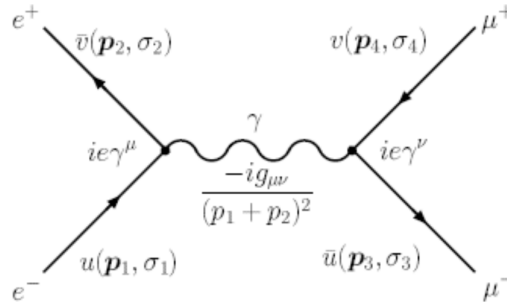


Electron-Positron Annihilation into a Muon-Antimuon Pair
With Everything (Hopefully) Spelled Out



Initial state:

$$e^+ \quad \bar{v}(p_2, \sigma_2)$$

$$e^- \quad u(p_1, \sigma_1)$$

Final state:

$$\mu^+ \quad v(p_4, \sigma_4)$$

$$\mu^- \quad \bar{u}(p_3, \sigma_3)$$

→ Matrix element:

$$T_{fi} = \frac{e^2}{q^2} \bar{u}(p_3, \sigma_3) \gamma_\mu v(p_4, \sigma_4) \bar{v}(p_2, \sigma_2) \gamma^\mu u(p_1, \sigma_1)$$

→ Squared matrix element:

$$|T_{fi}|^2 = \frac{e^4}{q^4} \left[\bar{u}(p_3, \sigma_3) \gamma_\mu v(p_4, \sigma_4) \bar{v}(p_2, \sigma_2) \gamma^\mu u(p_1, \sigma_1) \right] \\ \cdot \left[\bar{u}(p_3, \sigma_3) \gamma_\nu v(p_4, \sigma_4) \bar{v}(p_2, \sigma_2) \gamma^\nu u(p_1, \sigma_1) \right]^*$$

Observe:

$$\left[\bar{u}(p_3, \sigma_3) \gamma_\nu v(p_4, \sigma_4) \bar{v}(p_2, \sigma_2) \gamma^\nu u(p_1, \sigma_1) \right]^* =$$

$$\left[\bar{u}(p_3, \sigma_3) \gamma_\nu v(p_4, \sigma_4) \right]^* \left[\bar{v}(p_2, \sigma_2) \gamma^\nu u(p_1, \sigma_1) \right]^*$$

$$\text{Complex number} = 1 \times 1 \text{ 'matrix'} \rightarrow []^* \equiv []^\dagger$$

$$\rightarrow \left[\bar{u}(p_3, \sigma_3) \gamma_\nu v(p_4, \sigma_4) \right]^* = \left[u^\dagger(p_3, \sigma_3) \gamma_0 \gamma_\nu v(p_4, \sigma_4) \right]^\dagger = \left[v^\dagger(p_4, \sigma_4) \gamma_\nu^\dagger \gamma_0^\dagger u(p_3, \sigma_3) \right]$$

$$\gamma_0^2 = 1, \gamma_0^\dagger = \gamma_0, \bar{\gamma}_\nu \equiv \gamma_0 \gamma_\nu^\dagger \gamma_0$$

$$\rightarrow \left[v^\dagger(p_4, \sigma_4) \gamma_\nu^\dagger \gamma_0^\dagger u(p_3, \sigma_3) \right] = \left[v^\dagger(p_4, \sigma_4) \gamma_0^2 \gamma_\nu^\dagger \gamma_0 u(p_3, \sigma_3) \right]$$

$$= \left[\bar{v}(p_4, \sigma_4) \gamma_0 \gamma_\nu^\dagger \gamma_0 u(p_3, \sigma_3) \right] = \left[\bar{v}(p_4, \sigma_4) \bar{\gamma}_\nu u(p_3, \sigma_3) \right]$$

$$\rightarrow \left[\bar{u}(p_3, \sigma_3) \gamma_\nu v(p_4, \sigma_4) \right]^* = \left[\bar{v}(p_4, \sigma_4) \bar{\gamma}_\nu u(p_3, \sigma_3) \right]$$

Similarly:

$$\left[\bar{v}(p_2, \sigma_2) \gamma^\nu u(p_1, \sigma_1) \right]^* = \left[\bar{u}(p_1, \sigma_1) \bar{\gamma}^\nu v(p_2, \sigma_2) \right]$$

$$\rightarrow |T_{fi}|^2 = \frac{e^4}{q^4} \left[\bar{u}(p_3, \sigma_3) \gamma_\mu v(p_4, \sigma_4) \bar{v}(p_2, \sigma_2) \gamma^\mu u(p_1, \sigma_1) \right]$$

$$\cdot \left[\bar{v}(p_4, \sigma_4) \bar{\gamma}_\nu u(p_3, \sigma_3) \bar{u}(p_1, \sigma_1) \bar{\gamma}^\nu v(p_2, \sigma_2) \right]$$

Averaging and summing over initial and final spins:

$$\langle |T_{fi}|^2 \rangle = \frac{1}{4} \sum_{\sigma_1, \sigma_2, \sigma_3, \sigma_4} |T_{fi}|^2$$

$$\langle |T_{fi}|^2 \rangle = \frac{1}{4} \frac{e^4}{q^4} \sum_{\sigma_1, \sigma_2, \sigma_3, \sigma_4} \left[\bar{u}(p_3, \sigma_3) \gamma_\mu v(p_4, \sigma_4) \bar{v}(p_2, \sigma_2) \gamma^\mu u(p_1, \sigma_1) \right]$$

$$\cdot \left[\bar{v}(p_4, \sigma_4) \bar{\gamma}_\nu u(p_3, \sigma_3) \bar{u}(p_1, \sigma_1) \bar{\gamma}^\nu v(p_2, \sigma_2) \right]$$

$$\langle |T_{fi}|^2 \rangle = \frac{1}{4} \frac{e^4}{q^4} \sum_{\sigma_1, \sigma_2, \sigma_3, \sigma_4} \left[\bar{u}(p_3, \sigma_3) \gamma_\mu v(p_4, \sigma_4) \right] \left[\bar{v}(p_2, \sigma_2) \gamma^\mu u(p_1, \sigma_1) \right]$$

$$\cdot \left[\bar{v}(p_4, \sigma_4) \bar{\gamma}_\nu u(p_3, \sigma_3) \right] \left[\bar{u}(p_1, \sigma_1) \bar{\gamma}^\nu v(p_2, \sigma_2) \right]$$

$$\langle |T_{fi}|^2 \rangle = \frac{1}{4} \frac{e^4}{q^4} \sum_{\sigma_1, \sigma_2} \left[\bar{u}(p_1, \sigma_1) \bar{\gamma}^\nu v(p_2, \sigma_2) \right] \left[\bar{v}(p_2, \sigma_2) \gamma^\mu u(p_1, \sigma_1) \right]$$

$$\cdot \sum_{\sigma_3, \sigma_4} \left[\bar{u}(p_3, \sigma_3) \gamma_\mu v(p_4, \sigma_4) \right] \left[\bar{v}(p_4, \sigma_4) \bar{\gamma}_\nu u(p_3, \sigma_3) \right]$$

This sum of products can be reduced in several ways. Simplest trick:

a) Write explicitly spinor indexes:

$$\sum_{\sigma_1, \sigma_2} \sum_{\alpha, \beta, \gamma, \delta} \bar{u}_\alpha(p_1, \sigma_1) \bar{\gamma}_{(\alpha\beta)}^\nu v_\beta(p_2, \sigma_2) \bar{v}_\gamma(p_2, \sigma_2) \gamma_{(\gamma\delta)}^\mu u_\delta(p_1, \sigma_1) \quad (\& \text{ similar for } \sum_{\sigma_3, \sigma_4})$$

b) Move u_δ to the front in each term (OK because it is just a number)

$$\sum_{\sigma_1, \sigma_2} \sum_{\alpha, \beta, \gamma, \delta} u_\delta(p_1, \sigma_1) \bar{u}_\alpha(p_1, \sigma_1) \bar{\gamma}_{(\alpha\beta)}^\nu v_\beta(p_2, \sigma_2) \bar{v}_\gamma(p_2, \sigma_2) \gamma_{(\gamma\delta)}^\mu$$

c) Exchange the sum order and split the spin sums:

$$\sum_{\alpha, \beta, \gamma, \delta} \sum_{\sigma_1} u_\delta(p_1, \sigma_1) \bar{u}_\alpha(p_1, \sigma_1) \bar{\gamma}_{(\alpha\beta)}^\nu \sum_{\sigma_2} v_\beta(p_2, \sigma_2) \bar{v}_\gamma(p_2, \sigma_2) \gamma_{(\gamma\delta)}^\mu$$

d) Observe the appearance of

$$\sum_{\sigma_1} u_\delta(p_1, \sigma_1) \bar{u}_\alpha(p_1, \sigma_1) = (\not{p}_1 + m_1)_{\delta\alpha}$$

$$\sum_{\sigma_2} v_\beta(p_2, \sigma_2) \bar{v}_\gamma(p_2, \sigma_2) = (\not{p}_2 - m_2)_{\beta\gamma}$$

e) Then

$$\sum_{\alpha, \beta, \gamma, \delta} (\not{p}_1 + m_1)_{\delta\alpha} \bar{\gamma}_{(\alpha\beta)}^\nu (\not{p}_2 - m_2)_{\beta\gamma} \gamma_{(\gamma\delta)}^\mu = \text{Tr}[(\not{p}_1 + m_1) \bar{\gamma}^\nu (\not{p}_2 - m_2) \gamma^\mu]$$

Similarly for the sum over σ_3, σ_4 :

$$\sum_{\alpha, \beta, \gamma, \delta} (\not{p}_3 + m_3)_{\delta\alpha} \bar{\gamma}_{\nu(\alpha\beta)} (\not{p}_4 - m_4)_{\beta\gamma} \gamma_{\mu(\gamma\delta)} = \text{Tr}[(\not{p}_3 + m_3) \bar{\gamma}_\nu (\not{p}_4 - m_4) \gamma_\mu]$$

Now

$\bar{\gamma}_\nu \equiv \gamma_0 \gamma_\nu^\dagger \gamma_0 = \gamma_\nu$ by Dirac matrices algebra

$$\begin{aligned} \rightarrow \langle |T_{fi}|^2 \rangle &= \frac{1}{4} \frac{e^4}{q^4} \text{Tr}[(\not{p}_1 + m_1) \gamma^\nu (\not{p}_2 - m_2) \gamma^\mu] \text{Tr}[(\not{p}_3 + m_3) \gamma_\nu (\not{p}_4 - m_4) \gamma_\mu] \\ &= \frac{1}{4} \frac{e^4}{q^4} \text{Tr}[\not{p}_1 \gamma^\nu \not{p}_2 \gamma^\mu - \not{p}_1 \gamma^\nu m_2 \gamma^\mu + m_1 \gamma^\nu \not{p}_2 \gamma^\mu - m_1 \gamma^\mu m_2 \gamma^\nu] \\ &\quad \cdot \text{Tr}[(\not{p}_3 \gamma_\nu \not{p}_4 \gamma_\mu - \not{p}_3 \gamma_\nu m_4 \gamma_\mu + m_3 \gamma_\nu \not{p}_4 \gamma_\mu - m_3 \gamma_\nu m_4 \gamma_\mu)] \end{aligned}$$

$$m_1 = m_2 = m, m_3 = m_4 = M$$

$$\begin{aligned} \rightarrow \langle |T_{fi}|^2 \rangle &= \frac{1}{4} \frac{e^4}{q^4} \text{Tr}[\not{p}_1 \gamma^\nu \not{p}_2 \gamma^\mu - m \not{p}_1 \gamma^\nu \gamma^\mu + m \gamma^\nu \not{p}_2 \gamma^\mu - m^2 \gamma^\mu \gamma^\nu] \\ &\quad \cdot \text{Tr}[(\not{p}_3 \gamma_\nu \not{p}_4 \gamma_\mu - M \not{p}_3 \gamma_\nu \gamma_\mu + M \gamma_\nu \not{p}_4 \gamma_\mu - M^2 \gamma_\nu \gamma_\mu)] \end{aligned}$$

Trace of product of odd number of γ matrices = 0

$$\rightarrow \text{Tr}(\not{p}_1 \gamma^\nu \gamma^\mu) = \text{Tr}(\gamma^\nu \not{p}_2 \gamma^\mu) = \text{Tr}(\not{p}_3 \gamma_\nu \gamma_\mu) = \text{Tr}(\gamma_\nu \not{p}_4 \gamma_\mu) = 0$$

$$\rightarrow \langle |T_{fi}|^2 \rangle = \frac{1}{4} \frac{e^4}{q^4} \text{Tr}[\not{p}_1 \gamma^\nu \not{p}_2 \gamma^\mu - m^2 \gamma^\mu \gamma^\nu] \cdot \text{Tr}[(\not{p}_3 \gamma_\nu \not{p}_4 \gamma_\mu - M^2 \gamma_\nu \gamma_\mu)]$$

$$\text{Tr}(\gamma_\nu \gamma_\mu) = 4g_{\nu\mu}, \text{ similar for contravariant}$$

$$\rightarrow \langle |T_{fi}|^2 \rangle = \frac{1}{4} \frac{e^4}{q^4} [\text{Tr}(\not{p}_1 \gamma^\nu \not{p}_2 \gamma^\mu) - 4m^2 g^{\mu\nu}] \cdot [\text{Tr}(\not{p}_3 \gamma_\nu \not{p}_4 \gamma_\mu) - 4M^2 g_{\nu\mu}]$$

$$\text{Tr}(\not{p}_3 \gamma_\nu \not{p}_4 \gamma_\mu) = \text{Tr}(p_{3\alpha} \gamma_\alpha \gamma_\nu p_{4\beta} \gamma_\beta \gamma_\mu) = p_{3\alpha} p_{4\beta} \text{Tr}(\gamma_\alpha \gamma_\nu \gamma_\beta \gamma_\mu)$$

$$\rightarrow \text{Tr}(\not{p}_3 \gamma_\nu \not{p}_4 \gamma_\mu) = p_{3\alpha} p_{4\beta} 4(g_{\alpha\nu} g_{\beta\mu} - g_{\alpha\beta} g_{\nu\mu} + g_{\alpha\mu} g_{\nu\beta})$$

$$\rightarrow \text{Tr}(\not{p}_3 \gamma_\nu \not{p}_4 \gamma_\mu) = 4p_{3\nu} p_{4\mu} + 4p_{3\mu} p_{4\nu} - 4p_3 p_4 g_{\nu\mu}$$

Similarly:

$$\text{Tr}(\not{p}_1 \gamma^\nu \not{p}_2 \gamma^\mu) = 4p^{1\nu} p^{2\mu} + 4p^{1\nu} p^{2\mu} - 4p_1 p_2 g^{\nu\mu}$$

$$\begin{aligned} \rightarrow \langle |T_{fi}|^2 \rangle &= \frac{1}{4} \frac{e^4}{q^4} [4p_1^\nu p_2^\mu + 4p_1^\mu p_2^\nu - 4(p_1 p_2 + m^2) g^{\mu\nu}] \\ &\quad \cdot [4p_{3\nu} p_{4\mu} + 4p_{3\mu} p_{4\nu} - 4(p_3 p_4 + M^2) g_{\nu\mu}] \end{aligned}$$

$$s = (p_1 + p_2)^2 = p_1^2 + p_2^2 + 2p_1 p_2 = 2m^2 + 2p_1 p_2$$

$$\rightarrow 4(p_1 p_2 + m^2) = 4(p_3 p_4 + M^2) = 2s$$

$$\rightarrow \langle |T_{fi}|^2 \rangle = \frac{1}{4} \frac{e^4}{q^4} [4p_1^\nu p_2^\mu + 4p_1^\mu p_2^\nu - 2s g^{\mu\nu}] \cdot [4p_{3\nu} p_{4\mu} + 4p_{3\mu} p_{4\nu} - 2s g_{\nu\mu}]$$

$$\rightarrow \langle |T_{fi}|^2 \rangle = \frac{1}{4} \frac{e^4}{q^4} [32(p_1 p_4)(p_2 p_3) + 32(p_1 p_3)(p_2 p_4) - 16s(p_1 p_2) - 16s(p_4 p_3) + 16s^2]$$

$$\rightarrow \langle |T_{fi}|^2 \rangle = \frac{1}{4} \frac{e^4}{q^4} [32(p_1 p_4)(p_2 p_3) + 32(p_1 p_3)(p_2 p_4)$$

$$+ 16s \left[\underbrace{(p_1 p_2 + m^2)}_{s/2} + \underbrace{(p_4 p_3 + M^2)}_{s/2} - (p_1 p_2) - (p_4 p_3) \right]]$$

$$q^2 \equiv s$$

$$\rightarrow \langle |T_{fi}|^2 \rangle = \frac{1}{4} \frac{e^4}{s^2} [32(p_1 p_4)(p_2 p_3) + 32(p_1 p_3)(p_2 p_4) + 16s(m^2 + M^2)]$$

$$\rightarrow \langle |T_{fi}|^2 \rangle = \frac{4e^4}{s^2} [2(p_1 p_4)(p_2 p_3) + 2(p_1 p_3)(p_2 p_4) + s(m^2 + M^2)]$$

When $m \approx 0$:

$$\langle |T_{fi}|^2 \rangle = \frac{4e^4}{s^2} [2(p_1 p_4)(p_2 p_3) + 2(p_1 p_3)(p_2 p_4) + sM^2]$$

When $m \approx 0, M \approx 0$:

$$\langle |T_{fi}|^2 \rangle = \frac{8e^4}{s^2} [(p_1 p_4)(p_2 p_3) + (p_1 p_3)(p_2 p_4)]$$

By neglecting m :

$$CM: \begin{cases} |\mathbf{p}_1| = |\mathbf{p}_2| = |\mathbf{p}_{in}| = E \\ |\mathbf{p}_3| = |\mathbf{p}_4| = |\mathbf{p}_{out}| = \sqrt{E^2 - M^2} \end{cases}$$

$$\rightarrow \begin{cases} p_1 = (E \ 0 \ 0 \ E), p_2 = (E \ 0 \ 0 \ -E) \\ p_3 = (E \ |\mathbf{p}_{out}| \sin \theta \ 0 \ |\mathbf{p}_{out}| \cos \theta), p_4 = (E \ -|\mathbf{p}_{out}| \sin \theta \ 0 \ -|\mathbf{p}_{out}| \cos \theta) \end{cases}$$

$$\begin{cases} p_1 p_4 = E^2 + E\sqrt{E^2 - M^2} \cos \theta \\ p_2 p_3 = E^2 + E\sqrt{E^2 - M^2} \cos \theta \end{cases}$$

$$\rightarrow p_1 p_4 p_2 p_3 = \left(E^2 + E\sqrt{E^2 - M^2} \cos \theta \right)^2 = E^4 + E^2 (E^2 - M^2) \cos^2 \theta + 2E^2 E\sqrt{E^2 - M^2} \cos \theta$$

$$\rightarrow p_1 p_4 p_2 p_3 = E^4 \left[1 + \left(1 - \frac{M^2}{E^2} \right) \cos^2 \theta + 2\sqrt{1 - \frac{M^2}{E^2}} \cos \theta \right]$$

$$\begin{cases} p_1 p_3 = E^2 - E\sqrt{E^2 - M^2} \cos \theta \\ p_2 p_4 = E^2 - E\sqrt{E^2 - M^2} \cos \theta \end{cases} \rightarrow p_1 p_3 p_2 p_4 = \left(E^2 - E\sqrt{E^2 - M^2} \cos \theta \right)^2$$

$$\rightarrow p_1 p_3 p_2 p_4 = E^4 \left[1 + \left(1 - \frac{M^2}{E^2} \right) \cos^2 \theta - 2\sqrt{1 - \frac{M^2}{E^2}} \cos \theta \right]$$

$$\rightarrow \langle |T_{fi}|^2 \rangle = \frac{8e^4}{s^2} \left[(p_1 p_4)(p_2 p_3) + (p_1 p_3)(p_2 p_4) + \frac{1}{2} s M^2 \right]$$

$$\rightarrow \langle |T_{fi}|^2 \rangle = \frac{8e^4}{16E^4} \left[\begin{aligned} & E^4 \left[1 + \left(1 - \frac{M^2}{E^2} \right) \cos^2 \theta + 2\sqrt{1 - \frac{M^2}{E^2}} \cos \theta \right] + \\ & + E^4 \left[1 + \left(1 - \frac{M^2}{E^2} \right) \cos^2 \theta - 2\sqrt{1 - \frac{M^2}{E^2}} \cos \theta \right] \\ & + 2E^2 M^2 \end{aligned} \right]$$

$$\rightarrow \langle |T_{fi}|^2 \rangle = \frac{e^4}{2} \left[2 + 2 \left(1 - \frac{M^2}{E^2} \right) \cos^2 \theta + 4 \frac{M^2}{E^2} \right]$$

$$\rightarrow \langle |T_{fi}|^2 \rangle = \frac{e^4}{2} \left[2 + 2 \left(1 - \frac{M^2}{E^2} \right) \cos^2 \theta + 2 \frac{M^2}{E^2} \right]$$

$$\rightarrow \langle |T_{fi}|^2 \rangle = e^4 \left[1 + \frac{M^2}{E^2} + \left(1 - \frac{M^2}{E^2} \right) \cos^2 \theta \right]$$

Therefore, for the cross-section in the CM frame:

$$\begin{aligned} \frac{d\sigma}{d\Omega} &= \frac{(2\pi)^{4-2\cdot 3}}{4[(p_1 \cdot p_2)^2 - m^4]^{1/2}} \langle |T_{fi}|^2 \rangle \underbrace{\frac{|\mathbf{p}_f|}{4\sqrt{s}}}_{\text{phase space factor}} = \frac{1}{4\pi^2} \frac{1}{4 \cdot 2E^2} \langle |T_{fi}|^2 \rangle \frac{\sqrt{E^2 - M^2}}{4(2E)} \\ &\rightarrow \frac{d\sigma}{d\Omega} = \frac{1}{32\pi^2 E^2} \langle |T_{fi}|^2 \rangle \frac{1}{8} \sqrt{1 - \frac{M^2}{E^2}} \\ &\rightarrow \frac{d\sigma}{d\Omega} = \frac{e^4}{32\pi^2 E^2} \frac{\sqrt{E^2 - M^2}}{E} \frac{1}{8} \left[1 + \frac{M^2}{E^2} + \left(1 - \frac{M^2}{E^2} \right) \cos^2 \theta \right] \\ \alpha &= \frac{e^2}{4\pi} \rightarrow \frac{e^4}{16\pi^2} = \alpha^2 \\ &\rightarrow \frac{d\sigma}{d\Omega} = \frac{\alpha^2}{16E^2} \sqrt{1 - \frac{M^2}{E^2}} \left[1 + \frac{M^2}{E^2} + \left(1 - \frac{M^2}{E^2} \right) \cos^2 \theta \right] \end{aligned}$$

Reminder on 2 - body phase space :

P total 4 - momentum, CM components: $(\sqrt{s} \ 0 \ 0 \ 0)$

p_1, p_2, m_1, m_2 4-momentum & mass of final state particles

$$d^2\Phi_2 = \int \frac{d^3\mathbf{p}_1}{2E_1} \frac{d^3\mathbf{p}_2}{2E_2} \delta^4(P - p_1 - p_2)$$

$$\rightarrow d^2\Phi_2 = \int \delta^4(P - p_1 - p_2) \delta[p_1^2 - m_1^2] \delta[p_2^2 - m_2^2] d^4p_1 d^4p_2$$

$$\rightarrow d^2\Phi_2 = \int \delta^4(P - p_1 - p_2) \delta[p_1^2 - m_1^2] \delta[p_2^2 - m_2^2] d^4p_1 d^4p_2$$

$$\rightarrow d^2\Phi_2 = \int \delta[p_1^2 - m_1^2] \delta[(P - p_1)^2 - m_2^2] d^4p_1$$

$$\delta[(P - p_1)^2 - m_2^2] = \delta\left[(\sqrt{s} - E_1^*)^2 - \mathbf{p}_1^{*2} - m_2^2\right] = \delta\left[(s + E_1^{*2} - 2\sqrt{s}E_1^*) - \mathbf{p}_1^{*2} - m_2^2\right]$$

$$= \delta\left[s - 2\sqrt{s}E_1^* + m_1^2 - m_2^2\right]$$

$$\rightarrow d^2\Phi_2 = \int \frac{d^3\mathbf{p}_1^*}{2E_1^*} \delta\left[s - 2\sqrt{s}E_1^* + m_1^2 - m_2^2\right]$$

$$E^2 = p^2 + m^2 \rightarrow EdE = pdp \rightarrow dp = \frac{E}{p} dE$$

$$d^3 \mathbf{p} = p^2 dp d\Omega = p^2 \frac{E}{p} dE d\Omega = p E dE d\Omega$$

$$\rightarrow d^2 \Phi_2 = \int \frac{d^3 \mathbf{p}_1^*}{2E_1^*} \delta \left[d - 2\sqrt{s} E_1^* + m_1^2 - m_2^2 \right] = \int \frac{p_1^* E_1^*}{2E_1^*} \delta \left[s - 2\sqrt{s} E_1^* + m_1^2 - m_2^2 \right] dE_1^* d\Omega_1^*$$

$$\rightarrow d^2 \Phi_2 = \int \frac{1}{2} p_1^* \delta \left[s - 2\sqrt{s} E_1^* + m_1^2 - m_2^2 \right] dE_1^* d\Omega_1^*$$

$$E_1^* = \frac{s + m_1^2 - m_2^2}{2\sqrt{s}} \rightarrow \delta \left[s - 2\sqrt{s} E_1^* + m_1^2 - m_2^2 \right] = \frac{1}{2\sqrt{s}} \delta \left[E_1^* - \frac{s + m_1^2 - m_2^2}{2\sqrt{s}} \right]$$

$$p_1^* = \sqrt{E_1^{*2} - m_1^2} \equiv p^* = \sqrt{\frac{(s + m_1^2 - m_2^2)^2 - 4m_1^2 s}{4s}} = \sqrt{\frac{(s - m_1^2 - m_2^2)^2}{4s}} = \frac{s - m_1^2 - m_2^2}{2\sqrt{s}}$$

$$\rightarrow d^2 \Phi_2 = \frac{p^*}{\underbrace{4\sqrt{s}}_{\text{phase space factor}}} d\Omega_1^*$$

Reminder on Dirac representation:

$$\begin{cases} \gamma_0^\dagger = \gamma_0 \\ \gamma_k^\dagger = -\gamma_k \end{cases}$$

$$\rightarrow \gamma_\nu^\dagger = \gamma_0 \gamma_\nu \gamma_0$$

$$\rightarrow \gamma_0 \gamma_\nu^\dagger \gamma_0 = \gamma_0 \gamma_0 \gamma_\nu \gamma_0 \gamma_0$$

$$\gamma_0^2 = 1$$

$$\rightarrow \gamma_0 \gamma_\nu^\dagger \gamma_0 = \gamma_\nu$$