

Oscillatore armonico

$$\frac{d^2u}{dt^2} + \omega_0^2 u = 0$$

Punto di partenza: Teorema di esistenza e unicità'

(per eq. diff. ordinarie: Se si trova una soluzione generale, quella è l'unica:

Funzione di prova:

$$u(t) = Ae^{\lambda t}$$

$$\rightarrow \frac{du}{dt} = A\lambda e^{\lambda t}$$

$$\rightarrow \frac{d^2u}{dt^2} = A\lambda^2 e^{\lambda t}$$

$$\rightarrow A\lambda^2 e^{\lambda t} + \omega_0^2 Ae^{\lambda t} = 0$$

$$\rightarrow \lambda^2 + \omega_0^2 = 0 \rightarrow \lambda = \pm i\omega_0$$

$$\rightarrow u_{\pm}(t) = A_{\pm} e^{\pm i\omega_0 t}$$

$$\rightarrow u(t) = A_+ e^{+i\omega_0 t} + A_- e^{-i\omega_0 t}$$

$$\rightarrow v(t) = \frac{du}{dt} = i\omega_0 A_+ e^{+i\omega_0 t} - i\omega_0 A_- e^{-i\omega_0 t} = i\omega_0 (A_+ e^{+i\omega_0 t} - A_- e^{-i\omega_0 t})$$

A_{\pm} costanti arbitrarie complesse (= 4 costanti reali)

Restrizione: $u(t), v(t)$ funzioni reali

$$\rightarrow \operatorname{Im} A_+ \sin(\omega_0 t) = -\operatorname{Im} A_- \sin(\omega_0 t) \rightarrow \operatorname{Im} A_+ = -\operatorname{Im} A_-$$

$$\rightarrow \operatorname{Re} A_+ \cos(\omega_0 t) = \operatorname{Re} A_- \cos(\omega_0 t) \rightarrow \operatorname{Re} A_+ = \operatorname{Re} A_-$$

$$\rightarrow A_+ = A_-^* = A = C_0 e^{i\varphi_0}$$

$$\rightarrow u(t) = Ae^{+i\omega_0 t} + A^* e^{-i\omega_0 t} = C_0 (e^{i\varphi_0} e^{+i\omega_0 t} + e^{-i\varphi_0} e^{-i\omega_0 t}) = 2C_0 \cos(\omega_0 t + \varphi_0)$$

$$\rightarrow v(t) = i\omega_0 (Ae^{+i\omega_0 t} - A^* e^{-i\omega_0 t}) = i\omega_0 C_0 (e^{i\varphi_0} e^{+i\omega_0 t} - e^{-i\varphi_0} e^{-i\omega_0 t}) = -2\omega_0 C_0 \sin(\omega_0 t + \varphi_0)$$

Cost. arbitrarie in termini delle cond. iniziali:

$$\left. \begin{array}{l} u(0) = a \rightarrow 2C_0 \cos(\varphi_0) = a \\ v(0) = b \rightarrow -2\omega_0 C_0 \sin(\varphi_0) = b \end{array} \right\} \rightarrow \tan(\varphi_0) = -\frac{b}{a\omega_0}, C_0 = \frac{1}{2} \sqrt{\frac{b^2}{\omega_0^2} + a^2}$$

Oscillatore armonico smorzato

Attrito viscoso:

$$F = -\beta v$$

Massa fissata a molla, in moto in un fluido viscoso:

$$F = -kx - \beta v$$

$$\rightarrow m \frac{d^2x}{dt^2} = -kx - \beta \frac{dx}{dt}$$

$$\rightarrow \frac{d^2x}{dt^2} + \frac{\beta}{m} \frac{dx}{dt} + \frac{k}{m} x = 0$$

$$\rightarrow \frac{d^2x}{dt^2} + 2 \underbrace{\frac{\beta}{2m}}_{\gamma} \frac{dx}{dt} + \underbrace{\frac{k}{m}}_{\omega_0^2} x = 0$$

Funzione di prova:

$$x(t) = A e^{\lambda t}$$

$$\rightarrow \frac{dx}{dt} = A \lambda e^{\lambda t}, \frac{d^2x}{dt^2} = A \lambda^2 e^{\lambda t}$$

$$\frac{d^2x}{dt^2} + 2\gamma \frac{dx}{dt} + \omega_0^2 x = 0 \rightarrow A \lambda^2 e^{\lambda t} + 2\gamma A \lambda e^{\lambda t} + \omega_0^2 A e^{\lambda t} = 0$$

$$\rightarrow \lambda^2 + 2\gamma\lambda + \omega_0^2 = 0$$

$$\rightarrow \lambda = -\gamma \pm \sqrt{\gamma^2 - \omega_0^2}$$

$$\rightarrow x(t) = A_1 e^{(-\gamma + \sqrt{\gamma^2 - \omega_0^2})t} + A_2 e^{(-\gamma - \sqrt{\gamma^2 - \omega_0^2})t}$$

Integrale generale

$$\omega_0^2 > \gamma^2$$

$$\begin{aligned} \sqrt{\gamma^2 - \omega_0^2} &= \pm i\sqrt{\omega_0^2 - \gamma^2} \\ \rightarrow x_{1,2}(t) &= A_{1,2} e^{(-\gamma \pm i\sqrt{\omega_0^2 - \gamma^2})t} = A_{1,2} e^{-\gamma t} e^{\pm i\sqrt{\omega_0^2 - \gamma^2}t} \\ \rightarrow x_{1,2}(t) &= A_{1,2} e^{-\frac{\beta}{2m}t} \left[\cos(\sqrt{\omega_0^2 - \gamma^2}t) \pm i \sin(\sqrt{\omega_0^2 - \gamma^2}t) \right] \end{aligned}$$

Comb. lineare:

$$x_1 + x_2 = e^{-\gamma t} \cdot \left\{ A_1 \left[\cos(\sqrt{\omega_0^2 - \gamma^2}t) + i \sin(\sqrt{\omega_0^2 - \gamma^2}t) \right] + A_2 \left[\cos(\sqrt{\omega_0^2 - \gamma^2}t) - i \sin(\sqrt{\omega_0^2 - \gamma^2}t) \right] \right\}$$

Imponendo che $x_1 + x_2$ sia reale:

$$\begin{aligned} \rightarrow x_1 + x_2 &= e^{-\gamma t} \left(\underbrace{A_1 + A_2}_{=B, \text{reale}} \right) \cos(\sqrt{\omega_0^2 - \gamma^2}t) + i \left(\underbrace{A_1 - A_2}_{=-iC, \text{immaginario}} \right) \sin(\sqrt{\omega_0^2 - \gamma^2}t) \\ \rightarrow x(t) &= e^{-\gamma t} \left[B \cos(\sqrt{\omega_0^2 - \gamma^2}t) + C \sin(\sqrt{\omega_0^2 - \gamma^2}t) \right] \end{aligned}$$

B, C determinati da condizioni iniziali

Oscillazioni smorzate

$$\omega_0^2 < \gamma^2$$

$$\rightarrow x_{1,2}(t) = A_{1,2} e^{(-\gamma \pm \sqrt{\gamma^2 - \omega_0^2})t}$$

Comb. lineare:

$$x_1 + x_2 = e^{-\gamma t} \left\{ A_1 e^{\sqrt{\gamma^2 - \omega_0^2} t} + A_2 e^{-\sqrt{\gamma^2 - \omega_0^2} t} \right\}$$

A_1, A_2 determinati da condizioni iniziali

No oscillazione; andamento smorzato

$$\omega_0^2 = \gamma^2$$

Caso intermedio fra i due: smorzamento critico

$$\frac{d^2x}{dt^2} + 2\gamma \frac{dx}{dt} + \omega_0^2 x = 0$$

$$\frac{d^2x}{dt^2} + \gamma \frac{dx}{dt} + \gamma \frac{dx}{dt} + \omega_0^2 x = 0$$

$$\frac{d^2x}{dt^2} + \gamma \frac{dx}{dt} = \frac{d}{dt} \left(\frac{dx}{dt} + \gamma x \right) \equiv \frac{dz}{dt}$$

$$\gamma \frac{dx}{dt} + \omega_0^2 x = \gamma \frac{dx}{dt} + \gamma^2 x = \gamma z$$

$$\rightarrow \frac{dz}{dt} + \gamma z = 0$$

$$\rightarrow z(t) = Ae^{-\gamma t}$$

$$\rightarrow \frac{dx}{dt} + \gamma x = Ae^{-\gamma t}$$

$$\rightarrow e^{\gamma t} \frac{dx}{dt} + e^{\gamma t} \gamma x = A$$

$$\rightarrow \frac{d}{dt} (e^{\gamma t} x) = A$$

$$\rightarrow e^{\gamma t} x = At + B$$

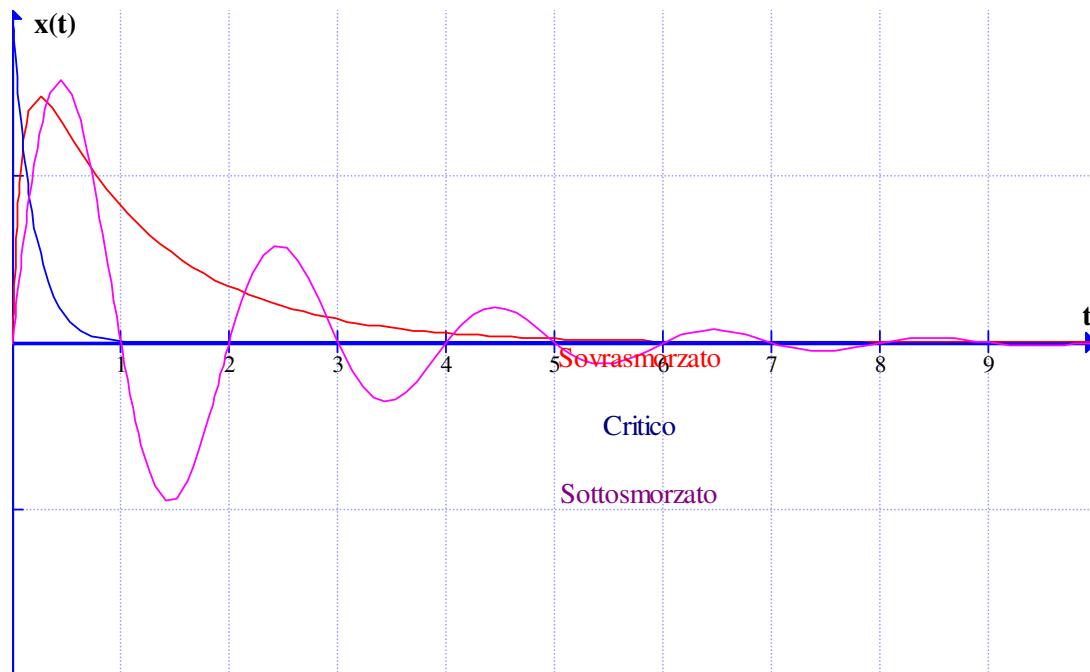
$$\rightarrow x(t) = Ate^{-\gamma t} + Be^{-\gamma t} = e^{-\gamma t} (At + B)$$

Smorzamento critico:

Minimo tempo di ritorno a zero

Come esempi:

Andamenti temporali nei 3 casi
(Condizioni iniziali arbitrarie e non specificate)



En. meccanica : Oscillazioni non smorzate

$$\frac{d^2x}{dt^2} + \omega_0^2 x = 0$$

$$\rightarrow \frac{dx}{dt} \frac{d^2x}{dt^2} + \omega_0^2 x \frac{dx}{dt} = 0$$

$$\rightarrow \frac{1}{2} \frac{d}{dt} \left(\frac{dx}{dt} \right)^2 + \frac{1}{2} \frac{d}{dt} (\omega_0^2 x^2) = 0$$

$$\rightarrow \frac{d}{dt} \left(\underbrace{\frac{1}{2} mv^2}_{E_k} + \underbrace{\frac{1}{2} kx^2}_U \right) = 0 \rightarrow E = E_k + U = \text{cost}$$

$$x(t) = A \sin \omega_0 t + B \cos \omega_0 t$$

$$\langle U \rangle_{\text{periodo}} = \left\langle \frac{1}{2} kx^2 \right\rangle = \frac{1}{2} k \left\langle (A \sin \omega_0 t + B \cos \omega_0 t)^2 \right\rangle$$

$$\langle E_k \rangle_{\text{periodo}} = \left\langle \frac{1}{2} mv^2 \right\rangle = \frac{1}{2} m \underbrace{\omega_0^2}_k \left\langle (A \cos \omega_0 t - B \sin \omega_0 t)^2 \right\rangle$$

$$\langle \cos^2 \omega_0 t \rangle = \langle \sin^2 \omega_0 t \rangle = \frac{1}{2}, \langle \sin \omega_0 t \cos \omega_0 t \rangle = 0$$

$$\rightarrow \langle U \rangle_{\text{periodo}} = \frac{1}{2} k \left(\frac{1}{2} A^2 + \frac{1}{2} B^2 \right) = \frac{1}{4} k (A^2 + B^2) = \langle E_k \rangle_{\text{periodo}}$$

$$E = E_k + U = \frac{1}{2} k (A^2 + B^2)$$

En. meccanica : Oscillazioni smorzate

$$\frac{d^2x}{dt^2} + \frac{\beta}{m} \frac{dx}{dt} + \frac{k}{m} x = 0$$

$$\frac{dx}{dt} \frac{d^2x}{dt^2} + \frac{\beta}{m} \left(\frac{dx}{dt} \right)^2 + \frac{k}{m} \frac{dx}{dt} x = 0$$

$$\frac{1}{2} \frac{d}{dt} \left(\frac{dx}{dt} \right)^2 + \frac{1}{2} \frac{k}{m} \frac{d}{dt} (x^2) = -\frac{\beta}{m} \left(\frac{dx}{dt} \right)^2$$

$$\frac{d}{dt} \left[\frac{1}{2} m \left(\frac{dx}{dt} \right)^2 \right] + \frac{d}{dt} \left(\frac{1}{2} kx^2 \right) = -\beta \left(\frac{dx}{dt} \right)^2$$

$$\rightarrow \frac{dE}{dt} = -\beta \left(\frac{dx}{dt} \right)^2 < 0$$

$$\rightarrow \frac{dE}{dt} = -\frac{\beta}{m} m \left(\frac{dx}{dt} \right)^2 < 0$$

v smorzata esponenzialmente

→ E tende a 0

Eccitazione esterna dell'oscillatore: termine forzante

Es: sinusoidale con frequenza angolare Ω

$$\frac{d^2x}{dt^2} + 2\frac{\beta}{2m}\frac{dx}{dt} + \frac{k}{m}x = \frac{F(t)}{m} = \frac{F_0}{m}\cos\Omega t$$

$$\rightarrow \frac{d^2x}{dt^2} + 2\gamma\frac{dx}{dt} + \omega_0^2x = \frac{F(t)}{m} = \frac{F_0}{m}\cos\Omega t$$

Formalismo esponenziali complessi:

variabile ausiliaria 'posizione complessa': $\tilde{x}(t) = \alpha e^{i\Omega t}$

$$\rightarrow x(t) = \operatorname{Re} \tilde{x}(t) = \operatorname{Re}(\alpha e^{i\Omega t})$$

variabile ausiliaria 'forza complessa': $\tilde{F}(t) = F_0 e^{i\Omega t}$

$$\rightarrow F(t) = \operatorname{Re} \tilde{F}(t)$$

$$\rightarrow \frac{d^2\tilde{x}}{dt^2} + 2\gamma\frac{d\tilde{x}}{dt} + \omega_0^2\tilde{x} = \frac{\tilde{F}}{m}$$

$$\rightarrow \alpha(-\Omega^2)e^{i\Omega t} + 2i\alpha\gamma\Omega e^{i\Omega t} + \alpha\omega_0^2 e^{i\Omega t} = \frac{F_0}{m} e^{i\Omega t}$$

$$\rightarrow \alpha[-\Omega^2 + 2i\gamma\Omega + \omega_0^2] = \frac{F_0}{m}$$

$$\rightarrow \alpha = \frac{F_0}{m} \frac{1}{(\omega_0^2 - \Omega^2) + 2i\gamma\Omega} \text{ funzione complessa di } \Omega$$

$$x(t) = \operatorname{Re} \tilde{x}(t) = \operatorname{Re}(\alpha e^{i\Omega t})$$

$$\alpha \equiv |\alpha| e^{i\phi}$$

$$\rightarrow x(t) = \operatorname{Re}(\alpha e^{i\Omega t}) = \operatorname{Re}(|\alpha| e^{i(\Omega t + \phi)}) = |\alpha| \cos(\Omega t + \phi)$$

$$\rightarrow \alpha = \frac{F_0}{m} \frac{1}{(\omega_0^2 - \Omega^2) + 2i\gamma\Omega} = F_0 \frac{(\omega_0^2 - \Omega^2) - 2i\gamma\Omega}{(\omega_0^2 - \Omega^2)^2 - (2i\gamma\Omega)^2}$$

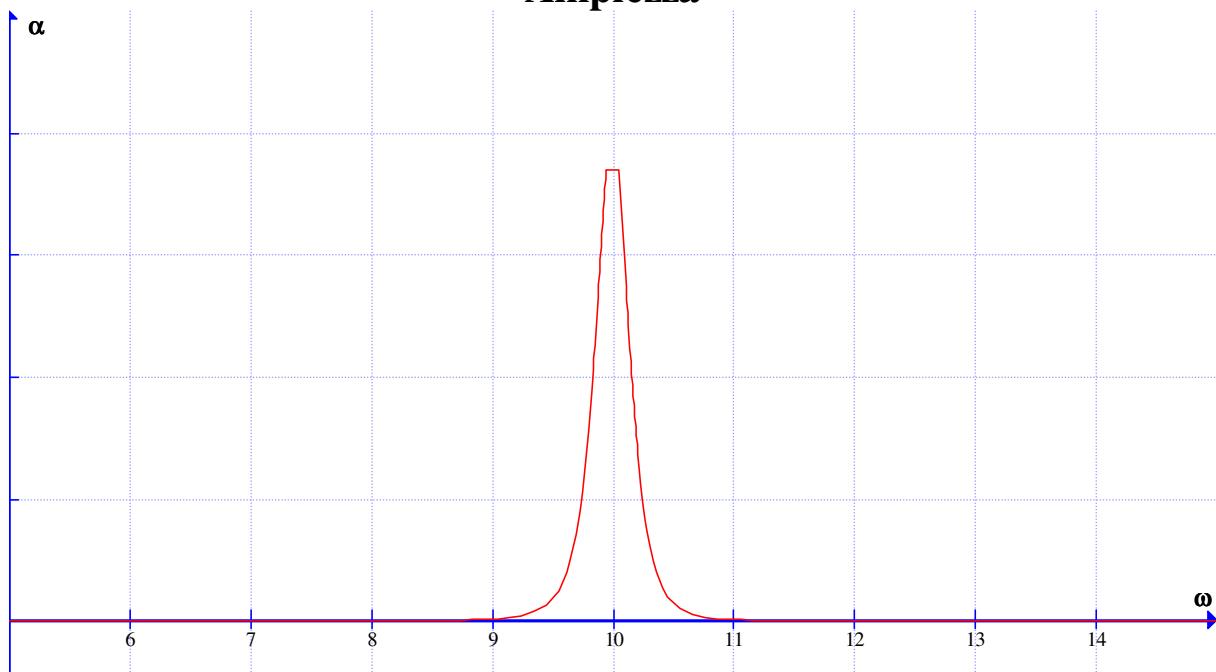
$$|\alpha|^2 = \frac{F_0^2}{m^2} \frac{[\omega_0^2 - \Omega^2]^2 + (2\lambda\Omega)^2}{[(\omega_0^2 - \Omega^2)^2 + 4\lambda^2\Omega^2]^2} = \frac{F_0^2}{(\omega_0^2 - \Omega^2)^2 + 4\lambda^2\Omega^2}$$

$$\left. \begin{aligned} |\alpha| &= \frac{F_0}{m\sqrt{(\omega_0^2 - \Omega^2)^2 + 4\lambda^2\Omega^2}} \quad \text{modulo} \\ \phi &= \arctan \frac{2\lambda\Omega}{\Omega^2 - \omega_0^2} \quad \text{fase} \end{aligned} \right\} \text{funzioni di } \Omega$$

$$\mathbf{Max} \ |\alpha| : \omega_0 = \Omega$$

$$|\alpha|_{\max} = \frac{F_0}{\lambda m \omega_0}, \quad \phi_{\max} = -\frac{\pi}{2}$$

Aampiezza



Fase

