3 \times 3 \times 3 = 27$ combinations
 But: Only 10 different quark contents

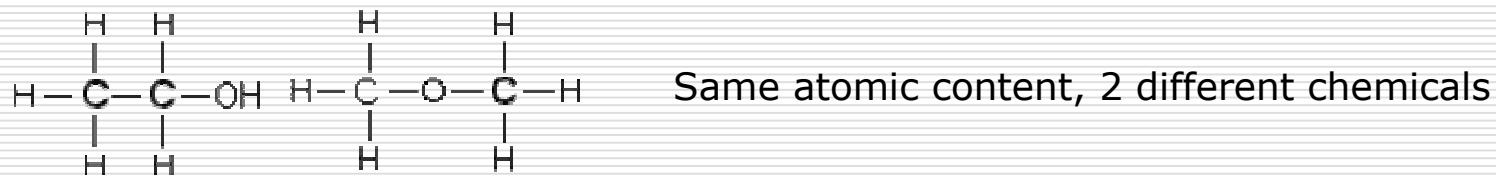
$3 + 3 \cdot 2 + 1 = 10$: $uuu, ddd, sss, uud, uus, ddu, dds, ssu, ssd, uds$

Remember:

Same composition does not imply same quantum state

Somewhat similar to difference between *raw* and *structural* formulae

Examples:



$\frac{1}{\sqrt{2}}(|pn\rangle + |np\rangle)$ symmetric
 No bound states

Same nuclear content, 2 different states

$\frac{1}{\sqrt{2}}(|pn\rangle - |np\rangle)$ antisymmetric
 One bound state

The Light Baryons - II

SU(3) Multiplets: **1, 8, 8, 10**

Reminder:

What about different quark masses?
Well, that's all out of SU(3) *breaking*..

Quarks of different flavor to be taken as *different states of identical particles* (like electrons with spin up, down)

→ Multi-quark states expected to have definite *exchange symmetry*

Can derive flavor exchange symmetry of each multiplet

1 – Singlet
Fully antisymmetric

8 – Two Octets
Undefined symmetry

10 – Decuplet
Fully symmetric

Great fun with $SU(3)$ and S_3 !

Special Unitary in 3 D
Symmetric Group of 3 objects } Tightly bound...

The Moral:

Learn some group theory while still young..

The Light Baryons - III

Now look at the remaining part of the wave function:

$$|a\rangle = |space\rangle |spin\rangle |flavor\rangle$$

NB: This expression is incomplete! See later

Space: Expect S-Wave $\rightarrow$ *Symmetric*

Difficult to guess an effective potential originating a ground state with $L \neq 0$

Spin: Quarks are Fermions

Combine 3 spin $1/2$:

$$1/2 \oplus 1/2 = \begin{cases} 0 \rightarrow 0 \oplus 1/2 = 1/2 & 2 \text{ sub-states} \\ 1 \rightarrow 1 \oplus 1/2 = 1/2, 3/2 & 2+4 \text{ sub-states} \end{cases}$$

$\rightarrow$ Expect 1 quartet, 2 doublets

$$\left. \begin{aligned} |3/2, +3/2\rangle &= (\uparrow\uparrow\uparrow), & |3/2, -3/2\rangle &= (\downarrow\downarrow\downarrow) \\ |3/2, +1/2\rangle &= 1/\sqrt{3}(\uparrow\uparrow\downarrow + \uparrow\downarrow\uparrow + \downarrow\uparrow\uparrow), & |3/2, -1/2\rangle &= 1/\sqrt{3}(\downarrow\downarrow\uparrow + \downarrow\uparrow\downarrow + \uparrow\downarrow\downarrow) \end{aligned} \right\} \text{Quartet - Symmetric}$$

$$|1/2, +1/2\rangle_A = 1/\sqrt{2}(\uparrow\downarrow - \downarrow\uparrow)\uparrow, \quad |1/2, -1/2\rangle_A = 1/\sqrt{2}(\uparrow\downarrow - \downarrow\uparrow)\downarrow \quad \text{Doublet - Antisymmetric 1-2}$$

$$|1/2, +1/2\rangle_S = 1/\sqrt{2}\uparrow(\uparrow\downarrow - \downarrow\uparrow), \quad |1/2, -1/2\rangle_S = 1/\sqrt{2}\downarrow(\uparrow\downarrow - \downarrow\uparrow) \quad \text{Doublet - Antisymmetric 2-3}$$

The Light Baryons - IV

Can use another bit of group theory to write succinctly:

$$2 \otimes 2 \otimes 2 = 4 \oplus 2_S \oplus 2_A \quad \text{spin}$$

$$3 \otimes 3 \otimes 3 = 1 \oplus 8_S \oplus 8_A \oplus 10 \quad \text{flavor}$$

Summary of flavor, spin symmetry of different representations:

SU(3)	Flavor	Symmetry	Spin	Symmetry	SU(2)
	10_S	S	4_S	S	
	$8_{M,S}$	<i>n.a.; symmetric 1-2</i>	$2_{M,S}$	<i>n.a.; symmetric 1-2</i>	
	$8_{M,A}$	<i>n.a.; antisymmetric 1-2</i>	$2_{M,A}$	<i>n.a.; antisymmetric 1-2</i>	
	1_A	A			

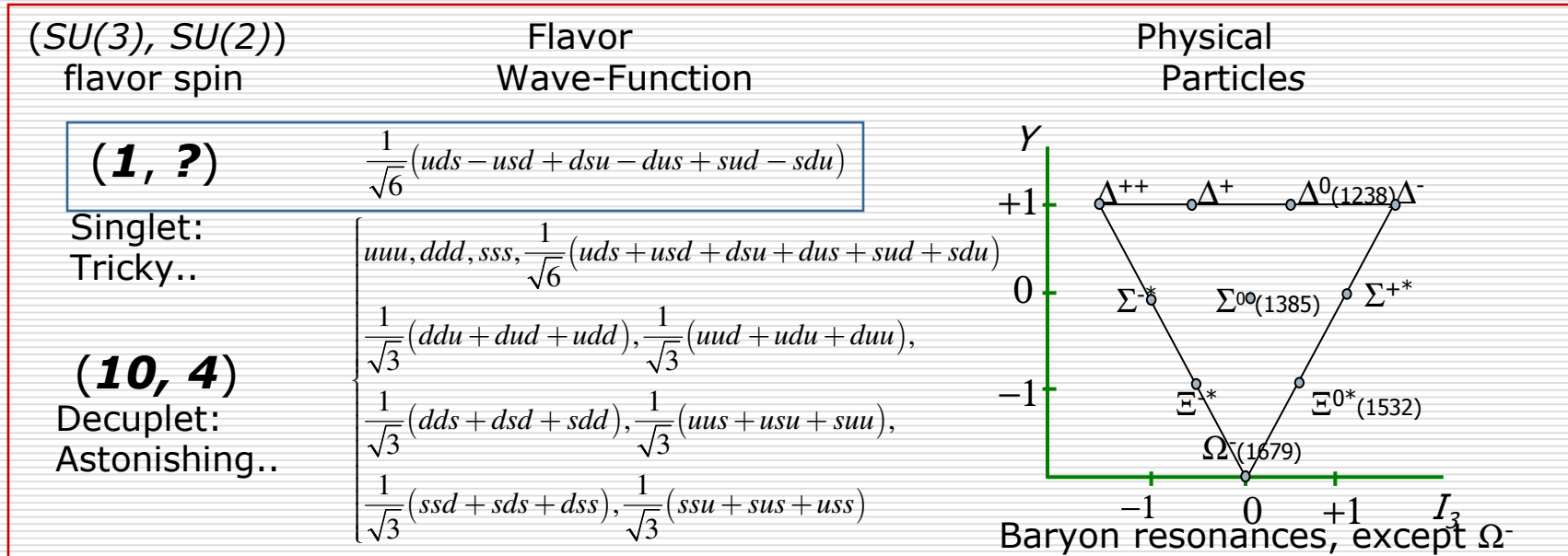
Now combine flavor *and* spin:

S, A, M referring to *flavor*spin*

	10_S	$8_{M,S}$	$8_{M,A}$	1_A
4_S	$(10, 4) \ S$	$(8, 4) \ M$	$(8, 4) \ M$	$(1, 4) \ A$
$2_{M,S}$	$(10, 2) \ M$	$(8, 2) \ M$	$(8, 2) \ M$	$(1, 2) \ M$
$2_{M,A}$	$(10, 2) \ M$	$(8, 2) \ M$	$(8, 2) \ M$	$(1, 2) \ M$

The Light Baryons: Singlet, Decuplet

Observed multiplets



Most unexpected:

Total wave function appears to be exchange symmetric for decuplet!

Would expect it *anti-symmetric* for a bundle of identical fermions

Are we forgetting something in this game?

The Light Baryons: Octet - I

Assume a globally *symmetric* wave-function for octet too
 Would be very difficult to account for a multiplet-dependent symmetry!
 Guess the symmetric spin-flavor part:

Antisymmetric 1↔2

Flavor: Two sets, 8 states each

Antisymmetric 2↔3

$\left. \begin{aligned} &\frac{1}{\sqrt{2}}(ud - du)d, \frac{1}{\sqrt{2}}(ud - du)u, \\ &\frac{1}{\sqrt{2}}(ds - sd)d, \frac{1}{\sqrt{2}}(ds - sd)s, \\ &\frac{1}{\sqrt{2}}(us - su)u, \frac{1}{\sqrt{2}}(us - su)s, \\ &\frac{1}{2}[(us - su)d + (ds - sd)u], \\ &\frac{1}{\sqrt{12}}[2(ud - du)s + (us - su)d - (ds - sd)u] \end{aligned} \right\} \varphi_{A12}^{(i)}, i = 1, 8$	$\left. \begin{aligned} &\frac{1}{\sqrt{2}}d(ud - du), \frac{1}{\sqrt{2}}u(ud - du), \\ &\frac{1}{\sqrt{2}}d(ds - sd), \frac{1}{\sqrt{2}}s(ds - sd), \\ &\frac{1}{\sqrt{2}}u(us - su), \frac{1}{\sqrt{2}}s(us - su), \\ &\frac{1}{2}[d(us - su) + u(ds - sd)], \\ &\frac{1}{\sqrt{12}}[2s(ud - du) + d(us - su) - u(ds - sd)] \end{aligned} \right\} \varphi_{A23}^{(i)}, i = 1, 8$
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Spin: Two sets, 2 states each

$\left. \begin{aligned} & 1/2, +1/2\rangle_A = 1/\sqrt{2}(\uparrow\downarrow - \downarrow\uparrow)\uparrow \\ & 1/2, -1/2\rangle_A = 1/\sqrt{2}(\uparrow\downarrow - \downarrow\uparrow)\downarrow \end{aligned} \right\} \chi_{A12}^{(j)}, j = 1, 2$	$\left. \begin{aligned} & 1/2, +1/2\rangle_S = 1/\sqrt{2}\uparrow(\uparrow\downarrow - \downarrow\uparrow) \\ & 1/2, -1/2\rangle_S = 1/\sqrt{2}\downarrow(\uparrow\downarrow - \downarrow\uparrow) \end{aligned} \right\} \chi_{A23}^{(j)}, j = 1, 2$
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The Light Baryons: Octet - II

Can also write down a 3rd set of 8 (flavor) or 2 (spin) wave functions, antisymmetric wrt $1 \leftrightarrow 3$: $\varphi_{A13}^{(i)}, i=1,8, \chi_{A13}^{(j)}, j=1,2$ Not independent from the others

The question:

What is the spin-flavor wave function of, say, a proton with spin up?

The answer:

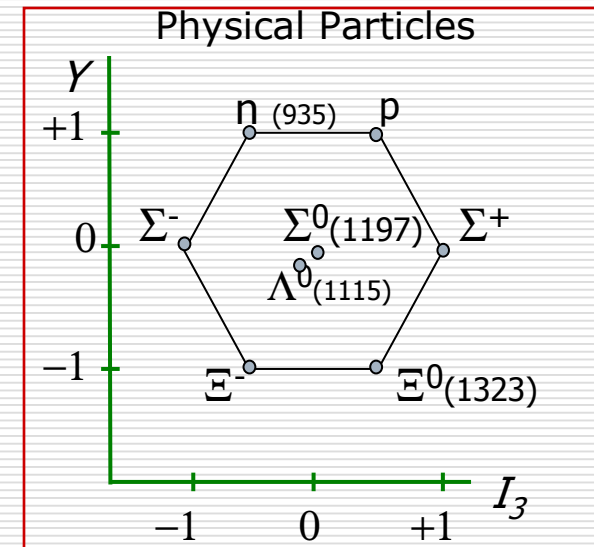
Must consider all symmetric spin-flavor products with the proper quark content and s_z

Say the appropriate functions are n.2 (flavor) and n.1 (spin)

$$\varphi = \begin{cases} \varphi_{A12}^{(2)} = \frac{1}{\sqrt{2}}(ud - du)u \\ \varphi_{A23}^{(2)} = \frac{1}{\sqrt{2}}u(ud - du) \\ \varphi_{A13}^{(2)} = \frac{1}{\sqrt{2}}(uud - duu) \end{cases} \quad \chi = \begin{cases} \chi_{A12}^{(1)} = \frac{1}{\sqrt{2}}(\uparrow\downarrow - \downarrow\uparrow)\uparrow \\ \chi_{A23}^{(1)} = \frac{1}{\sqrt{2}}\uparrow(\uparrow\downarrow - \downarrow\uparrow) \\ \chi_{A13}^{(1)} = \frac{1}{\sqrt{2}}(\uparrow\uparrow\downarrow - \downarrow\uparrow\uparrow) \end{cases}$$

Products:

$$\begin{cases} \varphi_{A12}^{(2)}\chi_{A12}^{(1)} = \frac{1}{\sqrt{2}}(ud - du)u \frac{1}{\sqrt{2}}(\uparrow\downarrow - \downarrow\uparrow)\uparrow \\ \varphi_{A23}^{(2)}\chi_{A23}^{(1)} = \frac{1}{\sqrt{2}}u(ud - du) \frac{1}{\sqrt{2}}\uparrow(\uparrow\downarrow - \downarrow\uparrow) \\ \varphi_{A13}^{(2)}\chi_{A13}^{(1)} = \frac{1}{\sqrt{2}}(uud - duu) \frac{1}{\sqrt{2}}(\uparrow\uparrow\downarrow - \downarrow\uparrow\uparrow) \end{cases}$$



The Light Baryons: Octet - III

In order to get a fully exchange-symmetric state, must take a linear combination of *all* contributions

$$|p, +1/2\rangle = \sum_{k=A12}^{A13} \varphi_k^{(2)} \chi_k^{(1)}$$

Unconvinced? Still it's true, because any sum of 2 products is equivalent to the 3rd...

$$|p, +1/2\rangle = \frac{1}{\sqrt{3}} \left[\frac{1}{\sqrt{2}} (ud - du) u \frac{1}{\sqrt{2}} (\uparrow\downarrow - \downarrow\uparrow) \uparrow + \frac{1}{\sqrt{2}} u (ud - du) \frac{1}{\sqrt{2}} \uparrow (\uparrow\downarrow - \downarrow\uparrow) + \frac{1}{\sqrt{2}} (uud - duu) \frac{1}{\sqrt{2}} (\uparrow\uparrow\downarrow - \downarrow\uparrow\uparrow) \right]$$

Finally: The *proton, spin up* wave function!

$$|p, +1/2\rangle = N \left(\begin{array}{c} 2u\uparrow d\downarrow u\uparrow + 2u\uparrow u\uparrow d\downarrow + 2d\downarrow u\uparrow u\uparrow \\ -u\downarrow d\uparrow u\uparrow - d\uparrow u\downarrow u\uparrow - u\uparrow u\downarrow d\uparrow - u\uparrow d\uparrow u\downarrow - u\downarrow u\uparrow d\uparrow - d\uparrow u\uparrow u\downarrow \end{array} \right)$$

N = Normalization constant

$$N = \frac{1}{\sqrt{6 \cdot 1^2 + 3 \cdot 2^2}} = \frac{1}{\sqrt{6+12}} = \frac{1}{\sqrt{18}}$$

The Light Baryons: $SU(6)$

Flavor symmetry: $SU(3)$

Spin 1/2 : $SU(2)$

→ Total symmetry: $SU(3) \otimes SU(2)$

Can think of extending to a larger group, $SU(6)$: Giant symmetry

$SU(6)$ includes $SU(3) \otimes SU(2)$ as a *subgroup*

[Just meaning $SU(6)$ has *extra* transformations wrt $SU(3) \otimes SU(2)$:

Generic $SU(6)$ operation can mix states sitting in *different* (flavor, spin) multiplets

Generic $SU(3) \otimes SU(2)$ operation only mixes states sitting in the *same* (flavor, spin) multiplet]

Example: Mesons

$$6 \otimes 6^* = 1 \oplus 35$$

$SU(3) \otimes SU(2)$ content:

$$35 = \underbrace{\{1, 3\}}_{3 \text{ states}} \oplus \underbrace{\{8, 1\}}_{8 \text{ states}} \oplus \underbrace{\{8, 3\}}_{24 \text{ states}}$$

Baryons

$$6 \otimes 6 \otimes 6 = 20 \oplus 56 \oplus 70 \oplus 70$$

$SU(3) \otimes SU(2)$ content:

$$56 = \underbrace{\{8, 2\}}_{16 \text{ states}} \oplus \underbrace{\{10, 4\}}_{40 \text{ states}} \text{ Symmetric!}$$

Observe: Situation similar to $SU(3)$ vs $SU(2) \otimes U(1)$

Different $SU(2)$ multiplets grouped into a single $SU(3)$ supermultiplet

Besides exchanging states within each $SU(2)$ multiplet, can exchange states among different $SU(2)$ multiplets, within the same $SU(3)$ representation

Summary: Decuplet

<i>State</i>	<i>Q</i>	<i>I</i> ₃	<i>I</i>	<i>S</i>	<i>J^{PC}=3/2⁺</i>
<i>uuu</i>	+2	+3/2	3/2	0	Δ^{++}
$1/\sqrt{3}(uud + udu + duu)$	+1	+1/2	3/2	0	Δ^{+}
$1/\sqrt{3}(udd + dud + duu)$	0	-1/2	3/2	0	Δ^0
<i>ddd</i>	-1	-1/2	3/2	0	Δ^{-}
$1/\sqrt{3}(uus + usu + suu)$	+1	+1	1	-1	Σ^{*+}
$1/\sqrt{6}(uds + sud + dsu + sdu + dus + usd)$	0	0	1	-1	Σ^{*0}
$1/\sqrt{3}(dds + +dsd + sdd)$	-1	-1	1	-1	Σ^{*-}
$1/\sqrt{3}(uss + sus + ssu)$	0	+1/2	1/2	-2	Ξ^{*0}
$1/\sqrt{3}(dss + sds + ssd)$	-1	-1/2	1/2	-2	Ξ^{*-}
<i>sss</i>	-1	0	0	-3	Ω^{-}

Wave functions

Summary: Octet

<i>Quarks</i>	Q	I_3	I	S	$J^C=1/2^+$
uud	$+1$	$+1/2$	$1/2$	0	p
udd	0	$-1/2$	$1/2$	0	n
dds	-1	-1	1	-1	Σ^-
uds	0	0	$1,0$	-1	Σ^0, Λ^0
uus	$+1$	$+1$	1	-1	Σ^+
dss	-1	$-1/2$	$1/2$	-2	Ξ^-
uss	0	$+1/2$	$1/2$	-2	Ξ^0

Quark content only
(no wave function)

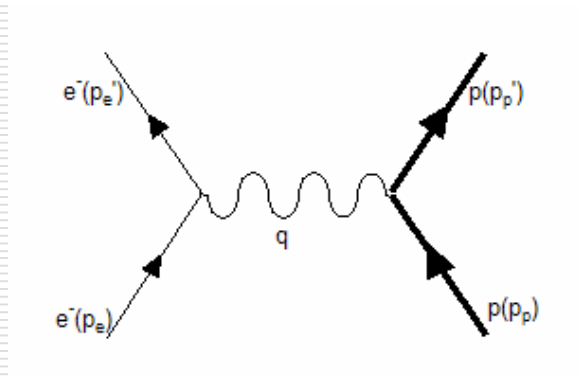
Tired of Symmetry, Multiplets etc?

Go for some dynamics...

Examine first electron-positron bound states: *Positronium*
Somewhat similar to mesons: *Particle-antiparticle bound state*

Can be dealt with by use of non-relativistic potential models
Useful insight by connecting to perturbation theory..

Start first from electron-proton interaction:



$$T_{fi} = e^2 \frac{\left(\bar{u}(p_e') \gamma^\mu u(p_e) \right) \left(\bar{u}(p_p') \gamma_\mu u(p) \right)}{q^2}$$

The e - p Effective Interaction - I

Expand matrix element to low speed approximation

$$T_{fi} \simeq -\frac{e^2}{q^2} \left[1 - \frac{\mathbf{p}_e'^2 + \mathbf{p}_e'^2}{8m_e^2} \right] \left[1 - \frac{\mathbf{p}_p'^2 + \mathbf{p}_p'^2}{8m_p^2} \right].$$

$$\left\{ \underbrace{\tilde{\chi}^{\dagger} \left[1 + \frac{\mathbf{p}_p' \cdot \mathbf{p}_p + i\boldsymbol{\sigma} \cdot (\mathbf{p}_p' \times \mathbf{p}_p)}{4m_p^2} \right]}_{\text{time section, p 4-current}} \underbrace{\tilde{\chi} \chi^{\dagger} \left[1 + \frac{\mathbf{p}_e' \cdot \mathbf{p}_e + i\boldsymbol{\sigma} \cdot (\mathbf{p}_e' \times \mathbf{p}_e)}{4m_e^2} \right]}_{\text{time section, e 4-current}} \chi + \right. \\ \left. - \underbrace{\tilde{\chi}^{\dagger} \left[\frac{\mathbf{p}_p' + \mathbf{p}_p - i\boldsymbol{\sigma} \times (\mathbf{p}_p' - \mathbf{p}_p)}{2m_p} \right]}_{\text{space section, p 4-current}} \underbrace{\tilde{\chi} \cdot \chi^{\dagger} \left[\frac{\mathbf{p}_e' + \mathbf{p}_e - i\boldsymbol{\sigma} \times (\mathbf{p}_e' - \mathbf{p}_e)}{2m_e} \right]}_{\text{space section, e 4-current}} \chi \right\}$$

Get a non-relativistic matrix element, where χ, χ' are 2-dimensional (Pauli) spinors for electron and proton

The Bottom Line:

At low speed/energy we can neglect radiation, pair production (real & virtual)

→Left with corrections:

*Relativistic Energy/Momentum
Magnetic Moments
More*

The e - p Effective Interaction - II

Reminder from potential scattering:

Transition amplitude = Fourier transform of the effective potential

→ Get effective e - p potential by anti-transforming the amplitude

$$V_C = -\frac{e^2}{r} (\tilde{\chi}^\dagger \tilde{\chi}) (\chi^\dagger \chi) \quad \text{Coulomb term}$$

$$V_{so} = \frac{e^2}{4m_e^2 r^3} (\tilde{\chi}^\dagger \tilde{\chi}) (\chi^\dagger \boldsymbol{\sigma} \cdot \mathbf{L} \chi) \quad \text{Spin-orbit}$$

$$V_D = \frac{e^2}{8m_e^2} 4\pi \delta^{(3)}(\mathbf{r}) (\tilde{\chi}^\dagger \tilde{\chi}) (\chi^\dagger \chi) \quad \text{'Darwin term'}$$

Fine structure terms

$$V_{dip-dip} = \underbrace{\frac{8\pi e^2}{3m_e m_p} g_p \mathbf{s}_e \cdot \mathbf{s}_p \delta^{(3)}(\mathbf{r})}_{\text{Hyperfine structure}} + \underbrace{\frac{e^2}{m_e m_p r^3} g_p \left[3(\mathbf{s}_e \cdot \mathbf{r})(\mathbf{s}_p \cdot \mathbf{r}) - \mathbf{s}_e \cdot \mathbf{s}_p \right]}_{\text{Tensor interaction}} \quad \text{Dipole-dipole interaction}$$

Valid for S states

Astonishing: Everything included in our modest 1-photon diagram...

The e - p Effective Interaction - III

Effect of hyperfine interaction on ground state energy:

$$\left\langle \frac{8\pi e^2}{3m_e m_p} g_p \mathbf{s}_e \cdot \mathbf{s}_p \delta^{(3)}(\mathbf{r}) \right\rangle_{\text{ground state}}$$

$$\rightarrow \Delta E_{\text{hyp}} = \frac{8\pi e^2}{3m_e m_p} g_p \cdot \frac{1}{2} [j(j+1) - s_e(s_e+1) - s_p(s_p+1)] \cdot |\psi(0)|^2$$

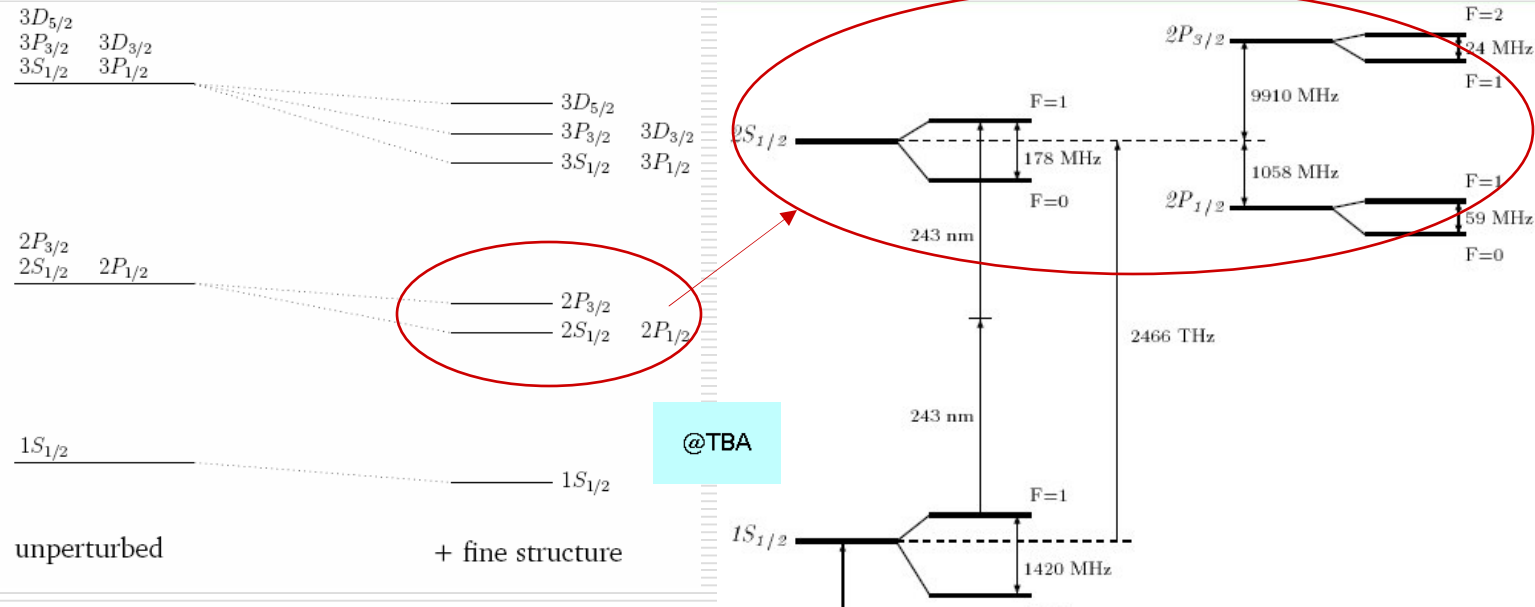
$$|\psi(0)|^2 = \frac{(m_e \alpha)^3}{\pi} \rightarrow \Delta E_{\text{hyp}} = \frac{8\pi e^2}{3m_e m_p} g_p \frac{1}{2} \frac{(m_e \alpha)^3}{\pi} \left[j(j+1) - \frac{3}{4} - \frac{3}{4} \right]$$

$$\rightarrow \Delta E_{\text{hyp}} = \frac{4}{3} g_p \frac{e^2}{m_e m_p} \left[j(j+1) - \frac{3}{2} \right] (m_e \alpha)^3$$

$$\Delta E_{\text{hyp}} = \begin{cases} \frac{4}{3} g_p \frac{e^2 m_e}{m_p} (2 - 3/2) (m_e \alpha^3) = \frac{4}{3} g_p \frac{m_e}{m_p} \frac{1}{2} (m_e \alpha^4) & \text{hyperfine shift -triplet} \\ \frac{4}{3} g_p \frac{2e^2 m_e}{m_p} (0 - 3/2) (m_e \alpha^3) = -\frac{12}{3} g_p \frac{1}{2} \frac{m_e}{m_p} (m_e \alpha^4) & \text{hyperfine shift - singlet} \end{cases}$$

$$\rightarrow \Delta(\Delta E_{\text{hyp}})_{\text{triplet-singlet}} = \frac{8}{3} g_p \frac{m_e}{m_p} (m_e \alpha^4)$$

Hydrogen: Fine & Hyperfine Structure



$$\Delta E_{l,1/2;j,m_j} = E_n \frac{\alpha^2}{n^2} \left(\frac{n}{j+1/2} - \frac{3}{4} \right)$$

Fine structure:
Spin-Orbit+Relativistic+Darwin
Splits j sublevels

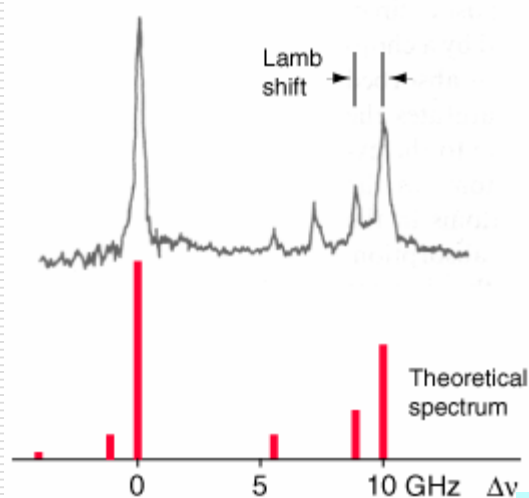
$$\Delta E = \frac{\mu_0 g_p e^2}{8\pi m_p m_e} \left\langle \frac{3 (\mathbf{S}_p \cdot \hat{\mathbf{r}}) (\mathbf{S}_e \cdot \hat{\mathbf{r}}) - \mathbf{S}_p \cdot \mathbf{S}_e}{r^3} \right\rangle + \frac{\mu_0 g_p e^2}{3 m_p m_e} \langle \mathbf{S}_p \cdot \mathbf{S}_e \rangle |\psi(0)|^2$$

Hyperfine structure:
Dipole-Dipole
Splits F sublevels
 $\mathbf{F} = \mathbf{j}_e + \mathbf{s}_p$

Total angular momentum electron+ proton

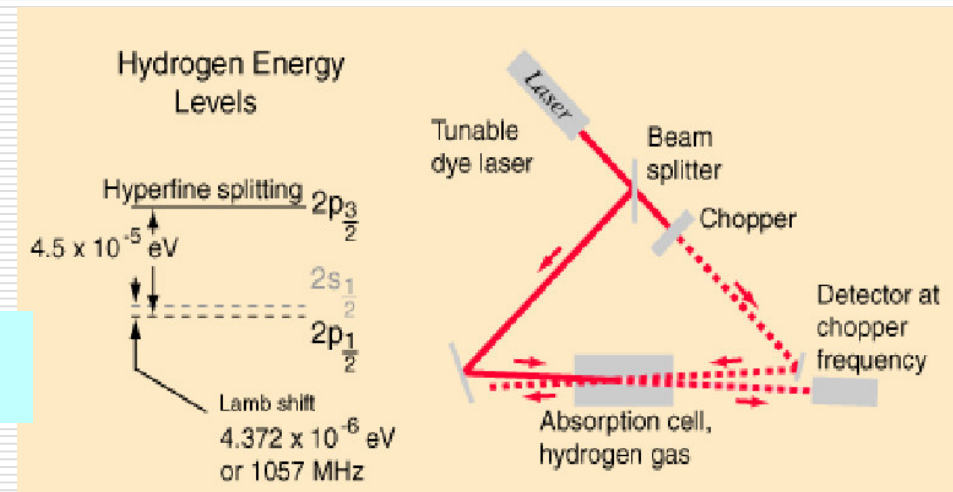
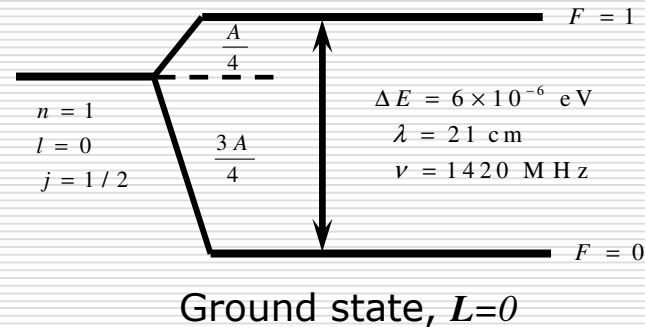
Hyperfine Splitting of Hydrogen

Ground state



Hydrogen fine structure and hyperfine structure for the $n=3 \rightarrow 2$ transition.
(After Ohanian, *Modern Physics*, Ch 7.,
spectrum from T. W. Hansch, Stanford Univ.)

@TBA



The 21 cm Line: A Cosmic Tune

Important radio-astronomical tool:
*Mapping the Hydrogen content of
the Universe*

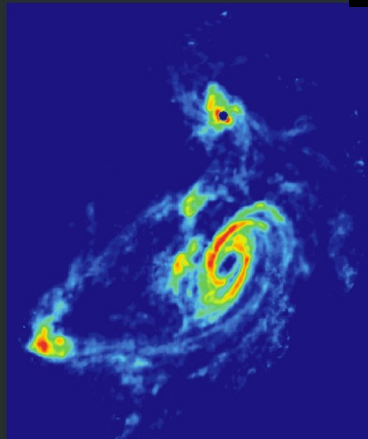
@TBA

TIDAL INTERACTIONS IN M81 GROUP

Stellar Light Distribution



21 cm HI Distribution



The Universe seen at 21 cm

Lots of physics and cosmology..

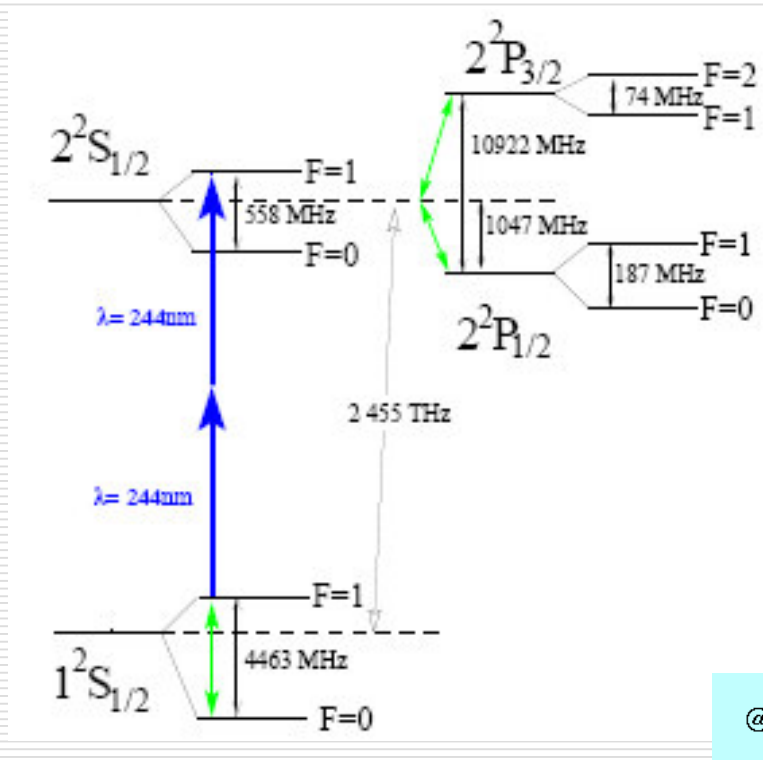
Example:

How is the transition excited?

A measurement of the galactic/
intergalactic temperature

Muonium

μ^+e^- 'atom'



Similar to Hydrogen

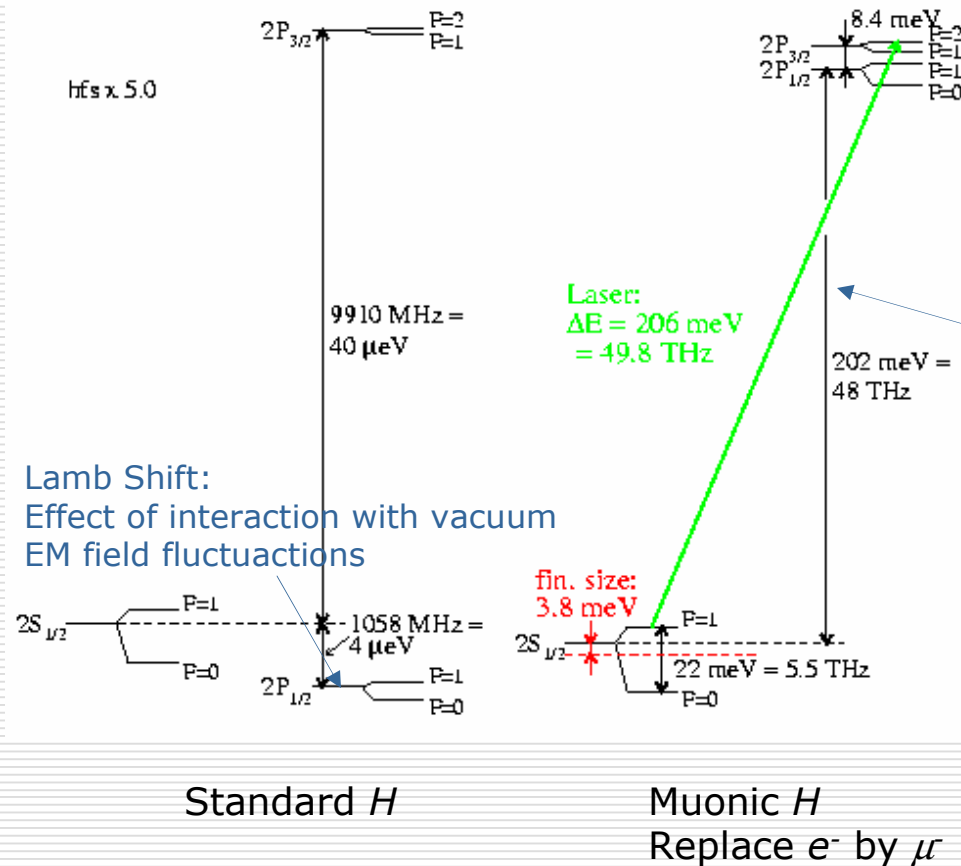
Different:

Reduced mass

Muon magnetic moment

@TBA

Muonic Hydrogen



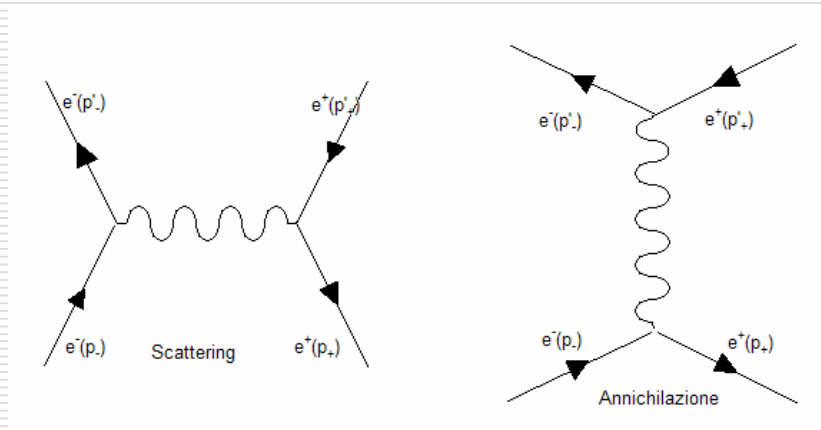
Simplest μ -mesic atom
 Made by stopping μ in
 hydrogenated matter

Huge Lamb shift:
 $\sim 45000 \times \text{Hydrogen!}$
 $\sim (m_\mu/m_e)^2$

@TBA

The $e^+ - e^-$ Effective Interaction

There are now 2 diagrams:

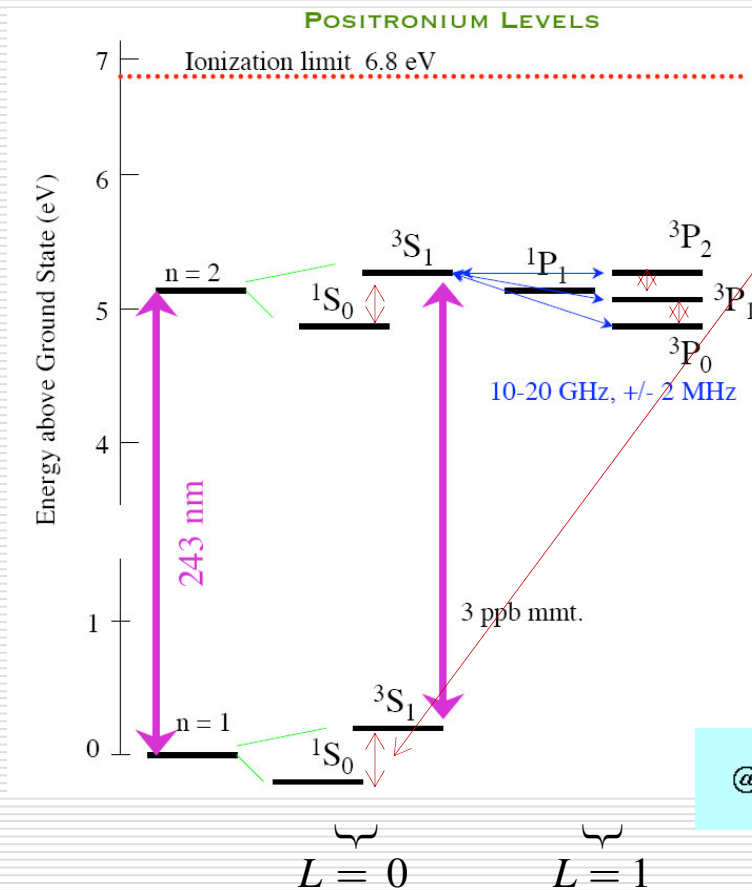


$$T_{fi} = e^2 \left[-\frac{(\bar{u}(p_-)\gamma^\mu u(p_-))(\bar{v}(p_+)\gamma_\mu v(p_+))}{(p_- - p_-')^2} + \frac{(\bar{v}(p_+)\gamma^\mu u(p_-))(\bar{u}(p_-)\gamma_\mu v(p_+))}{(p_+ + p_-)^2} \right]$$

Net effect: Add another term to the effective interaction

$$V_A = \frac{e^2 \pi}{2m^2} (3 + \boldsymbol{\sigma}_1 \cdot \boldsymbol{\sigma}_2) \delta^{(3)}(\mathbf{r}) \quad \text{Same structure as hyperfine term}$$

Positronium



Form of hyperfine term:

$$\Delta E_{hyp} = \left\langle \frac{8\pi e^2}{3m_e^2} \mathbf{s}_{e^-} \cdot \mathbf{s}_{e^+} \delta^{(3)}(\mathbf{r}) \right\rangle \quad \text{for the ground state}$$

More complicated for $n > 1$, $l > 0$

Observe:

Levels labeled by $^S L_J$
S: *Total* spin

Previous pictures:

Levels labeled by $^S L_J$
S: *Electron* spin
 Proton spin only
 in hyperfine term

Light Mesons: Masses

Observe large mass splitting between singlet and triplet mesons:
Guess effective strong interaction has some term similar to hyperfine electromagnetic

Can be shown to be true, to some extent:
When perturbative expansion can be granted, color interaction yields a *chromomagnetic term* with similar hyperfine structure

$$\Delta E = \frac{A}{m_1 m_2} (\mathbf{S}_1 \cdot \mathbf{S}_2)$$

Then expect for the hadron mass:

$$M = m_1 + m_2 + A \frac{(\mathbf{S}_1 \cdot \mathbf{S}_2)}{m_1 m_2}$$

$$\mathbf{J} = \mathbf{S}_1 + \mathbf{S}_2 \rightarrow J^2 = S_1^2 + S_2^2 + 2\mathbf{S}_1 \cdot \mathbf{S}_2$$

$$\rightarrow \mathbf{S}_1 \cdot \mathbf{S}_2 = 1/2 (J^2 - S_1^2 - S_2^2) = 1/2 (J(J+1) - 2S(S+1))$$

$$\mathbf{S}_1 \cdot \mathbf{S}_2 = \begin{cases} +1/4 & \text{triplets} \\ -3/4 & \text{singlets} \end{cases}$$

m_1, m_2 constituent quark mass
Somewhat difficult idea, basically similar to *effective mass* for electrons bound in a crystal.
Will be clarified when discussing QCD.

Masses, including hyperfine splitting

Free parameter counting:

3 quark masses (m_u, m_d, m_s) + 1 constant A

Hope to fit 7 meson masses:

Pseudoscalars + Vectors

→ Go for a 3 constraints fit

Results:

$$m_u = m_d \simeq 310 \text{ MeV}$$

$$m_s \simeq 483 \text{ MeV}$$

$$A \simeq 160 m_{u,d}^2 \text{ MeV}^3$$

Meson	ΔE_{HF}	Fitted mass (MeV)
π	$-\frac{3a}{m_u^2}$	140
K	$-\frac{3a}{m_u m_s}$	485
η	$-\frac{a}{m_u^2} - \frac{2a}{m_s^2}$	559
ρ, ω	$\frac{a}{m_u^2}$	780
K^*	$\frac{a}{m_u m_s}$	896
ϕ	$\frac{a}{m_s^2}$	1032

Light Baryons: Masses - I

Extend the idea to baryons: Sum over 3 quark pairs

$$m(q_1, q_2, q_3) = \sum_{i=1}^3 m_i + A' \frac{1}{2} \sum_{i,j=1}^3 \frac{\mathbf{S}_i \cdot \mathbf{S}_j}{m_i m_j}$$

As an exercise, first neglect differences between quark masses:

$$\mathbf{J} = \mathbf{S}_1 + \mathbf{S}_2 + \mathbf{S}_3 \rightarrow J^2 = (\mathbf{S}_1 + \mathbf{S}_2 + \mathbf{S}_3)^2$$

$$= S_1^2 + S_2^2 + S_3^2 + 2(\mathbf{S}_1 \cdot \mathbf{S}_2 + \mathbf{S}_1 \cdot \mathbf{S}_3 + \mathbf{S}_2 \cdot \mathbf{S}_3)$$

$$S^2 = S(S+1) = 3/4 \rightarrow S_1^2 + S_2^2 + S_3^2 = 9/4$$

$$\rightarrow \mathbf{S}_1 \cdot \mathbf{S}_2 + \mathbf{S}_1 \cdot \mathbf{S}_3 + \mathbf{S}_2 \cdot \mathbf{S}_3 = 1/2 [J^2 - 9/4] = 1/2 J(J+1) - 9/4$$

$$\rightarrow \frac{1}{2} \sum_{i,j=1}^3 \mathbf{S}_i \cdot \mathbf{S}_j = \begin{cases} +3/4 & j = 3/2 \text{ decuplet} \\ -3/4 & j = 1/2 \text{ octet} \end{cases}$$

Light Baryons: Masses - II

Now take into account different quark masses:

$$m(q_1, q_2, q_3) = \sum_{i=1}^3 m_i + A' \frac{1}{2} \sum_{i,j=1}^3 \frac{\mathbf{S}_i \cdot \mathbf{S}_j}{m_i m_j}$$

Must consider quark pairs in order to evaluate hyperfine contribute

$$J_{ik}^2 = (\mathbf{S}_i + \mathbf{S}_k)^2 = S_i^2 + S_k^2 + 2\mathbf{S}_i \cdot \mathbf{S}_k$$

$$\rightarrow \mathbf{S}_i \cdot \mathbf{S}_k = 1/2 [J_{ik}(J_{ik} + 1) - S_i(S_i + 1) - S_k(S_k + 1)]$$

Quarks i, k in a spin triplet state:

$$\mathbf{S}_i \cdot \mathbf{S}_k = 1/2 [1(1+1) - 1/2(1/2+1) - 1/2(1/2+1)]$$

$$\rightarrow \mathbf{S}_i \cdot \mathbf{S}_k = 1/4$$

Quarks i, k in a spin singlet state:

$$\mathbf{S}_i \cdot \mathbf{S}_k = 1/2 [0(0+1) - 1/2(1/2+1) - 1/2(1/2+1)]$$

$$\rightarrow \mathbf{S}_i \cdot \mathbf{S}_k = -3/4$$

Light Baryons: Masses - III

N : Only u, d quarks $\rightarrow$ Same mass

$$\rightarrow m_N = 3m_u - \frac{3A'}{4m_u^2}, \quad m_u = m_d$$

Λ : u, d spin & isospin singlet

$$m_\Lambda = 2m_u + m_s + A' \left(\frac{\mathbf{S}_u \cdot \mathbf{S}_d}{m_u m_d} + \frac{\mathbf{S}_u \cdot \mathbf{S}_s}{m_u m_s} + \frac{\mathbf{S}_s \cdot \mathbf{S}_d}{m_s m_d} \right) = 2m_u + m_s + A' \left(\frac{\mathbf{S}_u \cdot \mathbf{S}_d}{m_u^2} + \frac{\mathbf{S}_u \cdot \mathbf{S}_s + \mathbf{S}_s \cdot \mathbf{S}_d}{m_u m_s} \right)$$

$$\mathbf{S}_u \cdot \mathbf{S}_s + \mathbf{S}_s \cdot \mathbf{S}_d = \frac{1}{2} \sum_{i,j=1}^3 \mathbf{S}_i \cdot \mathbf{S}_j - \mathbf{S}_u \cdot \mathbf{S}_d = -3/4 - (-3/4) = 0 \rightarrow m_\Lambda = 2m_u + m_s - \frac{3A'}{4m_u^2}$$

Σ : u, d spin & isospin triplet

$$m_\Sigma = 2m_u + m_s + A' \left(\frac{\mathbf{S}_u \cdot \mathbf{S}_d}{m_u m_d} + \frac{\mathbf{S}_u \cdot \mathbf{S}_s}{m_u m_s} + \frac{\mathbf{S}_s \cdot \mathbf{S}_d}{m_s m_d} \right) = 2m_u + m_s + A' \left(\frac{\mathbf{S}_u \cdot \mathbf{S}_d}{m_u^2} + \frac{\mathbf{S}_u \cdot \mathbf{S}_s + \mathbf{S}_s \cdot \mathbf{S}_d}{m_u m_s} \right)$$

$$\mathbf{S}_u \cdot \mathbf{S}_s + \mathbf{S}_s \cdot \mathbf{S}_d = \frac{1}{2} \sum_{i,j=1}^3 \mathbf{S}_i \cdot \mathbf{S}_j - \mathbf{S}_u \cdot \mathbf{S}_d = -3/4 - (+1/4) = -1 \rightarrow m_\Sigma = 2m_u + m_s + A' \left(\frac{1}{4m_u^2} - \frac{1}{m_u m_s} \right)$$

Ξ : s_1, s_2 spin triplet ($\leftarrow$ Flavor w.function = ss Symmetric $\rightarrow$ Spin w.function must be symmetric too)

$$m_\Xi = 2m_s + m_u + A' \left(\frac{\mathbf{S}_u \cdot \mathbf{S}_{s1}}{m_u m_s} + \frac{\mathbf{S}_u \cdot \mathbf{S}_{s2}}{m_u m_s} + \frac{\mathbf{S}_{s1} \cdot \mathbf{S}_{s2}}{m_s^2} \right) = 2m_s + m_u + A' \left(\frac{\mathbf{S}_{s1} \cdot \mathbf{S}_{s2}}{m_s^2} + \frac{\mathbf{S}_u \cdot \mathbf{S}_{s1} + \mathbf{S}_u \cdot \mathbf{S}_{s2}}{m_u m_s} \right)$$

$$\mathbf{S}_u \cdot \mathbf{S}_{s1} + \mathbf{S}_u \cdot \mathbf{S}_{s2} = \frac{1}{2} \sum_{i,j=1}^3 \mathbf{S}_i \cdot \mathbf{S}_j - \mathbf{S}_{s1} \cdot \mathbf{S}_{s2} = -3/4 - 1/4 = -1 \rightarrow m_\Xi = 2m_s + m_u + A' \left(\frac{1}{4m_s^2} - \frac{1}{m_u m_s} \right)$$

Light Baryons: Masses - IV

Fit all octet + decuplet:
8 masses $\rightarrow$ 4 constraints

Interesting questions:

Is $A = A'$?

Are the quark masses the same in mesons as in baryons?

$$m_u = m_d \simeq 363 \text{ MeV}$$

$$m_s \simeq 538 \text{ MeV}$$

$$A' \simeq 50 m_{u,d}^2 \text{ MeV}^3$$

Baryons:

Masses $\sim +50 \text{ MeV} \sim 10\%$ higher

Constant $\sim 1/3$ Hyperfine splitting reduced

Baryon	ΔE^{HF}	Fitted mass (MeV)
$N(938)$	$-\frac{3a'}{m_{u,d}^2}$	939
$\Lambda(1116)$	$-\frac{3a'}{m_{u,d}^2}$	1114
$\Sigma(1193)$	$\frac{a'}{m_{u,d}^2} - \frac{4a'}{m_u m_s}$	1179
$\Xi(1318)$	$\frac{a'}{m_{u,d}^2} - \frac{4a'}{m_u m_s}$	1327
$\Delta(1232)$	$+\frac{3a'}{m_{u,d}^2}$	1239
$\Sigma^*(1384)$	$\frac{a'}{m_{u,d}^2} + \frac{4a'}{m_u m_s}$	1381
$\Xi^*(1533)$	$\frac{a'}{m_{u,d}^2} - \frac{4a'}{m_u m_s}$	1529
$\Omega(1672)$	$+\frac{3a'}{m_s^2}$	1682

Light Baryons: Magnetic Moments - I

Take total dipole moment operator as the sum of the single quark operators:

$$\boldsymbol{\mu} = \sum_{i=1}^3 \boldsymbol{\mu}_i \rightarrow \mu_p = \langle p, +1/2 | \boldsymbol{\mu} | p, +1/2 \rangle = \langle p, +1/2 | (\boldsymbol{\mu}_1 + \boldsymbol{\mu}_2 + \boldsymbol{\mu}_3) | p, +1/2 \rangle$$

Can this be really granted??

$$|p, +1/2\rangle = \frac{1}{\sqrt{18}} \begin{pmatrix} 2u \uparrow u \uparrow d \downarrow + 2d \downarrow u \uparrow u \uparrow + 2u \uparrow d \downarrow u \uparrow - \\ -u \downarrow d \uparrow u \uparrow - u \uparrow u \downarrow d \uparrow - u \downarrow u \uparrow d \uparrow - \\ -d \uparrow u \downarrow u \uparrow - u \uparrow d \uparrow u \downarrow - d \uparrow u \uparrow u \downarrow \end{pmatrix}$$

Each operator acting on the corresponding factor of the wave function

Now dive in some really dummy algebra:

$$4\langle u \uparrow u \uparrow d \downarrow | (\boldsymbol{\mu}_1 + \boldsymbol{\mu}_2 + \boldsymbol{\mu}_3) | u \uparrow u \uparrow d \downarrow \rangle = 4[\langle u | \boldsymbol{\mu}_1 | u \rangle + \langle u | \boldsymbol{\mu}_2 | u \rangle - \langle d | \boldsymbol{\mu}_3 | d \rangle] = 4[\mu_u + \mu_u - \mu_d] = 8\mu_u - 4\mu_d$$

$$4\langle d \downarrow u \uparrow u \uparrow | (\boldsymbol{\mu}_1 + \boldsymbol{\mu}_2 + \boldsymbol{\mu}_3) | d \downarrow u \uparrow u \uparrow \rangle = 4\langle u \uparrow d \downarrow u \uparrow | (\boldsymbol{\mu}_1 + \boldsymbol{\mu}_2 + \boldsymbol{\mu}_3) | u \uparrow d \downarrow u \uparrow \rangle = 8\mu_u - 4\mu_d$$

$$\langle u \downarrow d \uparrow u \uparrow | (\boldsymbol{\mu}_1 + \boldsymbol{\mu}_2 + \boldsymbol{\mu}_3) | u \downarrow d \uparrow u \uparrow \rangle = \langle u \uparrow u \downarrow d \uparrow | (\boldsymbol{\mu}_1 + \boldsymbol{\mu}_2 + \boldsymbol{\mu}_3) | u \uparrow u \downarrow d \uparrow \rangle$$

$$\langle d \uparrow u \downarrow u \uparrow | (\boldsymbol{\mu}_1 + \boldsymbol{\mu}_2 + \boldsymbol{\mu}_3) | d \uparrow u \downarrow u \uparrow \rangle = \dots = \mu_d$$

$$\rightarrow \langle p, +1/2 | (\boldsymbol{\mu}_1 + \boldsymbol{\mu}_2 + \boldsymbol{\mu}_3) | p, +1/2 \rangle = \frac{1}{18} [3(8\mu_u - 4\mu_d) + 6\mu_d] = \frac{1}{18} [24\mu_u - 6\mu_d] = \frac{1}{3} (4\mu_u - \mu_d) \equiv \mu_p$$

Then take neutron: Just swap $u \leftrightarrow d$

$$|n, +1/2\rangle = \frac{1}{\sqrt{18}} \begin{pmatrix} 2d \uparrow d \uparrow u \downarrow + 2u \downarrow d \uparrow d \uparrow + 2d \uparrow u \downarrow d \uparrow - \\ -d \downarrow u \uparrow d \uparrow - d \uparrow d \downarrow u \uparrow - d \downarrow d \uparrow u \uparrow - \\ -u \uparrow d \downarrow d \uparrow - d \uparrow u \uparrow d \downarrow - u \uparrow d \uparrow d \downarrow \end{pmatrix} \rightarrow \mu_n = \frac{1}{3} (4\mu_d - \mu_u)$$

Light Baryons: Magnetic Moments - II

Take quarks as Dirac particles:

$$\mu = \frac{e}{2m}$$

Can this be really granted??

$$\rightarrow \mu_u = \frac{2}{3} \frac{e}{2m_u}, \mu_d = -\frac{1}{3} \frac{e}{2m_d}, \mu_s = -\frac{1}{3} \frac{e}{2m_s}$$

Predict moment ratio:

$$\frac{\mu_n}{\mu_p} = \frac{4\mu_d - \mu_u}{4\mu_u - \mu_d} = \frac{4\left(-\frac{e}{3 \cdot 2m_d}\right) - \frac{2e}{3 \cdot 2m_u}}{4\frac{2e}{3 \cdot 2m_u} - \left(-\frac{e}{3 \cdot 2m_d}\right)} \simeq \frac{-4e - 2e}{8e + e} = -\frac{2}{3} \approx -0.667$$

Experimental ratio:

$$\frac{\mu_n}{\mu_p} \approx -0.685 \quad \text{Amazingly close!}$$

Absolute moments difficult to estimate, as involving unknown quark masses.

Nevertheless..

Magnetic Moments - Octet

...If one insists in believing the constituent quark masses have something to do with reality, can compute the expected magnetic moments for octet:

Not too bad for such a simple attempt...

We are taking baryons as composed only by valence quarks, which is wildly *incomplete*

<i>Baryon</i>	<i>Moment</i>	<i>Predicted</i>	<i>Observed</i>
p	$\frac{4}{3}\mu_u - \frac{1}{3}\mu_d$	2.79	2.793
n	$\frac{4}{3}\mu_d - \frac{1}{3}\mu_u$	-1.86	-1.913
Λ^0	μ_s	-0.58	-0.614
Σ^+	$\frac{4}{3}\mu_u - \frac{1}{3}\mu_s$	2.68	2.33
Σ^0	$\frac{2}{3}(\mu_u + \mu_d) - \frac{1}{3}\mu_s$	0.82	Unstable
Σ^-	$\frac{4}{3}\mu_d - \frac{1}{3}\mu_s$	-1.05	-1.41
Ξ^0	$\frac{4}{3}\mu_s - \frac{1}{3}\mu_u$	-1.40	-1.253
Ξ^-	$\frac{4}{3}\mu_s - \frac{1}{3}\mu_d$	-0.47	-0.69

Hyperon Magnetic Moments

Dipole rotation along a variable path length in a uniform magnetic field

$$\phi_{rot} \propto \mu_{\Lambda^0} \int B dl$$

Collect Λ^0 decays at different distances
Measure Λ^0 energy

Λ^0 produced polarized at high energy

$\mu_{\Lambda} \parallel s_{\Lambda}$: $s_{\Lambda} \perp \Lambda^0$ production plane

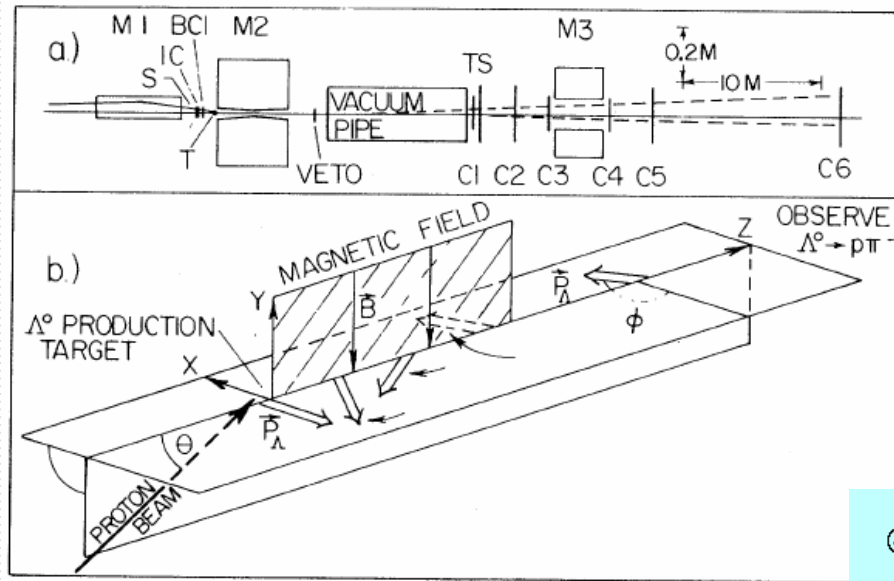
$\rightarrow s_{\Lambda}$ initial known

$\Lambda^0 \rightarrow p + \pi^0$ weak decay

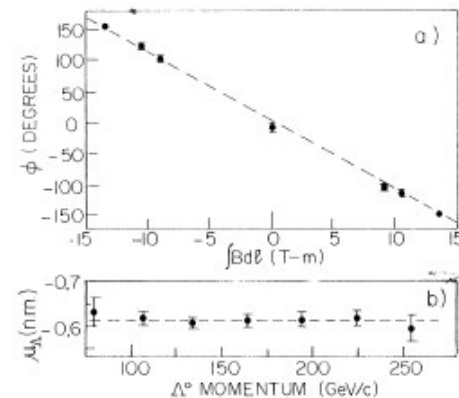
Parity violation

$\rightarrow p$ momentum preferentially $\angle s_{\Lambda}$

$\rightarrow s_{\Lambda}$ final known from $\mathbf{p}_p$ direction



@TBA



Mag. Moments of Unstable Particles

Cannot measure them with the same technique discussed for stable particles
 → Consider a different approach

Take as an example the Σ^0 decay (octet):

$\Sigma^0 \rightarrow \Lambda^0 + \gamma$ Parity conserving, electromagnetic decay

$\eta_P(\Sigma^0) = + = \eta_P(\Lambda^0) \rightarrow \eta_P(\gamma)$ must be +

$$\eta_P(\gamma) = \begin{cases} (-1)^{j+1} & \text{magnetic} \\ (-1)^j & \text{electric} \end{cases}, j=1 \rightarrow \text{magnetic} \leftarrow$$

Just meaning:
 This photon has
total angular momentum = 1
total parity = +1

Transition is *M1* (magnetic dipole)

Σ^0, Λ^0 : Same quark content *uds*

Σ^0 : I-spin triplet $\rightarrow u, d$ Spin triplet

→ Wave function = Sum of Permutations of $[(ud + du)s] \uparrow\uparrow\downarrow$

Λ^0 : I-spin singlet $\rightarrow u, d$ Spin singlet

→ Wave function = Sum of Permutations of $[(ud - du)s](\uparrow\downarrow - \downarrow\uparrow)\uparrow$

Amplitude = *A(Spin flip)* for $[u \text{ or } d]$

The Transition Magnetic Moment

Reminder: The neutron e.m. transition 4-current

$$j_n^\mu = e \bar{u}_n(p') \left(F_n(q^2) \gamma^\mu + G_n(q^2) i \kappa_n \frac{\sigma^{\mu\nu}}{2m} q_\nu \right) u_n(p)$$

Take the Σ^0 as a kind of neutron... (Well, that's $SU(3)$...)

This would be the current involved e.g. in electron DIS off a Σ^0

$$j_{\Sigma^0}^\mu = e \bar{u}_{\Sigma^0}(p') \left(F_{\Sigma^0}(q^2) \gamma^\mu + G_{\Sigma^0}(q^2) i \kappa_{\Sigma^0} \frac{\sigma^{\mu\nu}}{2m} q_\nu \right) u_{\Sigma^0}(p)$$

$\kappa_{\Sigma^0} = ?$ Not observable, the Σ^0 is unstable

Define an e.m. transition current for our process

$$j_{\Sigma^0 \Lambda^0}^\mu = e \bar{u}_{\Lambda^0}(p') \left(F_{\Sigma^0 \Lambda^0}(q^2) \gamma^\mu + G_{\Sigma^0 \Lambda^0}(q^2) i \kappa_{\Sigma^0 \Lambda^0} \frac{\sigma^{\mu\nu}}{2m} q_\nu \right) u_{\Sigma^0}(p)$$

$\kappa_{\Sigma^0 \Lambda^0} = ?$ Can be determined by the observed rate

In the static ($q^2=0$) limit (actually never reached in the transition) analog to the static magnetic dipole moment

Vector Mesons Radiative Decays

Take radiative decays of vector mesons to pseudoscalars:

$$V \rightarrow P + \gamma$$

$$1^{--} \rightarrow 0^{-+} + \gamma$$

$\rightarrow \gamma: 1^+ \rightarrow \text{magnetic dipole}$

For any magnetic dipole transition:

$$\text{Rate} \propto \omega^3, \omega: \text{Photon energy}$$

From quark model perspective: *Triplet* $\rightarrow$ *Singlet*, *S-wave*

As before: Spin flip of one quark

(I = Fraction of space overlap between initial/final meson wave functions)

Process	Rate in units of $\omega^3 I ^2 / 3\pi$	Prediction (keV)	Experiment (keV)	$ I ^2$
$\omega \rightarrow \pi^0 \gamma$	$(\mu_u - \mu_d)^2$	$1390 I ^2$	890 ± 50	0.64 ± 0.04
$\rho \rightarrow \pi \gamma$	$(\mu_u + \mu_d)^2$	$148 I ^2$	67 ± 7	0.45 ± 0.05
$\omega \rightarrow \eta \gamma$	$(\mu_u + \mu_d)^2 / 2$	$11 I ^2$	$3^{+2.5}_{-1.8}$	$0.27^{+0.23}_{-0.16}$
$\rho \rightarrow \eta \gamma$	$(\mu_u - \mu_d)^2 / 2$	$92 I ^2$	50 ± 13	0.54 ± 0.14
$\eta' \rightarrow \omega \gamma$	$3(\mu_u + \mu_d)^2 / 2$	$17 I ^2$	7.6 ± 3	0.45 ± 0.18
$\eta' \rightarrow \rho \gamma$	$3(\mu_u - \mu_d)^2 / 2$	$171 I ^2$	83 ± 30	0.48 ± 0.18
$\phi \rightarrow \eta \gamma$	$2\mu_s^2$	$110 I ^2$	62 ± 9	0.56 ± 0.08
$\phi \rightarrow \pi^0 \gamma$	0	0	5.7 ± 2	
$K^{*+} \rightarrow K^+ \gamma$	$(\mu_u + \mu_s)^2$	$153 I ^2$	60 ± 15	0.39 ± 0.10
$K^{*0} \rightarrow K^0 \gamma$	$(\mu_d - \mu_s)^2$	$224 I ^2$	75 ± 35	0.34 ± 0.16

Quite consistent with simple
SU(3) symmetry: Same overlap

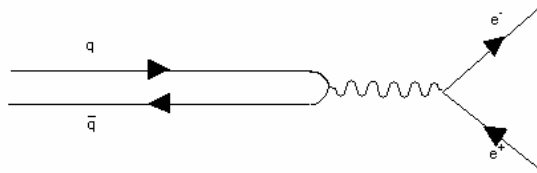
@TBA

Decays of Vector Mesons to $e^+ e^-$

$$\rho^0 \rightarrow e^+ + e^-$$

$$\omega \rightarrow e^+ + e^-$$

$$\varphi \rightarrow e^+ + e^-$$



$$\Gamma_{e^+e^-} = \frac{16\pi\alpha^2}{q^2 (=M_V^2)} |\psi(0)|^2 \left| \sum_i a_i Q_i \right|^2$$

Van Royen – Weisskopf Formula
Partial decay rate

$$\rho^0: \frac{1}{\sqrt{2}}(u\bar{u} - d\bar{d}) \rightarrow \left| \sum_i a_i Q_i \right|^2 = \left| \frac{1}{\sqrt{2}} \left(\frac{2}{3} - \left(-\frac{1}{3} \right) \right) \right|^2 = \frac{1}{2}$$

$$\omega: \frac{1}{\sqrt{2}}(u\bar{u} + d\bar{d}) \rightarrow \left| \sum_i a_i Q_i \right|^2 = \left| \frac{1}{\sqrt{2}} \left(\frac{2}{3} + \left(-\frac{1}{3} \right) \right) \right|^2 = \frac{1}{18}$$

$$\varphi: s\bar{s} \rightarrow \left| \sum_i a_i Q_i \right|^2 = \left| \left(-\frac{1}{3} \right) \right|^2 = \frac{1}{9}$$

$$\rightarrow \Gamma_{e^+e^-}(\rho^0) : \Gamma_{e^+e^-}(\omega) : \Gamma_{e^+e^-}(\varphi) = 9:1:2$$

Van-Royen - Weisskopf - I

Attempting to calculate the vector meson decay rate:

$\Gamma_V = |A_V|^2$, $A_V = \langle f | T | V \rangle$ Transition amplitude between V (initial), f (final) state

The meson is a bound state $\rightarrow$ Initial state *not* a plane wave!

Then expand the amplitude into plane waves:

$$A_V = \sum_p \underbrace{\langle f | T | p \rangle}_{A(\mathbf{p})} \underbrace{\langle p | V \rangle}_{\psi(\mathbf{p})}$$

$A(\mathbf{p})$ plane wave amplitude, $\psi(\mathbf{p})$ momentum space wave function

$$\rightarrow A_V = \int d^3\mathbf{p} A(\mathbf{p}) \psi(\mathbf{p}) = \frac{1}{(2\pi)^{3/2}} \int d^3\mathbf{p} A(\mathbf{p}) \int \psi(\mathbf{r}) e^{i\mathbf{p}\cdot\mathbf{r}} d^3\mathbf{r} = \frac{1}{(2\pi)^{3/2}} \int d^3\mathbf{r} \psi(\mathbf{r}) \int A(\mathbf{p}) e^{i\mathbf{p}\cdot\mathbf{r}} d^3\mathbf{p}$$

$$\text{Take } A(\mathbf{p}) \approx A = \text{const} \rightarrow A_V \approx \frac{A}{(2\pi)^{3/2}} \int d^3\mathbf{r} \psi(\mathbf{r}) \underbrace{\int e^{i\mathbf{p}\cdot\mathbf{r}} d^3\mathbf{p}}_{(2\pi)^3 \delta^3(\mathbf{r})} = (2\pi)^{3/2} A \psi(0)$$

$$\rightarrow \Gamma_V = |A_V|^2 \approx (2\pi)^3 |A|^2 |\psi(0)|^2$$

Why is $A(p) \approx \text{const}$?

Van-Royen - Weisskopf - II

Consider the process involving free quarks (plane wave):

$$q\bar{q} \rightarrow e^+e^-$$

Cross section, averaged over initial, summed over final spin projections:

$$\sigma_{q\bar{q} \rightarrow e^+e^-}(p) = \underbrace{|A(p)|^2}_{\text{flux}} \frac{1}{v}, v \text{ } q, \bar{q} \text{ relative velocity} \rightarrow \sigma_{q\bar{q} \rightarrow e^+e^-}(p) = |A(p)|^2 \frac{(2\pi)^3}{v}$$

Take 1 photon, annihilation, QED diagram:

$$\sigma_{q\bar{q} \rightarrow e^+e^-}(p) = \frac{\pi\alpha^2 Q^2}{s} \frac{p_e}{p_q} \left(1 + \frac{v^2}{3} + \frac{4m_q^2}{s} \right)$$

For small initial velocity:

$$s \approx (2m_q)^2, p_q \approx m_q \frac{v}{2}$$

$$\sigma_{q\bar{q} \rightarrow e^+e^-}(p) = \frac{\pi\alpha^2 Q^2}{s} \frac{p_e}{p_q} \left(1 + \frac{v^2}{3} + \frac{4m_q^2}{s} \right) \approx \frac{\pi\alpha^2 Q^2}{4m_q^2} \frac{p_e}{m_q \frac{v}{2}} \left(1 + \frac{v^2}{3} + 1 \right) \approx \frac{\pi\alpha^2 Q^2}{4m_q^2} \frac{p_e}{m_q v} 4 \approx \frac{\pi\alpha^2 Q^2}{m_q^3} \frac{p_e}{v}$$

Just the same as

$$e^+ + e^- \rightarrow \mu^+ + \mu^-$$

But:

Do not neglect rest mass!

Clumsy algebra..

Van-Royen - Weisskopf - III

Obtain the decay rate:

$$\rightarrow \sigma_{q\bar{q} \rightarrow e^+e^-}(p) = |A(p)|^2 \frac{(2\pi)^3}{v} = \frac{\pi\alpha^2 Q^2}{m_q^3} \frac{p_e}{v}$$

$$\rightarrow |A(p)|^2 = \frac{\pi\alpha^2 Q^2}{m_q^3} \frac{p_e}{v} \frac{v}{(2\pi)^3} = \frac{\pi\alpha^2 Q^2}{m_q^3} \frac{p_e}{(2\pi)^3}$$

$$p_e \approx m_q \rightarrow |A(p)|^2 \approx \frac{\pi\alpha^2 Q^2}{(2\pi)^3 m_q^2}$$

Neglect quark momentum, electron mass

$$m_q \approx \frac{M_V}{2} \rightarrow \Gamma_V \approx \frac{4\pi\alpha^2 Q^2}{(2\pi)^3 M_V^2} \quad p\text{-independent OK}$$

Can this really be granted?

We took the average over initial spins for our plane wave cross section, resulting in 3 (triplet) + 1 (singlet) = 4 states

Vector mesons have spin 1, so we should not count spin 0

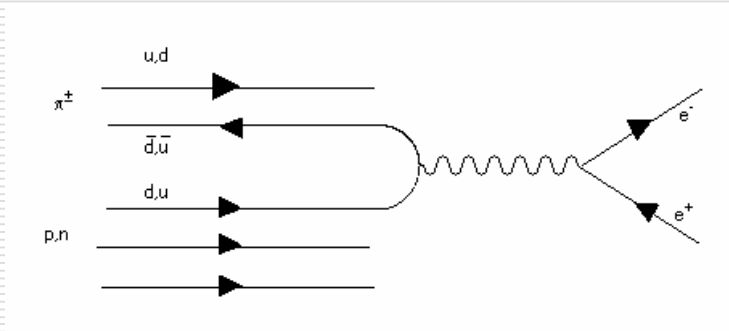
→ Get a further factor 4/3:

$$\Gamma_V \approx \frac{4}{3} \frac{4\pi\alpha^2 Q^2}{(2\pi)^3 M_V^2} = \frac{16}{3} \frac{\pi\alpha^2 Q^2}{(2\pi)^3 M_V^2}$$

This formula is still incomplete...
Missing factor 3 ??

Drell-Yan from Isoscalar Targets

Take production of electron pairs from pion beams: *Drell-Yan*



$$\text{Rate} \propto Q_{quark}^2$$

Cross section: Electromagnetic, counting antiquark content in π
 For isoscalar targets: $N_p = N_n \rightarrow N_u = N_d$

$$\left. \begin{aligned} \sigma(\pi^+) &\propto Q_d^2 = \frac{1}{9} \\ \sigma(\pi^-) &\propto Q_u^2 = \frac{4}{9} \end{aligned} \right\} \rightarrow \frac{\sigma(\pi^-)}{\sigma(\pi^+)} = 4$$

More Quarks

<i>Flavor</i>	<i>Mass</i>	Q	I	I_3	S	C	B	T
Up	5.6 MeV	2/3	$\frac{1}{2}$	+1/2	0	0	0	0
Down	9.9 MeV	-1/3	$\frac{1}{2}$	-1/2	0	0	0	0
strange	199 MeV	-1/3	0	0	1	0	0	0
charm	1350 MeV	2/3	0	0	0	1	0	0
bottom	4400 MeV	-1/3	0	0	0	0	1	0
Top	174 GeV	2/3	0	0	0	0	0	1

Charm

Predicted to exist in order to solve a major puzzle in weak interactions via the *GIM Mechanism* (see later)

Expected mass much higher than u, d, s

Phenomenology similar to strange quark s :

New breed of *charmed particles*, both mesons and baryons

Difference: Much larger mass

→ Many channels open to weak decays → *Shorter lifetime* $\sim 10^{-13}$ s

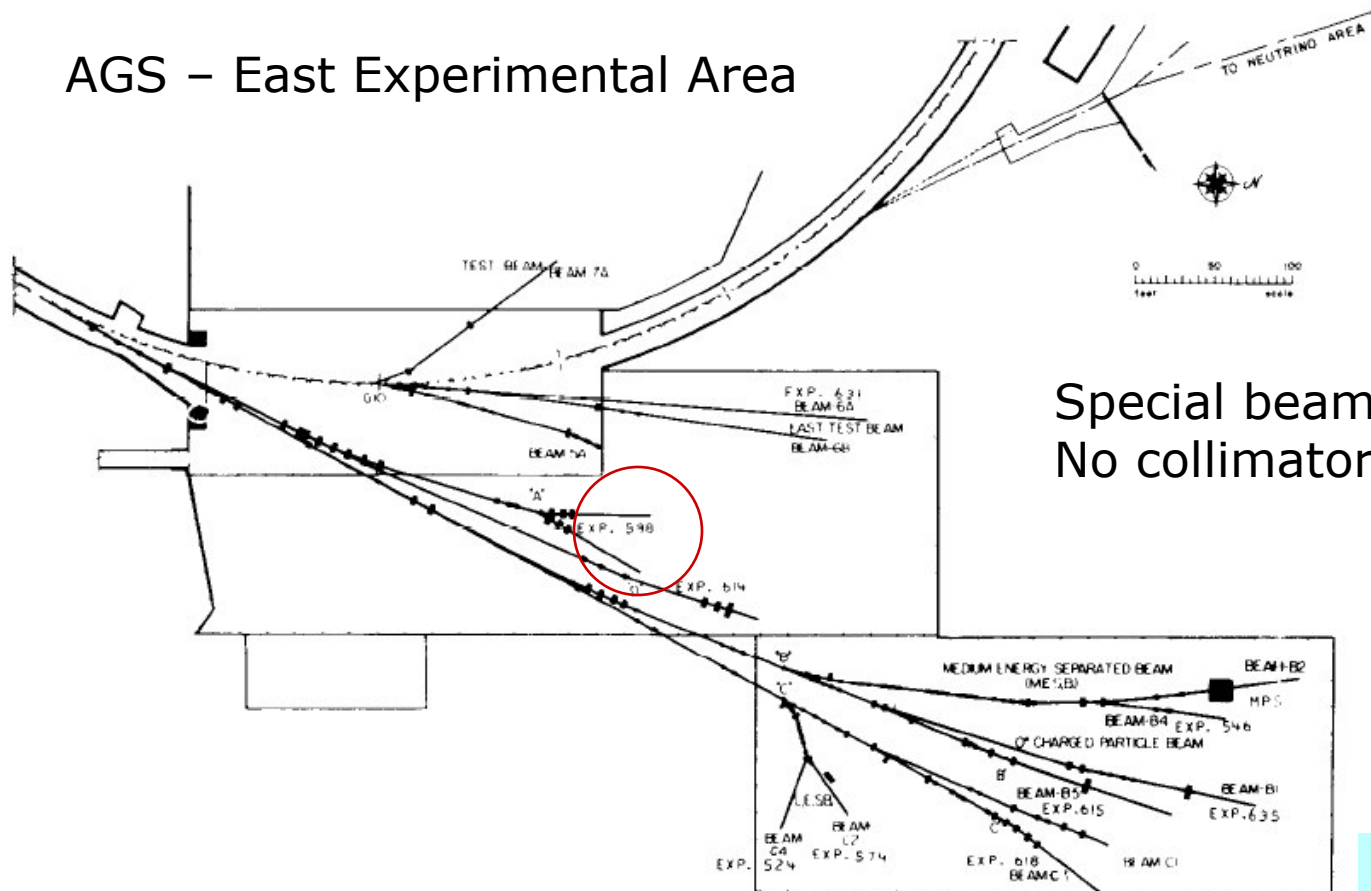
→ Extended symmetry severely broken → *SU(4) not useful*

Based on *asymptotic freedom* of QCD (see later), guess an exciting physics case for a *heavy, hidden charm bound state*

Discovered simultaneously at SLAC (Mark I) and BNL (E598)

The J/ψ Particle at Brookhaven - I

AGS - East Experimental Area



Special beam (Line A)
No collimators

@TBA

The J/ψ Particle at Brookhaven - II

2-Arm Spectrometer

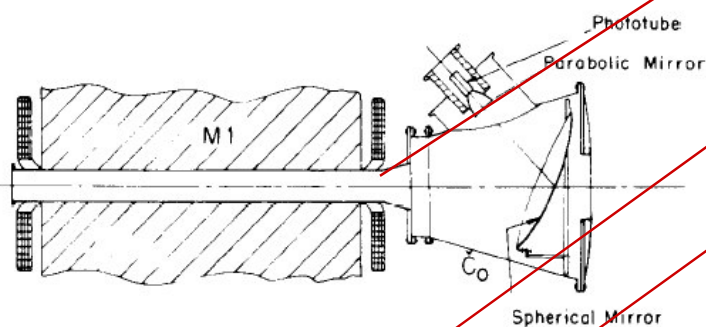
Magnetic analysis (3 dipoles)

Excellent electron identification (2 Cerenkov)

Very high intensity ($2 \cdot 10^{12}$ ppp)

Small spot size ($3 \times 6 \text{ mm}^2$)

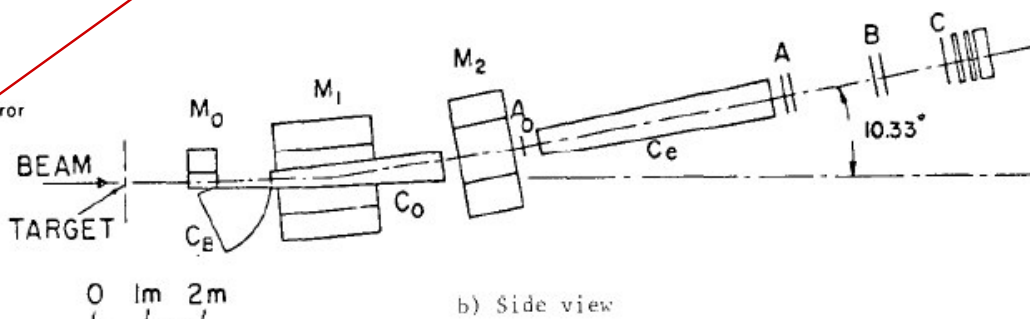
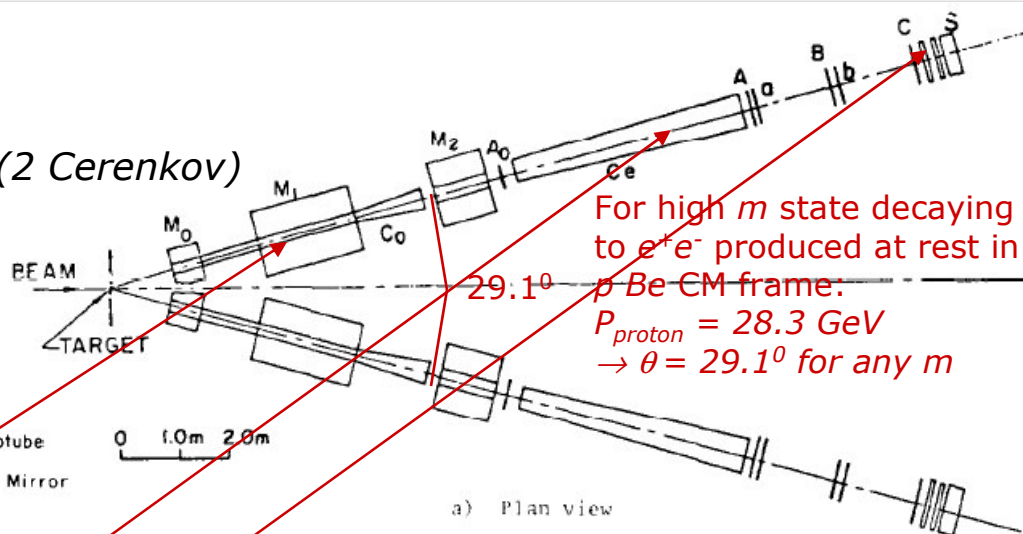
Electron identification



1st Cerenkov
(inside magnet)

2nd Cerenkov

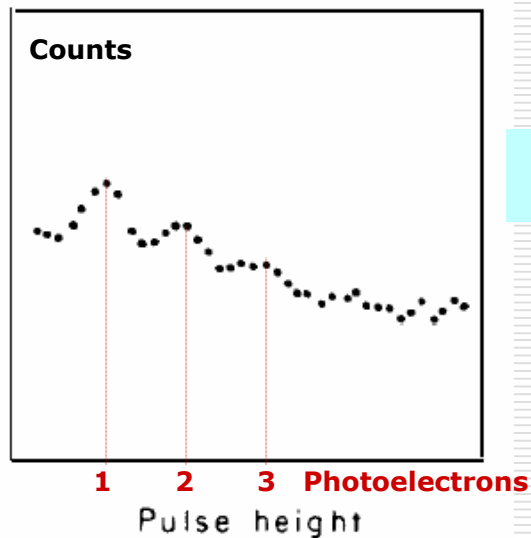
Calorimeter



The J/ψ Particle at Brookhaven - III

Detectors in the experimental setup at BNL

Cherenkov photomultiplier spectrum
Excellent single electron resolution

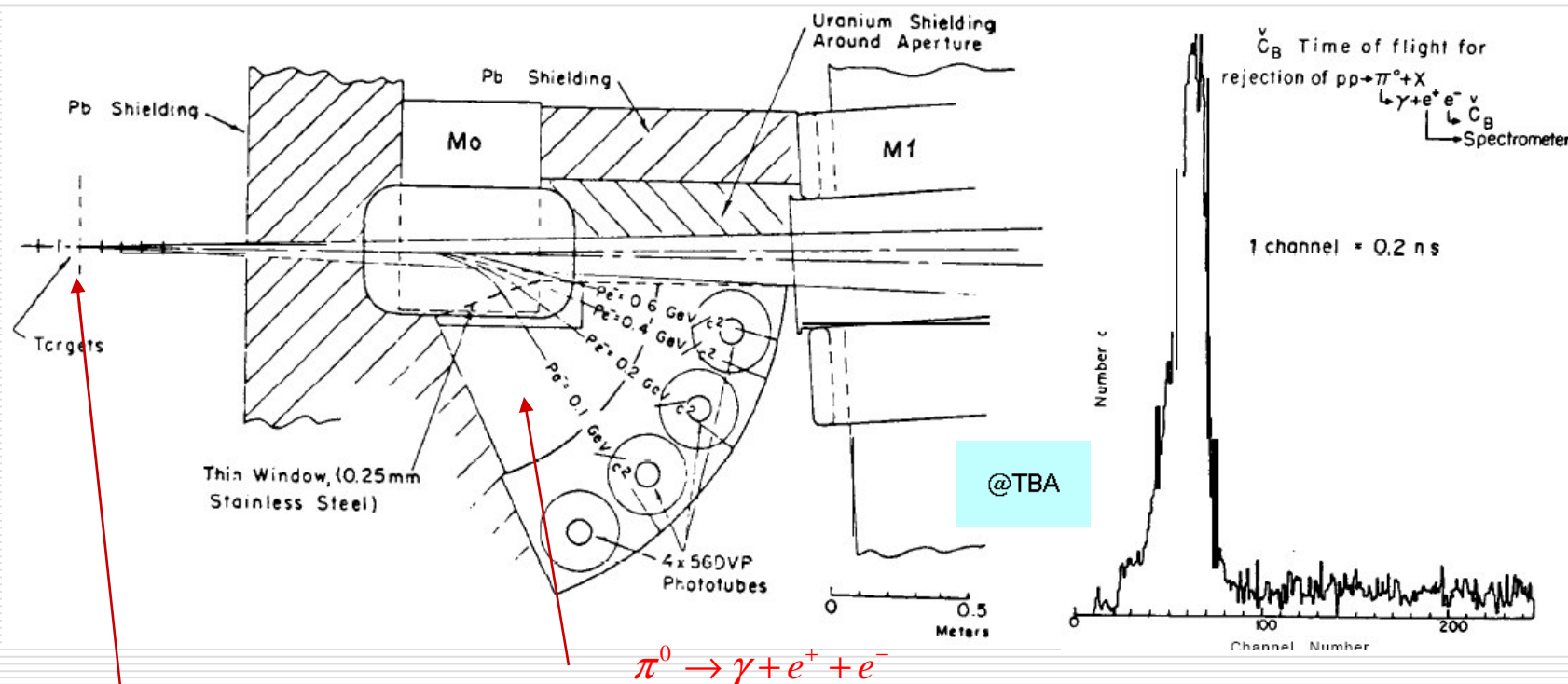


Good for signal efficiency, background rejection



The J/ψ Particle at Brookhaven - IV

Target area: Special Magnet (M_0) + Cerenkov (C_B) for calibration
 Want to be sure not to trigger on e^+e^- pairs from γ conversion



Segmented target
 Improved mass resolution

Beam Cerenkov
 Detects electrons from Dalitz decays
 Logic: $C_B * (C_0 + C_e) * Cal$

Time resolution $\sim 1 \text{ ns}$

The J/ψ Particle at Brookhaven - V

Large peak observed at $m = 3.1 \text{ GeV}$

Still present at the same mass after reducing magnet currents by 10%

@TBA

Very large mass

Wrong by $6/3100 \sim 2 \cdot 10^{-3}$

Excellent control of systematics!

Very small width

Consistent with experimental resolution

Totally inconsistent with standard picture of a high mass hadron:
In view of the many decay channels which are open, should be quite wide

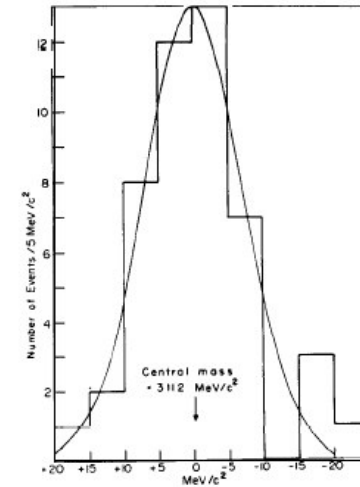
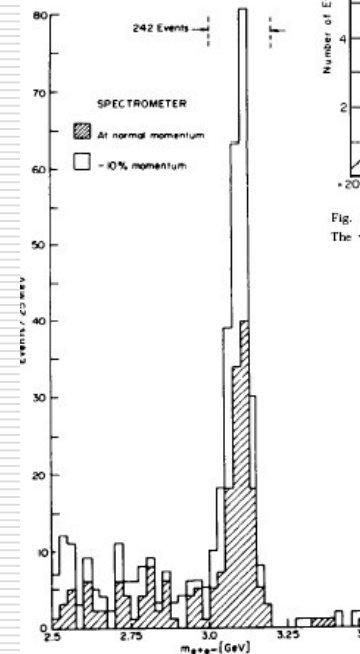
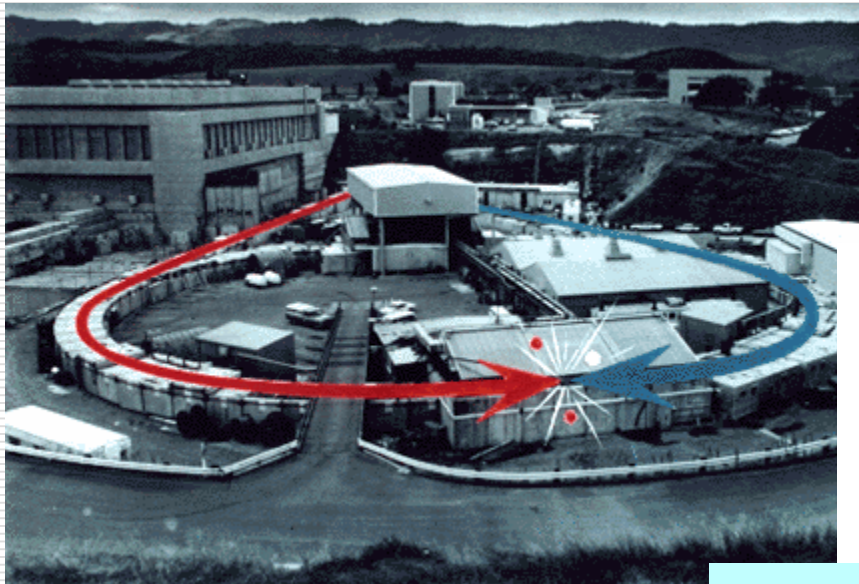


Fig. 12b. The measurement of the width of the J/ψ . The width is shown to be less than 5 MeV.

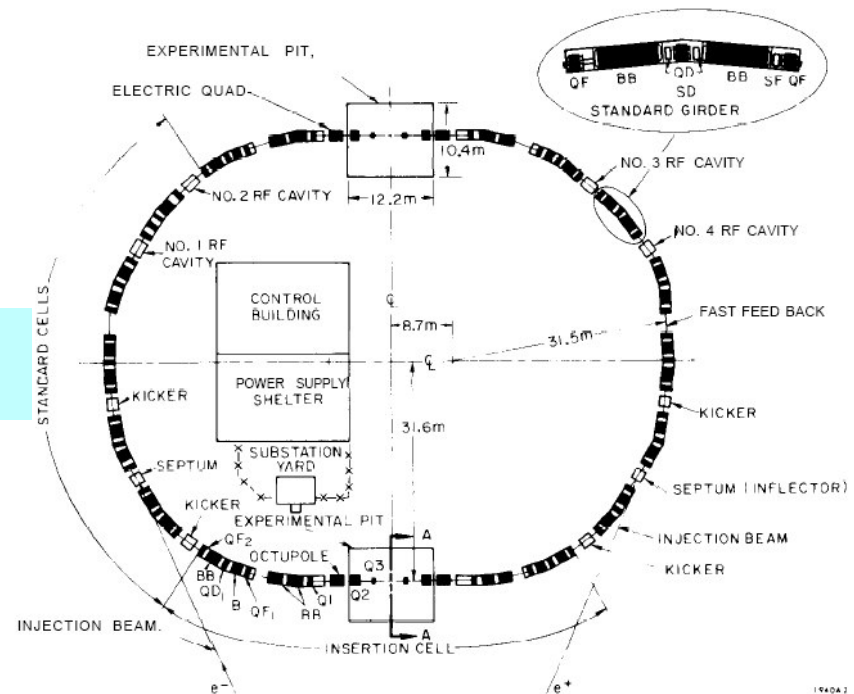
$$\sigma_m = 5 \text{ MeV}$$

The J/ψ Particle at SLAC - I

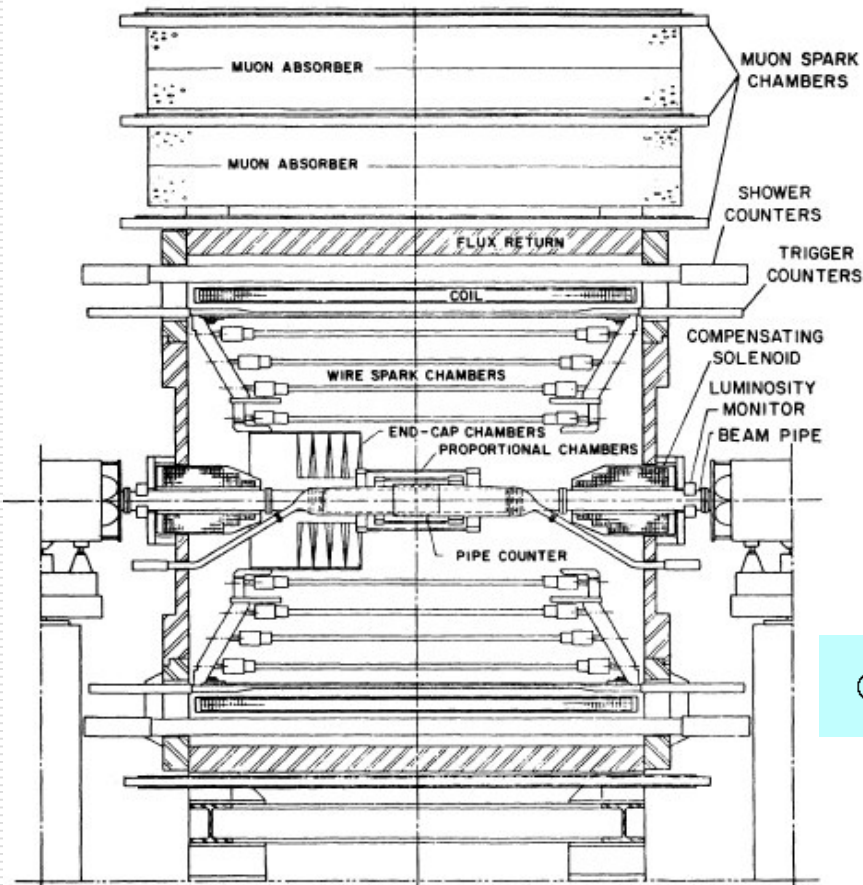


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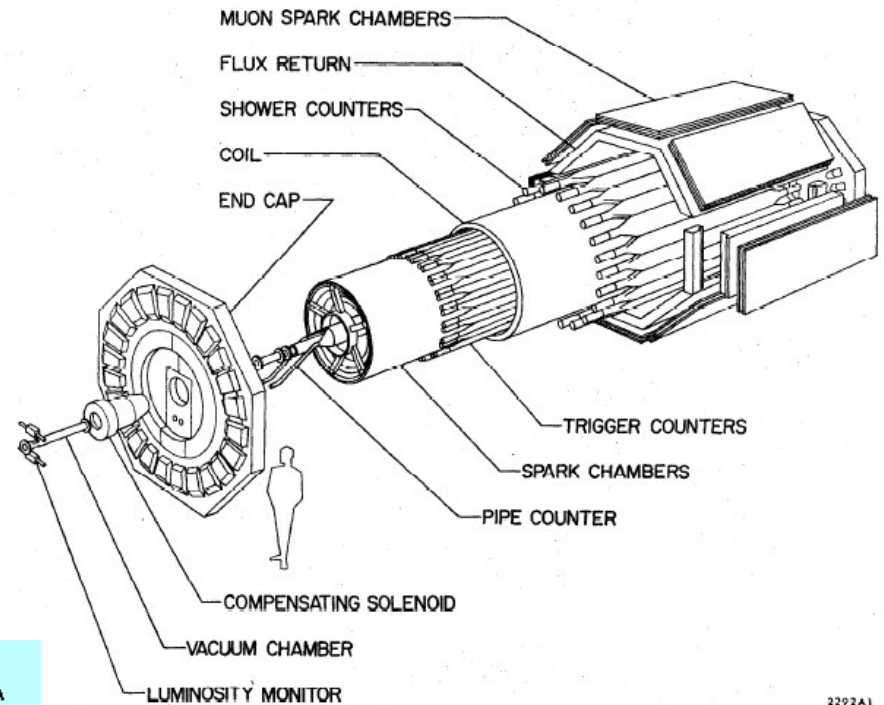
SPEAR



The J/ψ Particle at SLAC - II



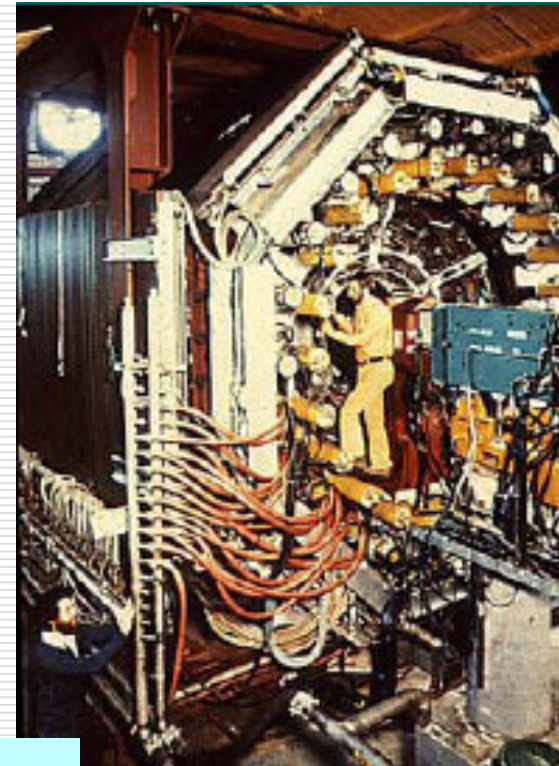
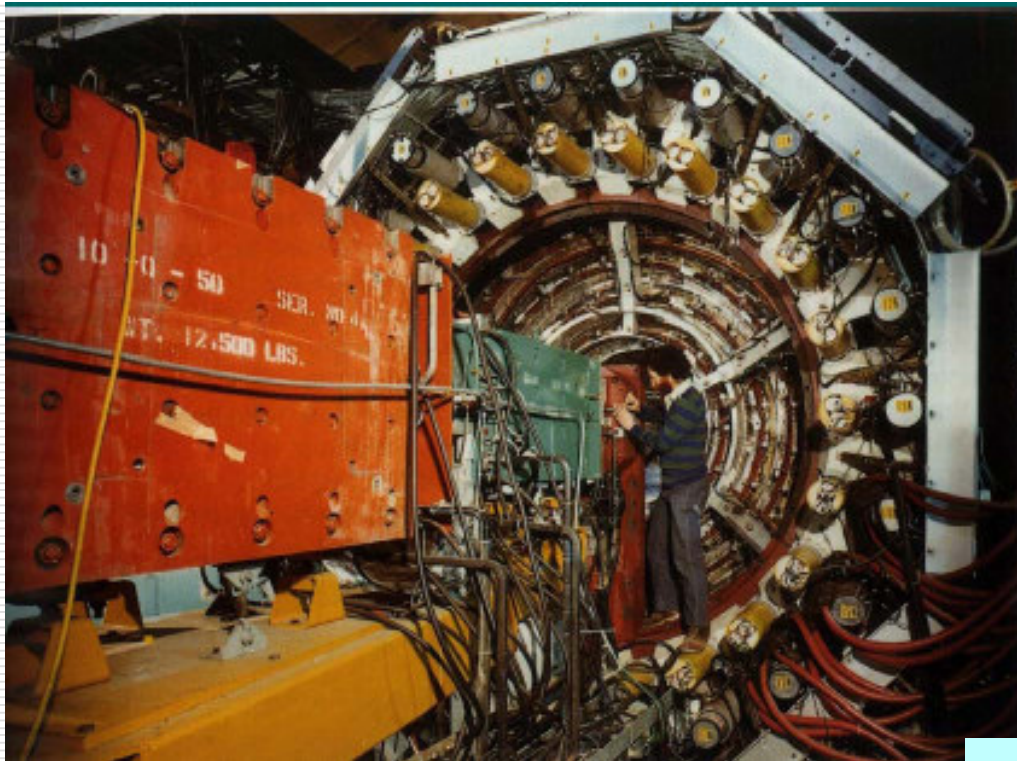
@TBA



The Mark I Detector

The J/ψ Particle at SLAC - III

Mark I: First example of multi-purpose, collider detector



@TBA

The J/ψ Particle at SLAC - IV

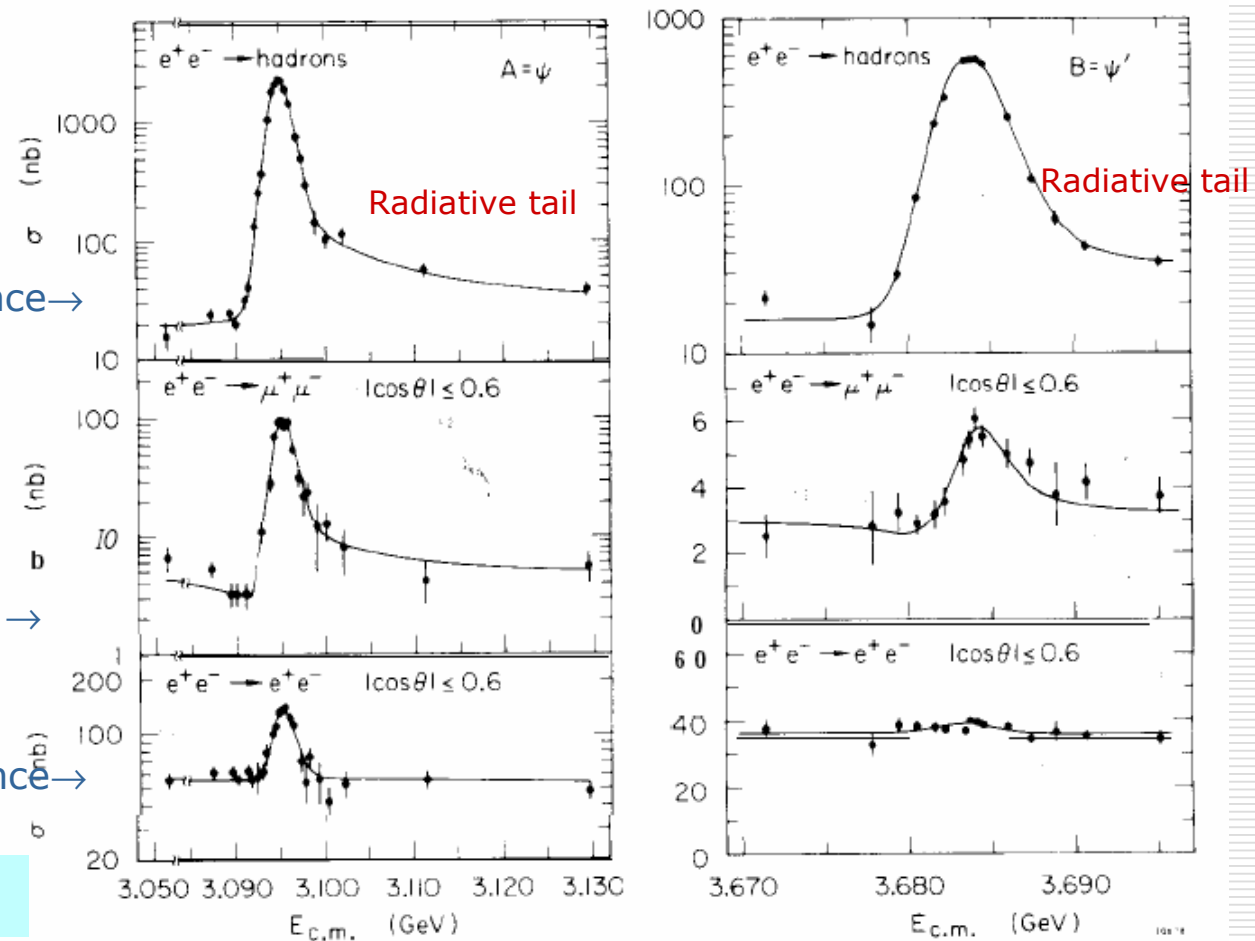
J/ψ and ψ' as seen
in different channels

No interference →

Interference →

No interference →

@TBA



What is the J/ψ ?

Quickly understood as the first, indirect evidence for charm
Bound state of $c, \bar{c}$ quark-antiquark pair

Another member of the vector mesons family

Main differences:

Charm quark has a large mass 1.5 GeV

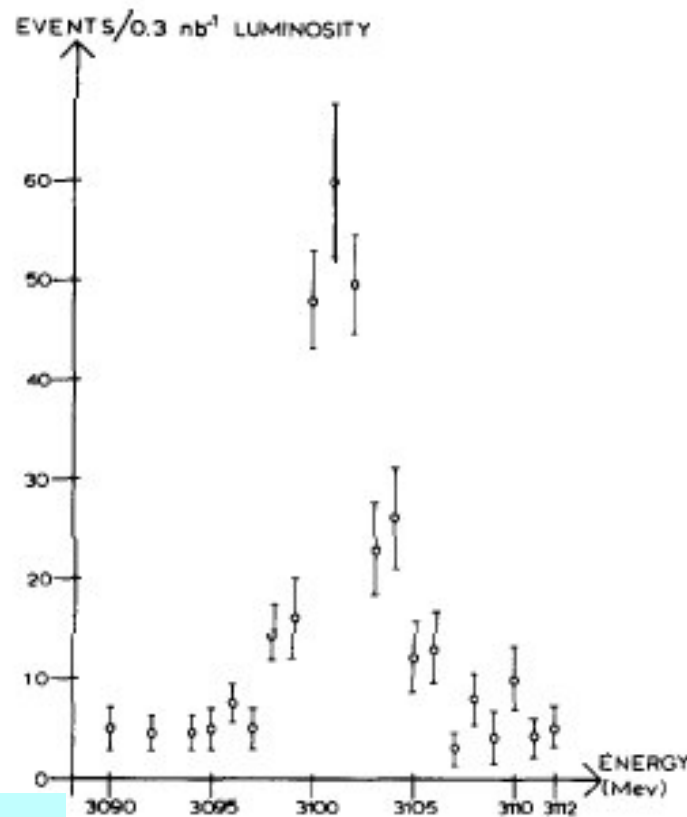
Lightest charmed particles are so heavy the J/ψ cannot decay into a pair of them → Most decays channels are closed

The J/ψ Particle at Frascati

One day of November, back in 1974, Frascati got a phone call from Brookhaven..

But: J/ψ mass was just beyond the energy range of ADONE

So what?
Push magnet currents up...

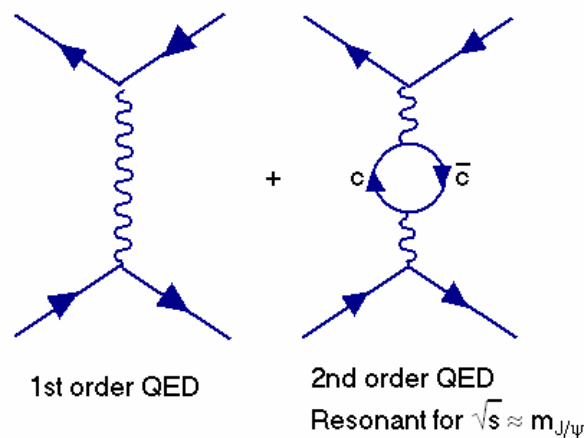


@TBA

J/ψ Quantum Numbers

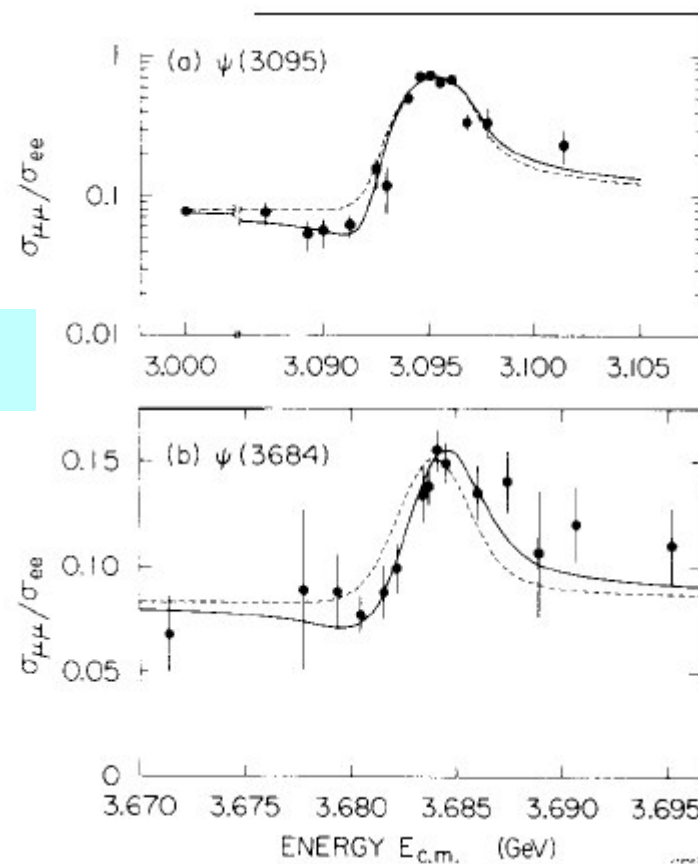
An interesting example of quantum interference

Take the 2 annihilation diagrams:



Interference between the 2 shows up in the total cross-section as a result of the resonant state being $J^{PC}=1^-$ like the photon

Take the ratio to minimize point-to-point luminosity systematics



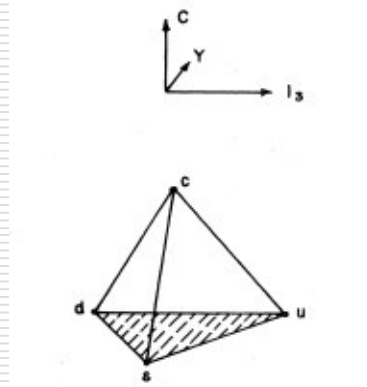
@TBA

$SU(4)$ Multiplets

Fundamental rep. **4, 4*, 6**

$4 \cdot 4 - 1 = 15$ generators

3 fundamental, non equivalent irr. reps.



@TBA

OK for spin 3/2

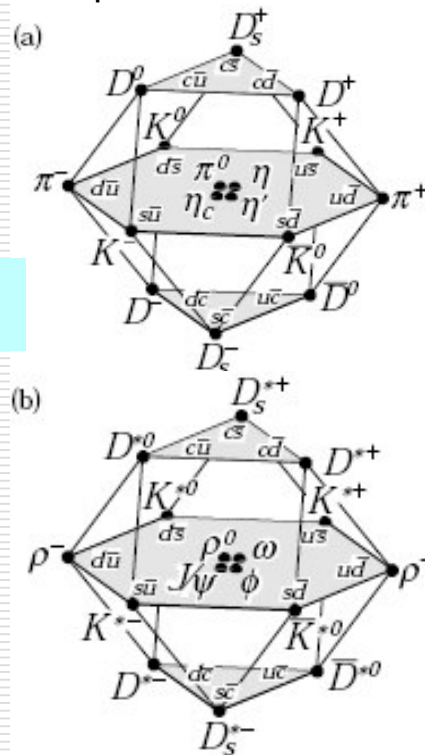
Spin 1/2

$$4 \otimes 4^* = 1 \oplus 15$$

$$4 \otimes 4 \otimes 4 = 20_S \oplus 20_M \oplus 20_M \oplus 4_A$$

Mesons

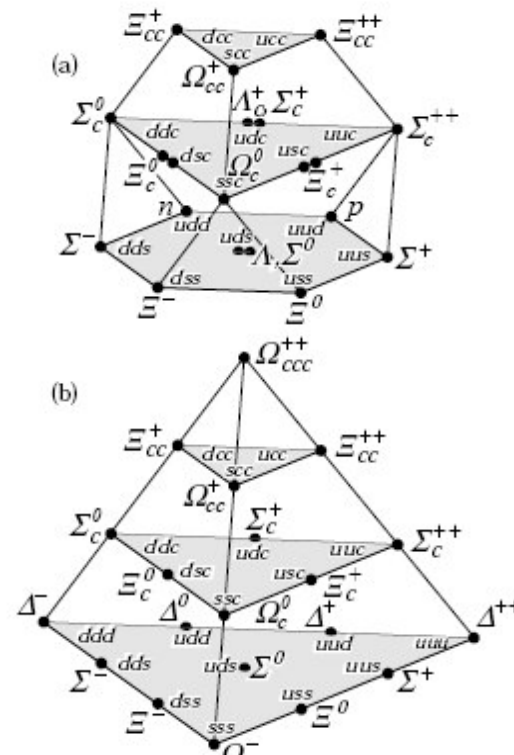
Pseudoscalars



Vectors

Baryons

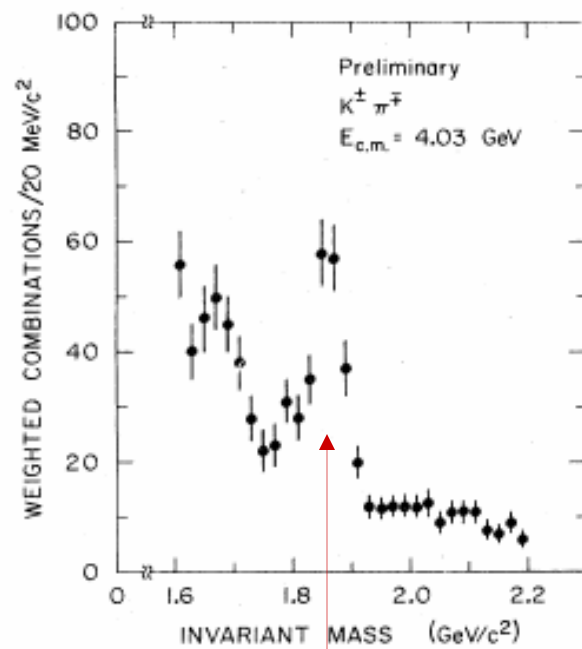
Spin 1/2



Spin 3/2

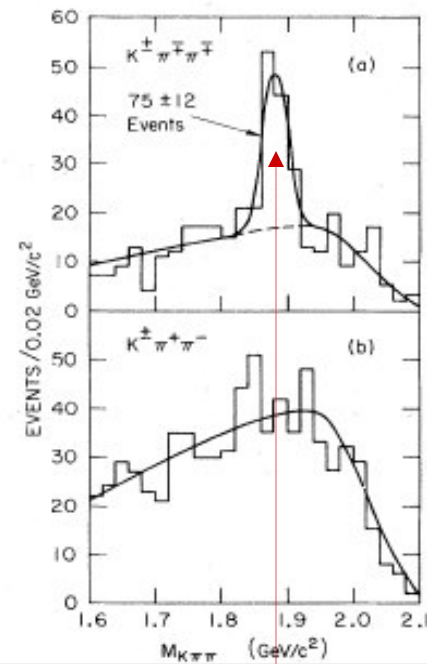
Open Charm

SLAC-LBL Collaboration – Mark I



D^0

@TBA



D^{0*}

The Charmed Zoo

Λ_c^+	****
$\Lambda_c(2593)^+$	***
$\Lambda_c(2625)^+$	***
$\Lambda_c(2765)^+$	*
$\Lambda_c(2880)^+$	**
$\Sigma_c(2455)$	****
$\Sigma_c(2520)$	***
$\Sigma_c(2800)$	***
Ξ_c^+	***
Ξ_c^0	***
$\Xi_c'^+$	***
$\Xi_c'^0$	***
$\Xi_c(2645)$	***
$\Xi_c(2790)$	***
$\Xi_c(2815)$	***
Ω_c^0	***
Ξ_{cc}^+	*

Baryons

@TBA

CHARMED (C = ± 1)	
• $D^\pm$	1/2(0 ⁻)
• D^0	1/2(0 ⁻)
• $D^*(2007)^0$	1/2(1 ⁻)
• $D^*(2010)^\pm$	1/2(1 ⁻)
• $D_0^*(2400)^0$	1/2(0 ⁺)
• $D_0^*(2400)^\pm$	1/2(0 ⁺)
• $D_1(2420)^0$	1/2(1 ⁺)
• $D_1(2420)^\pm$	1/2(1 ⁺)
• $D_1(2430)^0$	1/2(1 ⁺)
• $D_2^*(2460)^0$	1/2(2 ⁺)
• $D_2^*(2460)^\pm$	1/2(2 ⁺)
• $D^*(2640)^\pm$	1/2(1 ⁺)
CHARMED, STRANGE (C = S = ± 1)	
• $D_s^\pm$	0(0 ⁻)
• $D_s^{*\pm}$	0(1 ⁺)
• $D_{s0}^*(2317)^\pm$	0(0 ⁺)
• $D_{s1}(2460)^\pm$	0(1 ⁺)
• $D_{s1}(2536)^\pm$	0(1 ⁺)
• $D_{s2}(2573)^\pm$	0(2 ⁺)

Mesons

Bottom

3rd family (*Bottom*, *Top*) predicted in order to 'explain' (Better: Account for) CP violation (see later) – Kobayashi & Maskawa, 1973

Bound $b\bar{b}$ states first observed at Fermilab in 1977

Discovery subsequently confirmed at e^+e^- machines (DESY, Cornell)

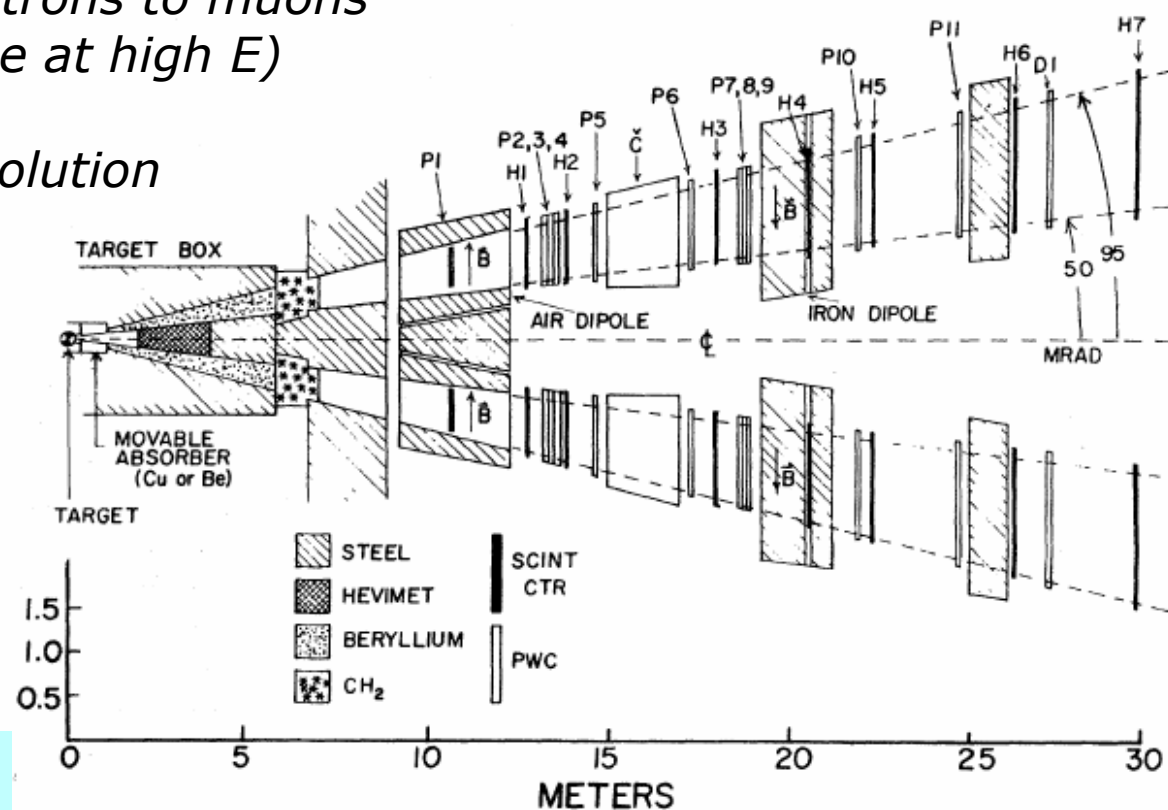
Several b -hadrons observed

Very large b -quark mass $\sim 4\text{-}5$ GeV

Situation similar to charm

The Υ Discovery at FNAL - I

Similar design as J/ψ experiment:
Switch from electrons to muons
(Easier to handle at high E)
High intensity
~Good mass resolution



@TBA

The Y Discovery at FNAL - II

Mass distribution for exclusive process:

$$p + Be \rightarrow \mu^+ + \mu^- + X$$

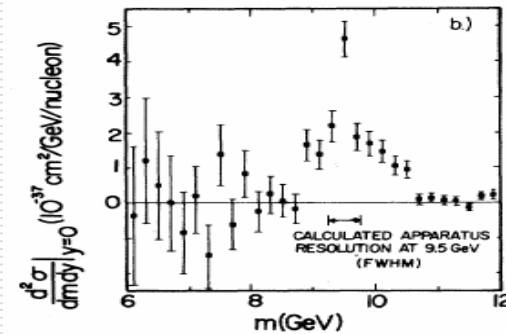
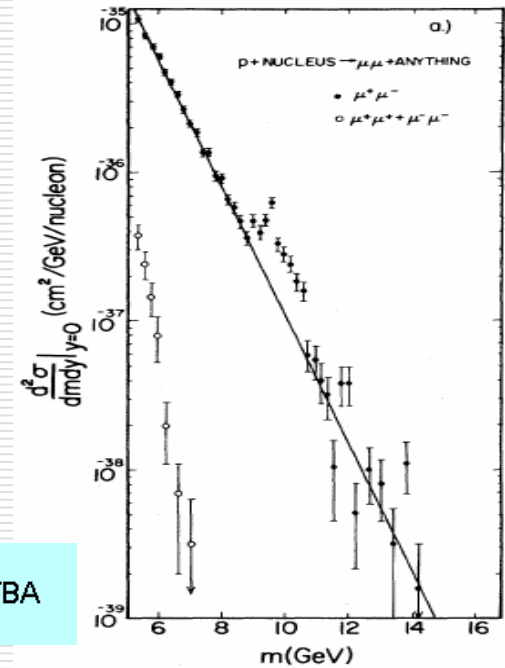
y : Pseudorapidity of the muon pair
(Related to CM angle)

$y=0$ Central region

High mass region shown

Exponential trend + peak

Mass resolution ~ 180 MeV



The Last (?) Zoo

$$\Lambda_b^0 \quad ***$$

$$\Xi_b^0, \Xi_b^- \quad *$$

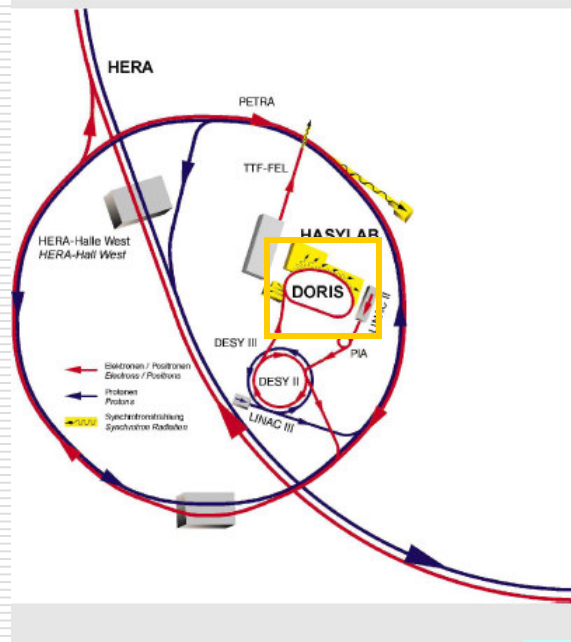
Baryons

@TBA

BOTTOM ($B = \pm 1$) $J^G(J^{PC})$	
• $B^\pm$	$1/2(0^-)$
• B^0	$1/2(0^-)$
• $B^\pm/B^0$ ADMIXTURE	
• $B^\pm/B^0/B_s^0/b$ -baryon ADMIXTURE	
V_{cb} and V_{ub} CKM Matrix Elements	
• B^*	$1/2(1^-)$
$B_J^*(5732)$	$?(?^?)$
BOTTOM, STRANGE ($B = \pm 1, S = \mp 1$)	
• B_s^0	$0(0^-)$
B_s^*	$0(1^-)$
$B_{sJ}^*(5850)$	$?(?^?)$
BOTTOM, CHARMED ($B = C = \pm 1$)	
• $B_c^\pm$	$0(0^-)$

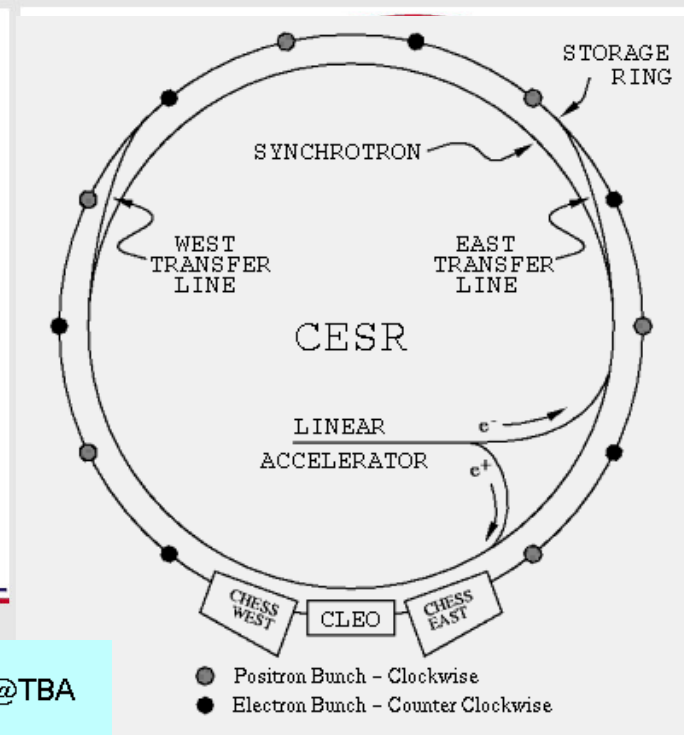
Mesons

DORIS & CESR



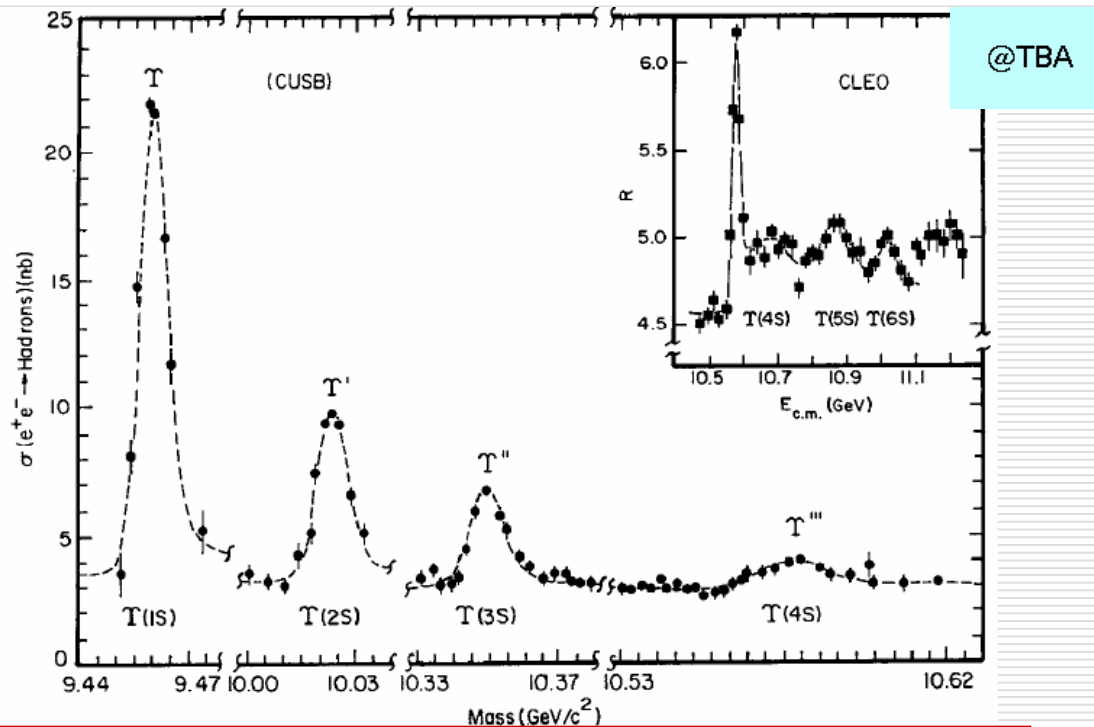
DESY, Hamburg

@TBA

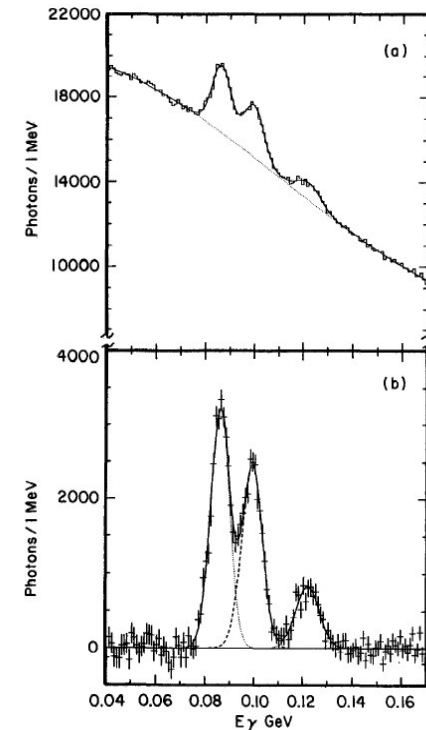


Cornell, Ithaca, NY

The Y Family



3 radial excitations of the Y observed as narrow peaks



Inclusive γ -ray spectrum from $Y(3S)$

ARGUS

One of the first examples
of modern collider detector
design

Large size

(6m $\varnothing$, 6m L: High p resol.)

Vertex chamber

(Aiming to short lifetimes..)

Good EM Calorimetry

(Electron/Photon detection)

Machine improvements

(Low β quads for luminosity)

Muon Chambers

Electromagnetic Calorimeter

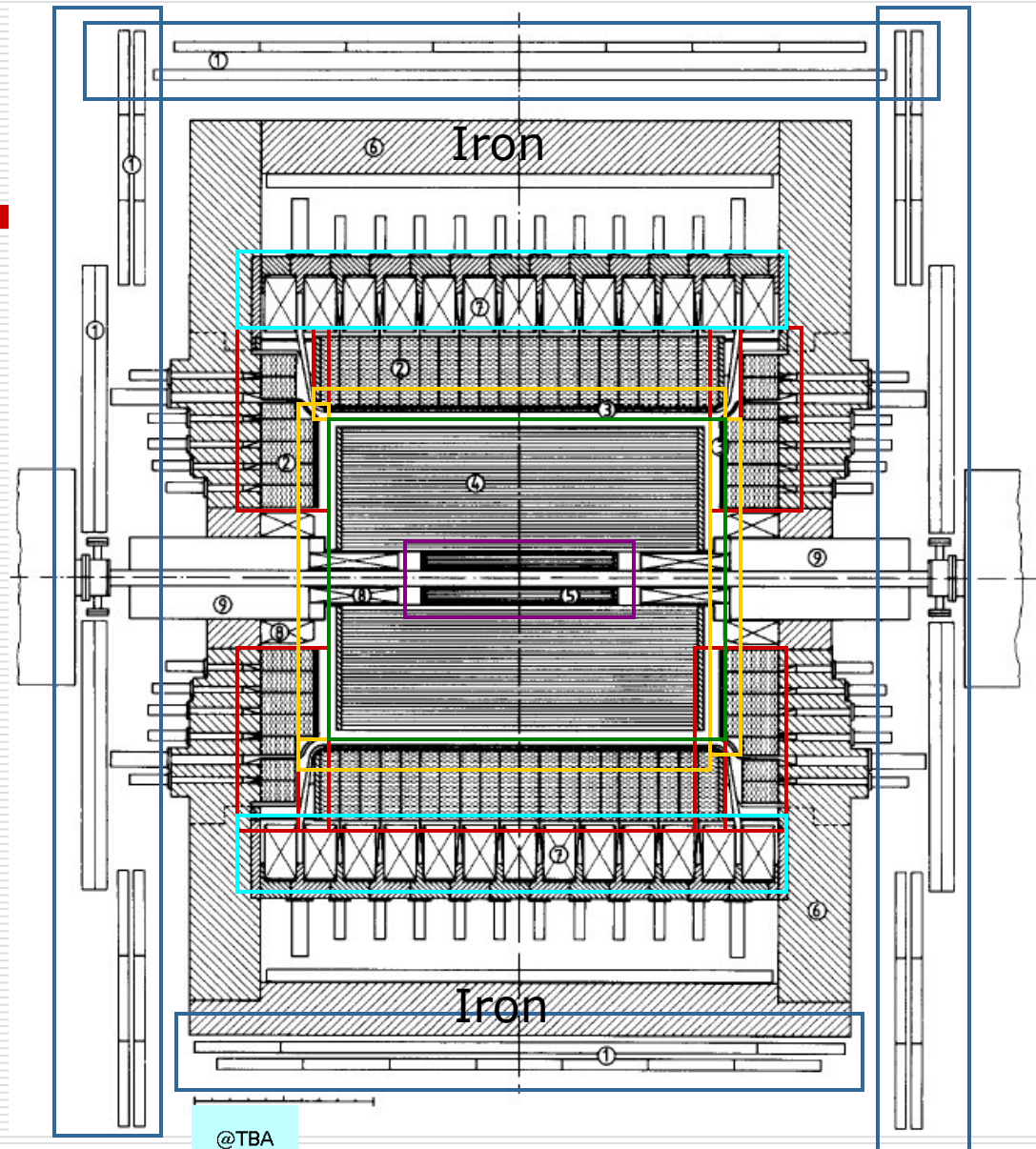
Time of Flight

Drift Chamber

Vertex Detector

Iron Yoke

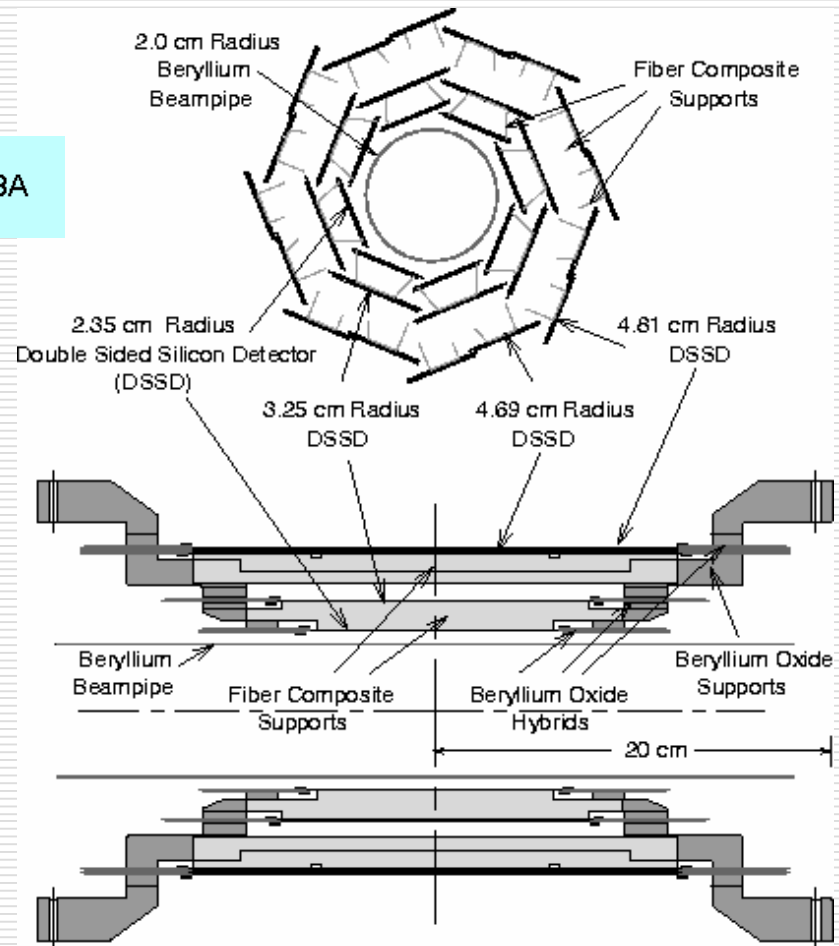
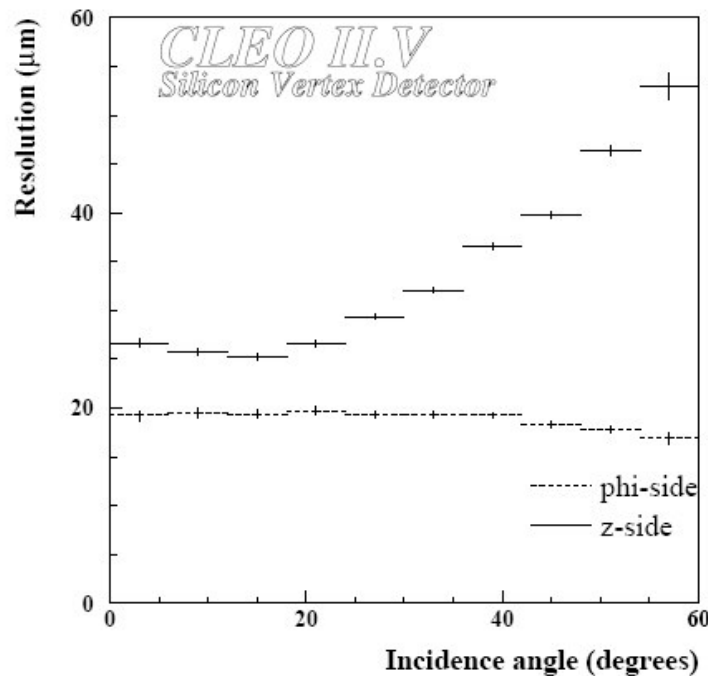
Solenoid



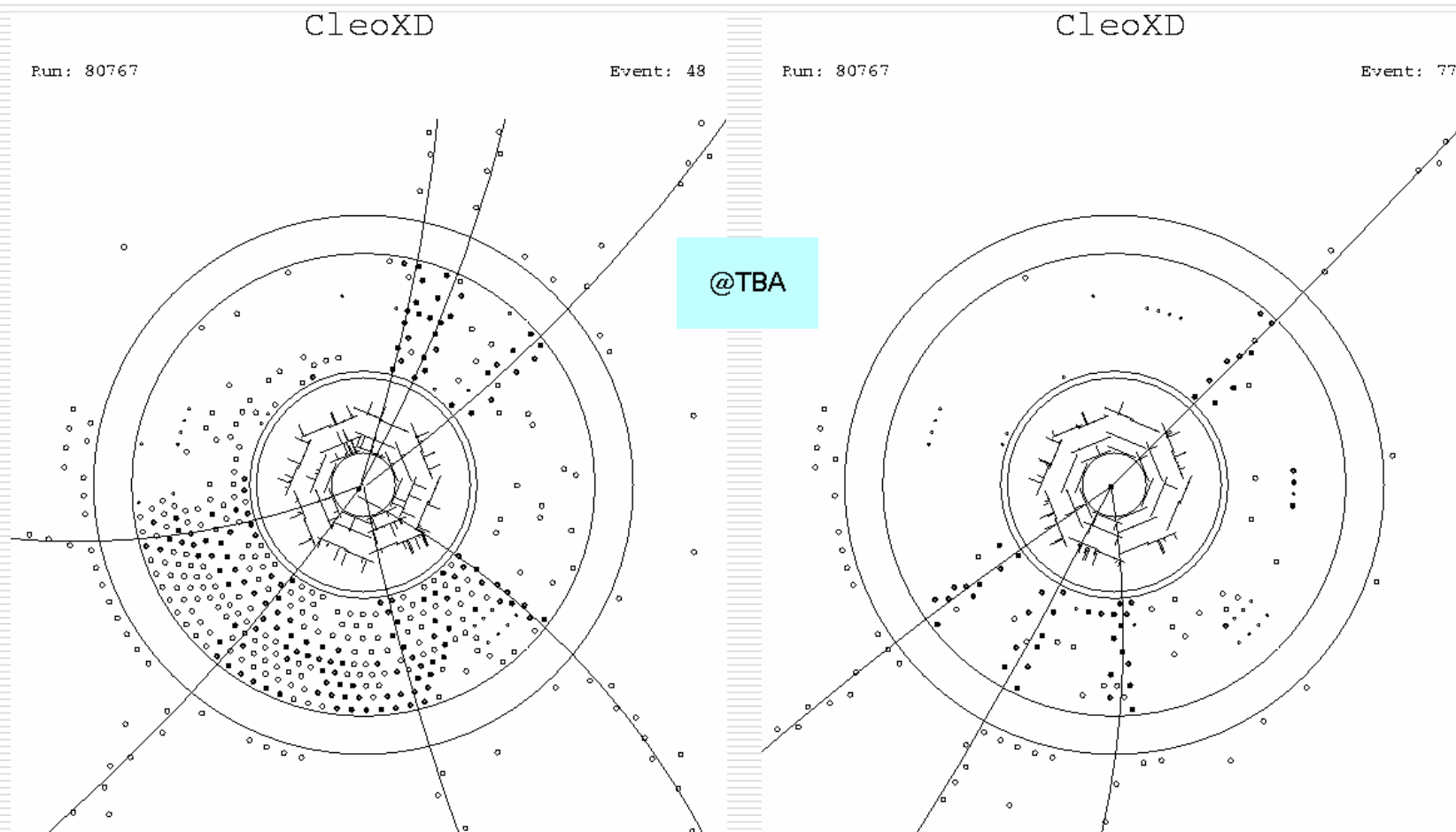
CLEO: The Vertex Detector

20208 strips

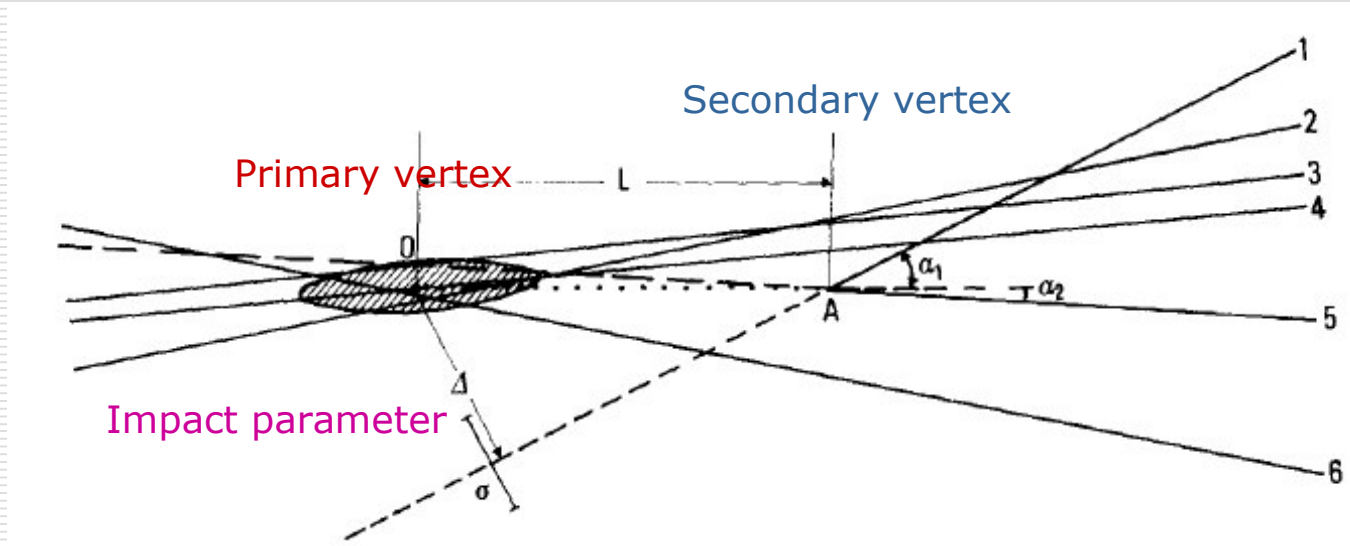
@TBA



Effects on Tracking



Vertex Detection - I



Plane defined by primary vertex, track direction

Consider a particle produced at primary vertex with speed β

When it decays to another particle, call speed β^* , decay angle in CM θ^*

$$\tan \theta = \frac{\sin \theta^*}{\gamma \cos \theta^* + \beta / \beta^*} \quad \text{Lorentz transformation to LAB}$$

Vertex Detection - II

$L = \beta\gamma\tau$ Decay length

Define impact parameter Δ in terms of decay length, L , and angle θ :

$$\Delta = L \sin \theta = L \frac{\tan \theta}{\sqrt{1 + \tan^2 \theta}} = L \frac{\frac{\sin \theta^*}{\gamma(\cos \theta^* + \beta/\beta^*)}}{\sqrt{1 + \left(\frac{\sin \theta^*}{\gamma(\cos \theta^* + \beta/\beta^*)}\right)^2}} = L \frac{\sin \theta^*}{\sqrt{(\gamma(\cos \theta^* + \beta/\beta^*))^2 + \sin^2 \theta^*}}$$

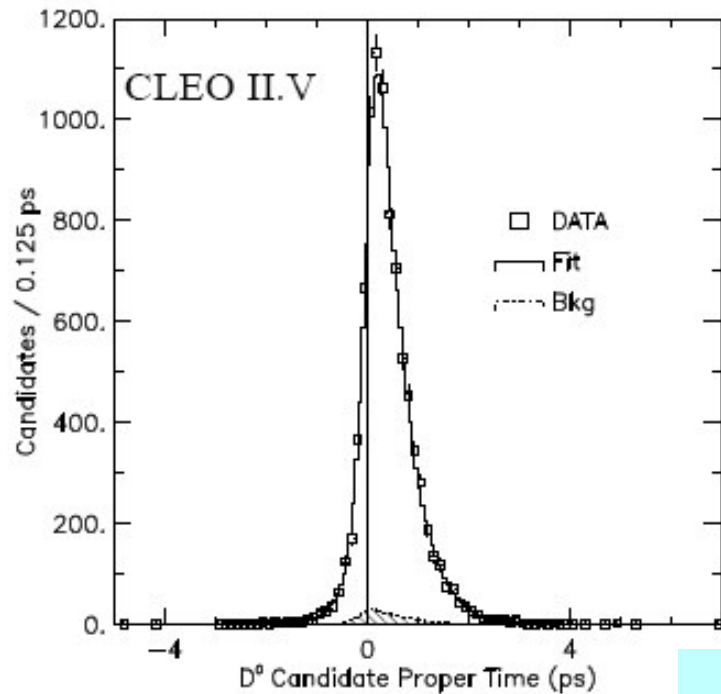
$$\rightarrow \Delta = L \frac{1}{\gamma} \frac{\sin \theta^*}{\sqrt{(\cos \theta^* + \beta/\beta^*)^2 + \frac{1}{\gamma^2} \sin^2 \theta^*}} = \beta\tau \frac{\sin \theta^*}{\sqrt{(\cos \theta^* + \beta/\beta^*)^2 + \frac{1}{\gamma^2} \sin^2 \theta^*}} \xrightarrow{\beta, \beta^* \rightarrow 1} \beta\tau \frac{\sin \theta^*}{1 + \cos \theta^*} = \beta\tau \tan \frac{\theta^*}{2}$$

$y \equiv \frac{\Delta}{\tau} \rightarrow$ Find statistical distribution of y for isotropic θ^* , exponential τ

$\rightarrow \langle y \rangle = \frac{\pi}{2} \rightarrow \langle \Delta \rangle = \frac{\langle \tau \rangle \pi}{2}$ Get a measurement of the decay lifetime

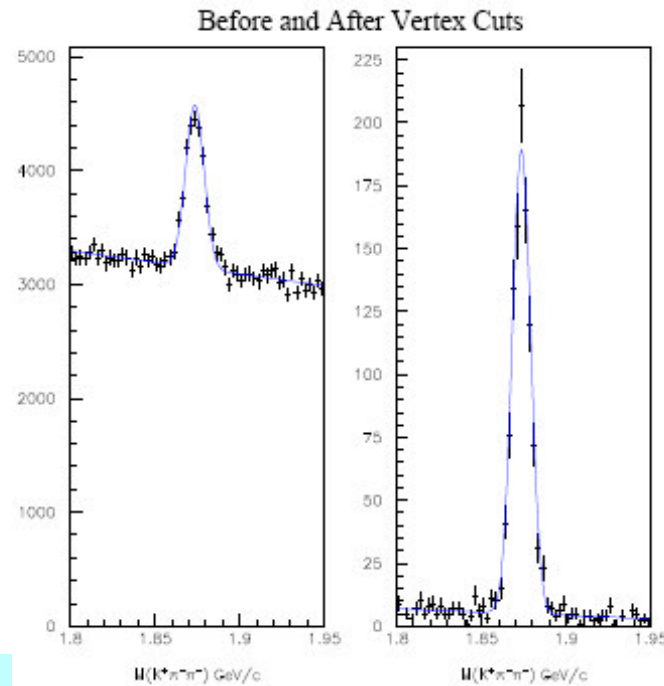
Only from impact parameter, in the limit of relativistic speeds!
Full decay reconstruction not required

Vertex Detection: Charm



D^0 Lifetime

@TBA

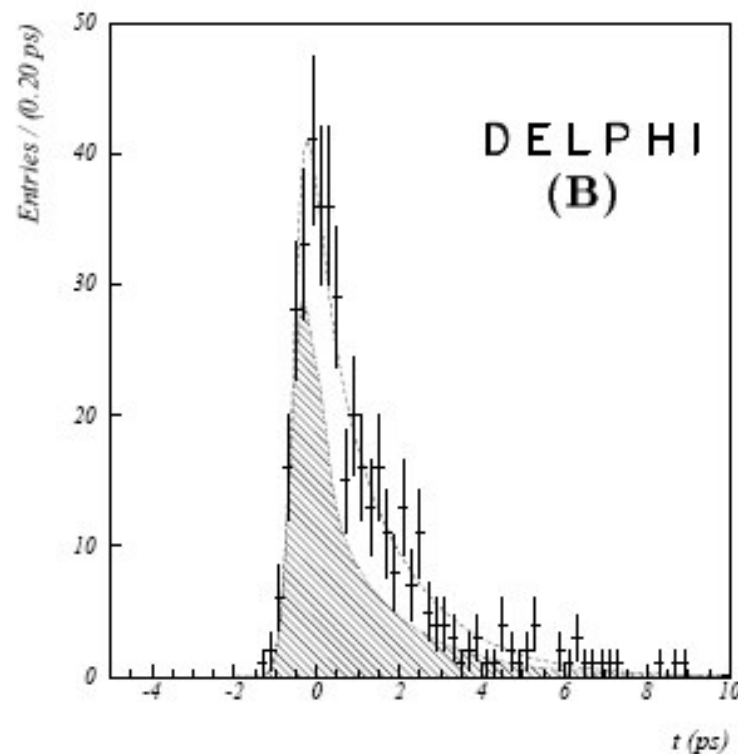


D^* selection,
with and without
secondary vertex

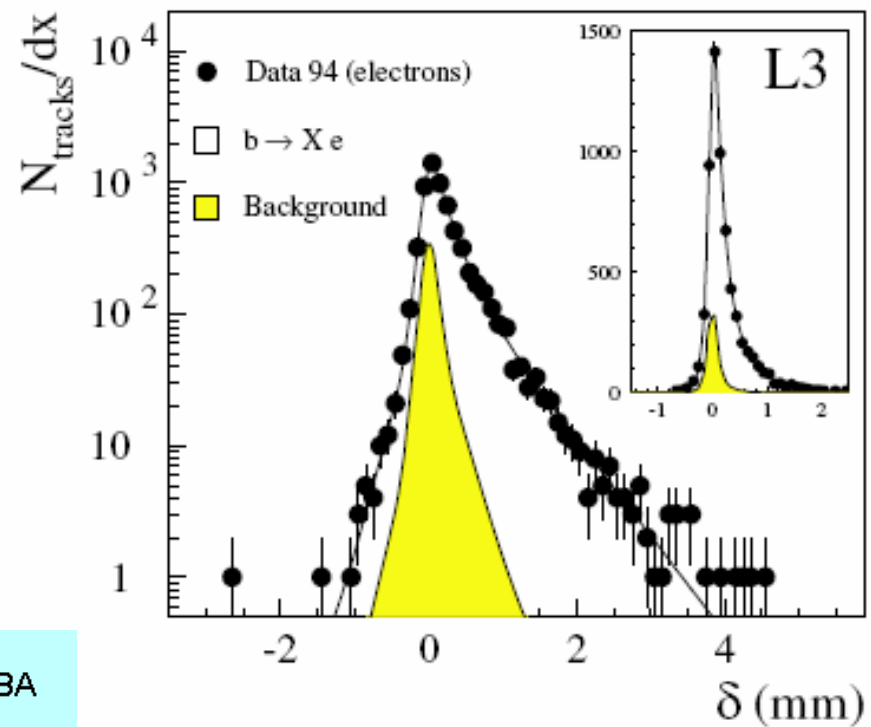
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DELPHI and L3: B Lifetime



@TBA



Top

Heaviest quark, predicted together with b as a member of the 3rd family

Finally found at Fermilab in 1995, after more than 20 years of theoretical and experimental hunting

Very peculiar quark, whose mass of 171 GeV is so large as to allow for weak decays into $b + \text{real } W, Z^0$

→ *Very large weak decay rate, short lifetime similar to strong interaction resonances*

→ *Does not bind into mesons, baryons*

Best understood while discussing weak interactions (see later)